# Asteroid 45 Eugenia: Lightcurves and the Pole Orientation ${ }^{1}$ 

R. C. TAYLOR,* P. V. BIRCH, $\dagger$ A. POSPIESZALSKA-SURDEJ, $\ddagger$ and J. SURDEJ $\dagger \S$<br>*Lunar and Planetary Laboratory, University of Arizona, Tucson, Arizona 85721; $\dagger$ Perth Observatory, Bickley, Western Australia; $\ddagger$ Institut d’Astrophysique, Université de Liège, Liège, Belgium; and §Chercheur Qualifié au Fonds National de la Recherche Scientifíque, Belgium

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#### Abstract

Nine lightcurves of asteroid 45 Eugenia, three from 1969 and six from 1984, are given. In 1984-1985 the $\boldsymbol{H}_{0}$ magnitude of Eugenia, corrected to the lightcurve maximum, was 7.47 and the slope parameter $\boldsymbol{G}_{0}$ was $\mathbf{0 . 0 4}$. The north pole of Eugenia is within $\pm 10^{\circ}$ of ecliptic longitude $106^{\circ}$ and latitude $+26^{\circ}$ (or $295^{\circ}$ and $+34^{\circ}$ ). This solution is consistent with an amplitude-aspect pole analysis. The sidereal period is $0.2374645 \pm 0.0000002$ day, or $5 \mathrm{hr} 41 \mathrm{~min} 56.93 \mathrm{sec} \pm 0.02$ sec and the sense of rotation is retrograde. When observations are closest to both the north and south poles ( $\sim 30^{\circ}$ ) only one maximum and one minimum are present in the lightcurves; at other oppositions there are two of each. It is suggested that this is caused by albedo features on the surface of Eugenia. © 1988 Acsedernic Press, Inc.


## I. INTRODUCTION

Vesely and Taylor (1985) published the remaining useful lightcurves from our files with a comment that lightcurves of 45 Eu genia would be published separately. Those lightcurves, three from 1969, are given in this paper along with five lightcurves from 1984. We give Eugenia's 1984-1985 absolute magnitude $H_{0}$ and slope parameter $G_{0}$ derived from lightcurve maxima. We determine the asteroid's pole orientation, sidereal period, and sense of rotation. This paper represents continuing research in the application of photometric astrometry (PA) to asteroids. PA is a method for determining the pole orientation, sidereal period, and sense of rotation of an asteroid. The most recent capsule summary of PA can be found in Section I of Taylor et al. (1987). PA is explained in detail in Taylor (1979) and Taylor and Tedesco (1983). Our PA solution is compared to results of an ampli-tude-aspect pole analysis.

[^0]Eugenia is a $250-\mathrm{km}$ U-type asteroid (Bowell et al. 1979). It is also grouped as a LASPA (large amplitude and short period) asteroid which leads Farinella et al. (1981) to suggest that Eugenia's shape may be a Jacobi ellipsoid in rotational equilibrium.

Eugenia was observed on four consecutive nights in May 1978 by Debehogne and Zappaià (1980). They refined the synodic period to 5 hr 41 min 56 sec , using a lightcurve obtained a month later by Harris and Young (1979). Eugenia was observed during the 1982 opposition by Debehogne et al. (1983) and by Weidenschilling et al. (1987). The latter group also obtained Eugenia lightcurves in the 1981-1982, 1983, 19841985, and 1985-1986 oppositions.

## II. THE OBSERVATIONS

Figure 1 gives the Eugenia lightcurves from 1969. Figure 2 is a composite of those lightcurves using a 5 hr 42 min period. Figures 3 and 4 are 1984 lightcurves. All lightcurves are in Universal Time not corrected for light time. The vertical scale is the differential $V$ magnitude in the sense of asteroid minus comparison star normalized to


Fig. 1. 1969 lightcurves of Eugenia. $V_{0}(1, \alpha)=8.37$ for June 9 and 8.34 for June 11.
zero at the lightcurve maximum. Table I Table III gives the observed $V$ magnigives information about the comparison tudes and colors of Eugenia in 1969 and stars, the observers, and the telescopes they used. Table II gives the aspect data at the midtime of each lightcurve. 1984. $V_{0}(1, \alpha)$ represents each $V$ magnitude corrected to both its lightcurve maximum and to unit distance from the Sun and the


Fig. 2. Composite of the Fig. 1 lightcurves. The horizontal time scale in hours does not pertain to any specific night but is representative of one rotation period. Open circles are repetitions of points just one rotational cycle earlier or later.



Fig. 3. 1984 lightcurves of Eugenia. In the upper lightcurve open circles are from September 28 and filled circles are from September 29. $V_{0}(1, \alpha)=8.23$ for September $28,8.22$ for September 29, and 7.85 for October 24.

Earth. The absolute magnitude $H_{0}$ and the slope parameter $G_{0}$ were calculated using the 1984-1985 observations. Data from other oppositions are not included because precise axial ratios of Eugenia are not known and therefore aspect corrections cannot be applied accurately. The data are from Table III and the results, using the method of Tedesco (1986), are $H_{0}=7.47$ and $G_{0}=0.04$. Figure 5 shows the $1984 V$ phase relation of Eugenia.

## III. PHOTOMETRIC ASTROMETRY

The details of the photometric astrometry method will not be reproduced in this
paper (see Section I for PA references). However, two refinements to PA are adopted and now discussed. First, in calculating the distance an asteroid moves across the sky $(\Delta \phi)$ we use the distance between the phase angle bisectors rather than the distance between the sub-Earth points. The phase angle bisector concept was introduced by Harris et al. (1984). This routine is used in calculating the asteroid's angular velocity $(\Delta \phi / \Delta t)$ across the sky in the estimated sidereal period analysis. Second, phase angle bisectors also replace "time shifts'" in the basic PA Eq. (1) of Taylor and Tedesco (1983). Both of these refinements

TABLE I
Comparison Stars and Observers

| Fig. <br> No. | UT Date | Observer | Comparison star | $\begin{gathered} \text { Observed } \\ V \\ (\mathrm{mag}) \end{gathered}$ | Observed $B-V$ | Observed $U-B$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1969 Jun 9 | D | $a$ | $10.10 \pm 0.02$ | $+0.63 \pm 0.02$ |  |
| 2 | 1969 Jun 10 | D | $b$ |  |  |  |
| 3 | 1969 Jun 11 | D | c | $12.04 \pm 0.03$ | $+0.82 \pm 0.01$ |  |
| 5 | 1984 Sep 28 | S | ${ }^{\text {d }}$ | $10.89 \pm 0.02$ | $+0.92 \pm 0.01$ | $+0.54 \pm 0.01$ |
| 5 | 1984 Sep 29 | S | ${ }^{\text {d }}$ | $10.89 \pm 0.02$ | $+0.92 \pm 0.01$ | $+0.54 \pm 0.01$ |
| 6 | 1984 Oct 24 | B | $\mathrm{BD}+2348$ | $10.03 \pm 0.01$ |  |  |
| 7 | 1984 Oct 31 | B | $\mathrm{BD}+2348$ | $10.08 \pm 0.04$ |  |  |
| 8 | 1984 Nov 21 | B | HD 12923 | $6.30 \pm 0.02$ |  |  |
| 9 | 1984 Nov 27 | B | HD 11037 | $5.92 \pm 0.02$ |  |  |

Note. D, Dunlap at the Steward $91-\mathrm{cm}$ telescope on Kitt Peak presently housing the Spacewatch Telescope; S, A. Pospieszalska-Surdej and J. Surdej at the European Southern Observatory 50-cm telescope. The observing and reduction procedures are described in Surdej et al. (1983); B, P. V. Birch at the Lowell-Perth $61-\mathrm{cm}$ reflector.
${ }^{a}$ Comparison star not catalogued; 1950 coordinates are $\mathrm{RA}=14 \mathrm{hr} 36.6 \mathrm{~min}$, and Decl. $=-4^{\circ} 45^{\prime}$.
${ }^{b}$ Same star June 9, 1969. No photometry done.
${ }^{c}$ Identification not available.
${ }^{d}$ Comparison star not catalogued; 1950 coordinates are RA $=2 \mathrm{hr} 24 \mathrm{~min} 19 \mathrm{sec}$ and Decl. $=$ $+5^{\circ} 42^{\prime} 20^{\prime \prime}$.
were used in the study of Herculina (Taylor et al. 1987). The results from the new and old techniques give identical results. However, we recognize that it is possible, given rarely seen relative positions of an asteroid with respect to its pole position, that discrepancies between the techniques could occur. The phase angle bisector is the more
appropriate concept and therefore it is now a part of PA. Because of this change, the formulas (14) for the cycle corrections found in the Appendix of Taylor (1979) should be modified; the coordinates of the phase angle bisector rather than the subearth point should be used.

Table IV lists the time intervals of the

TABLE II

| Aspect Data for Eugenia |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Observed UT date | RA | Decl. | Distance (AU) from the |  | Phase angle | Ecliptic |  |
|  |  |  |  |  |  | Long | Lat |
|  |  |  | Sun | Earth |  |  |  |
| 1969 Jun 9 | $14^{\text {h }} 35 \mathrm{~m} .5$ | $5-4^{\circ} 41^{\prime}$ | 2.498 | 1.642 | +15.47 | 218.0 | $+10.0$ |
| 1969 Jun 10 | 1435.5 | $5-441$ | 2.498 | 1.645 | +15.61 | 218.0 | +10.0 |
| 1969 Jun 11 | 1435.6 | $6-440$ | 2.498 | 1.649 | +15.75 | 218.0 | $+10.0$ |
| 1984 Sep 28 | 222.7 | +537 | 2.916 | 2.014 | -10.31 | 35.3 | $-8.1$ |
| 1984 Sep 29 | 222.1 | $1+532$ | 2.917 | 2.008 | - 9.98 | 35.1 | $-8.1$ |
| 1984 Oct 24 | 23.5 | $5+36$ | 2.927 | 1.940 | - 3.03 | 29.9 | $-8.9$ |
| 1984 Oct 31 | 157.7 | $7+230$ | 2.929 | 1.954 | + 4.42 | 28.3 | $-8.9$ |
| 1984 Nov 21 | 143.2 | $2+123$ | 2.935 | 2.075 | +11.27 | 24.4 | - 8.7 |
| 1984 Nov 27 | 140.5 | $5+117$ | 2.937 | 2.129 | +12.99 | 23.7 | $-8.5$ |

TABLE III
Magnitudes and Colors of Eugenia

| Observed <br> UT date | $V_{0}(1, \alpha)$ <br> $(\mathrm{mag})$ | $B-V$ <br> $(\mathrm{mag})$ | $U-B$ <br> $(\mathrm{mag})$ |
| :--- | :---: | :---: | :---: |
| 1969 Jun 9 | $8.37 \pm 0.03$ | $+0.68 \pm 0.02$ | $+0.29 \pm 0.02$ |
| 1969 Jun 11 | $8.34 \pm 0.03$ | $+0.67 \pm 0.01$ | $+0.27 \pm 0.01$ |
| 1984 Sep 28 | $8.23 \pm 0.03$ | $+0.65 \pm 0.02$ | $+0.24 \pm 0.02$ |
| 1984 Sep 29 | $8.22 \pm 0.03$ | $+0.67 \pm 0.02$ | $+0.23 \pm 0.02$ |
| 1984 Oct 24 | $7.85 \pm 0.02$ | - | - |
| 1984 Oct 31 | $7.86 \pm 0.04^{a}$ | - | - |
| 1984 Nov 21 | $8.28 \pm 0.03$ | - | - |
| 1984 Nov 27 | $8.31 \pm 0.03$ | - | - |
| 1985 Jan 17 | $8.62^{b}$ |  |  |

${ }^{a}$ Since the lightcurve maximum was not observed a 0.05 mag correction was applied (based upon an overlay of the Oct 24 and Oct 31 lightcurves).
${ }^{b}$ Observed at $19.50^{\circ}$ solar phase angle. From Weidenschilling et al. (1987).


Fig. 4. 1984 lightcurves of Eugenia. $V_{0}(1, \alpha)=7.86$ for October $31,8.28$ for November 21, and 8.31 for November 27.


FIG. 5. The $V$ phase relation of Eugenia at maximum light. The solid curve is the least-squares fitted Bowell-Lumme-Harris phase relation (Bowell et al. 1987).

1984-1985 lightcurves that were used to determine the "estimated sidereal period." Note that the times for January 17, 1985 are from the lightcurve of Weidenschilling et al. (1987). The synodic period when the angular velocity of the phase angle bisector ( $\Delta \phi / \Delta t$ ) is zero is the sidereal period. Therefore, the least-squares solution of the ordered pairs ( $\Delta \phi / \Delta t$, synodic period) gives the estimated sidereal period: 0.237464 day. The routine was done 10 additional times by varying the time intervals randomly by their estimated errors (the right column of Table IV). The results indicate that the er-
ror in the estimated sidereal period is $\pm 0.000004$ day $(1 \sigma)$. A negative slope from the least-squares routine indicates that if the asteroid is in the apparent retrograde loop of its orbit then the observed synodic periods are each larger than the sidereal period; a characteristic of a retrograde sense of rotation. The slope of the Eugenia solution is $-0.00018 \pm 0.00004$, clearly implying that the sense of rotation of Eugenia is retrograde.

The mean synodic period (MSP) is a constant which represents the mean of all synodic periods over the orbit of an asteroid. It enables the number of synodic cycles to be counted between similar lightcurve features over long time intervals. Eugenia's lightcurves usually have two maxima and two minima per rotation period of 5 hr 41.9 min which do not create any difficulties in finding the MSP. However, in 1983 and 19851986 Weidenschilling et al. (1987) observed the Eugenia lightcurve to have essentially only one maximum and one minimum. Those lightcurves have the smallest observed amplitudes which imply that the observations were the furthest from the equator of Eugenia. The difference in ecliptic longitude between the two oppositions is approximately $180^{\circ}$ so our first thought was that the single lightcurve minimum from

TABLE IV
Intervals from the 1984-1985 Data Used to Find the Estimated
Sidereal Period

| Date | UT—Date | UT | $\Delta \phi / \Delta t$ | Synodic <br> periods <br> (days) | Estimated <br> error in overlays <br> ( $\pm$ min) |
| :--- | ---: | :---: | :---: | :---: | :---: |
| Oct 24 | $16: 00$ —Sep 29 | $6: 11$ | -0.0207 | 0.237475 | 4 |
| Oct 31 | $15: 08$ —Sep 29 | $8: 00$ | -0.0190 | 0.237479 | 3 |
| Nov 21 | $15: 00$ —Sep 29 | $9: 41$ | -0.0108 | 0.237466 | 3 |
| Nov 27 | $14: 00$ —Sep 29 | $5: 12$ | -0.0083 | 0.237464 | 4 |
| Jan 17 | $5: 00$ —Sep 29 | $6: 17$ | +0.0551 | 0.237456 | 2 |
| Oct 24 | $14: 00$ —Nov 21 | $14: 29$ | -0.0018 | 0.237453 | 4 |
| Oct 24 | $15: 00$ —Nov 27 | $14: 00$ | +0.0057 | 0.237464 | 3 |
| Jan 17 | $5: 00$ —Oct 24 | $16: 05$ | +0.0745 | 0.237453 | 5 |
| Jan 17 | $7: 00$ —Oct 31 | $15: 29$ | +0.0824 | 0.237437 | 3 |
| Jan 17 | $5: 00$ —Nov 27 | $15: 03$ | +0.1207 | 0.237453 | 3 |

TABLE V
Epochs and Parameters Used in Photometric Astrometry

| $\begin{aligned} & \text { UT } \\ & \text { date } \end{aligned}$ | Minimum $m 1$ |  | Minimum $m 2$ |  | Use | Source | Aspect ${ }^{a}$ <br> (degree) | Ampl. <br> (mag) | Cycle corr. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | ID | UT | ID | UT |  |  |  |  |  |
| Jun 91969 | - | - | 1 | 5:53 | MSP | T | 75.2 | $0.34{ }^{\text {b }}$ | - |
| Jun 101969 | 2 | 7:30 | - | - | MSP, POLE | T | 75.2 | $.34{ }^{\text {b }}$ | -1 |
| May 41978 | - | - | 3 | 8:52 | POLE | D | 68.4 | . 31 | +1 |
| Jun 11978 | 4 | 6:30 | 5 | 9:26 | MSP | H | 72.9 | . $34{ }^{\text {c }}$ | - |
| May 211982 | 6 | 6:05 | 7 | 8:44 | MSP, POLE | W | 123.0 | . 20 | +2 |
| Jun 301983 | 8 | 7:01 | - | - | MSP, POLE | W | 34.0 | . 11 | +3 |
| Oct 311984 | 10 | 15:38 | - | - | MSP, POLE | T | 96.9 | - ${ }^{d}$ | +3 |
| Nov 271984 | - | - | 9 | 14:32 | MSP, POLE | T | 93.1 | . 45 | +3 |
| Oct 201985 | - | - | 11 | 3:03 | MSP | W | 146.3 | 0.15 | - |

Note. T, this paper; D, Debehogne and Zappalà (1980); H, Harris and Young (1979); W, Weidenschilling et al. (1987).
${ }^{a}$ As measured from the $106^{\circ}$ north pole solution.
${ }^{b}$ The amplitude is from the composite lightcurve (Fig. 2).
${ }^{c}$ From the composite lightcurve of 4 nights in Fig. 1 of Harris and Young (1979).
${ }^{d}$ The lightcurve does not cover the full rotation period.
each opposition represented the same surface feature on the asteroid. This is not true. In the MSP analysis it is apparent that the 1983 and 1985 lightcurves are switched with respect to each other. That is, the feature on the asteroid causing the single minimum in the 1983 lightcurve is not the same feature causing the single minimum in the 1985-1986 lightcurve. The two minima are separated by approximately $180^{\circ}$ in rotational phase. This same phenomenon may occur with asteroid Herculina; a lightcurve switch is predicted for 1988 (see Section V of Taylor et al. 1987).

Table V gives the Eugenia lightcurve epochs of minimum light used in the derivation of the MSP. The epochs with even number identifications include 1983 and are designated $m 1$. The odd number identifications $m 2$, which include 1985 , are one-half a rotation cycle from the $m 1$ set. In the MSP analysis a time interval between two epochs from the independent sets $m 1$ and $m 2$ is used only if the epochs are within $20^{\circ}$ solar phase angle of each other (see Section III of Taylor and Tedesco 1983), and the time interval is longer than the orbital period of Eugenia ( $\sim 1640$ days). The resulting
intervals, from Table V, are 2-4, 2-6, 2-10, $4-8,4-10,1-5,1-7,1-9,1-11,5-9$, and $5-11$. Table V also identifies the epochs which are used in the pole analysis and a footnote column giving the source of each lightcurve epoch. The last three columns of the table are explained below.

By examining a few test cases it was found that, for retrograde rotation, a MSP is 0.000035 day smaller than the sidereal period it generates. Since the estimated sidereal period is 0.237465 day $\pm 0.000004(1 \sigma)$ then a search was made for a MSP within $3 \sigma$ of the expected result, namely between 0.237418 and 0.237442 day. Only one MSP exists in that domain. The MSP of Eugenia, from both sets $m 1$ and $m 2$, is $0.2374296 \pm$ $0.0000003(1 \sigma)$ day, within $1 \sigma$ of the estimated value.

Table VI gives the time intervals, corrected for light time, used in determining the pole orientation and sidereal period of Eugenia. The identification numbers are the same as those found in Table V. The number of cycles is the quotient of the time interval and the MSP. The lightcurves were overlaid with similar minima superposed. Maxima are not used in this analysis be-

TABLE VI
Time Intervals Used to Determine the Pole Orientation and Sidereal Period of Eugenia

| ID of <br> interval | Time <br> interval <br> (days) | Number of <br> cycles | Error <br> (min) |
| :---: | :---: | :---: | ---: |
| $2-4$ | 3277.9588 | 13806 | $\pm 3$ |
| $2-6$ | 4727.9381 | 19913 | 4 |
| $2-8$ | 5132.9757 | 21619 | 5 |
| $2-10$ | 5622.3371 | 23680 | 3 |
| $4-8$ | 1855.0214 | 7813 | 4 |
| $4-10$ | 2344.3783 | 9874 | 4 |
| $1-3$ | 3251.1252 | 13693 | 2 |
| $1-7$ | 4729.1159 | 19918 | 5 |
| $1-9$ | 5650.3577 | 23798 | 3 |
| $1-11$ | 5977.2923 | 25175 | 5 |
| $3-9$ | 2399.2324 | 10105 | 5 |
| $3-11$ | 2726.1671 | 11482 | $\pm 5$ |

cause of the anomalies in the 1983 and 1985 lightcurves. The errors, in Table VI, are estimated uncertainties of each fit. We also estimate that the error in assuming lightcurve minima represent the same feature on the asteroid to be $\pm 2^{\circ}$ rotational phase, which for Eugenia is $\pm 3 \mathrm{~min}$. Equation (1) of Taylor and Tedesco (1983) was applied 20 times and for each trial the time interval was altered randomly by both uncertainties discussed above. PA results are given below; the "formal" pole error is the mean of the 20 angular differences between each trial pole and the adopted pole. The "approximate" uncertainty is our estimate based on the formal error.


In Table V the last column gives the cycle corrections needed for the $106^{\circ}$ pole. It is an additional synodic cycle which is
added for each orbital rotation of the asteroid (see Eq. (1) and Section V of Taylor and Tedesco 1983). The cycle correction is only listed for the epochs used in the pole analysis. In Table $V$ the third from the last column gives the aspect angle which is measured between the $106^{\circ}$ pole of the asteroid and the line of sight. The next column gives the lightcurve amplitude. A plot of amplitudes versus aspect (not shown) is internally consistent; they vary directly (see the next section).

## IV. THE AMPLITUDE-ASPECT POLE SOLUTIONS

Combining all available photometric lightcurves of 45 Eugenia that have been recorded at phase angles $\leq 15^{\circ}$ (cf. Table VII), we have applied the "revisited ampli-tude-aspect relation" in order to test the consistency of the pole solutions derived by PA. When doing so, we have first digitized all published "complete" lightcurves and measured the slope $D_{i}$ of the normalized curves ( $i=1$ to 8 in Table VII) representing $I_{\mathrm{r}}^{2}$ versus $\cos ^{2}(\psi)$, where $I_{\mathrm{r}}$ is the relative intensity of a measurement observed at phase $\psi$. Under the assumption of a threeaxis ellipsoid model, PSS $^{2}$ have shown that a simple but nonlinear relation does hold between the observed $D_{i}$, the ecliptic coordinates $\lambda_{i}, \beta_{i}$ of the minor planet and the pole coordinates $\left(\lambda_{0}, \beta_{0}\right)$ as well as the semi-axes ratios $a / b$ and $b / c$ of the bestfitted ellipsoid model. In this context, a set of at least four independent nonlinear equations must then be solved in order to determine the values of the four unknown parameters $\lambda_{0}, \beta_{0}, a / b$, and $b / c$. Since observations made at ecliptic longitudes $180^{\circ}$ apart are photometrically equivalent, Table VII indicates however that 45 Eugenia has only been observed within three dis-

[^1]TABLE VII
Aspect Information and Characteristics of the $I_{\mathrm{r}}^{2}-\cos ^{2}(\psi)$ Relation Used in the Amplitude-Aspect Analysis of Eugenia

| Observed UT date | Ref. ${ }^{\text {a }}$ | Ecliptic |  | Phase angle | $t$ |  | D | $R$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | Long (1 | ${ }_{(0)}^{\text {Lat }}$ |  |  |  |  |  |
| 1969 Jun 10 | PP | $217 \times 92$ | +10.02 | 15.7 | 1. 28567 | 0.68053 | $\pm 0.03977$ | 0.932 |
| 1978 May 4 | DZ | 226.49 | +10.79 | 4.5 | 4.60353 | 0.56598 | 0.01513 | 0.974 |
| 1978 Jun 1 | HY | 221.04 | + 10.51 | 12.0 | 8.01975 | 0.57853 | 0.03847 | 0.967 |
| 1983 Jun 30 | W | 294.81 | + 6.82 | 12.3 | $8.5{ }^{\text {b }}$ | 0.07847 | 0.02522 | $0.511^{\text {b }}$ |
| 1984 Sep 29 | PP | 35.07 | -8.16 | 9.9 | 7.14971 | 0.95478 | 0.02869 | 0.978 |
| 1984 Nov 21 | PP | 24.44 | -8.69 | 11.3 | 14.53327 | 1.05685 | 0.03457 | 0.991 |
| 1984 Nov 27 | PP | 23.75 | -8.53 | 13.0 | 13.08204 | 1.21469 | 0.05085 | 0.988 |
| 1986 Jan 17 | W | 116.65 | - 5.26 | 1.9 | $9.75{ }^{\text {b }}$ | 0.07416 | 0.03216 | $0.592^{\text {b }}$ |

${ }^{a}$ References: PP, Present paper; DZ, Debehogne and Zappalà (1980); HY, Harris and Young (1979); W, Weidenschilling et al. (1987).
${ }^{b}$ This lightcurve seriously departs from an ellipsoidal one (only one maximum and minimum).
tinct longitude ranges, namely $\lambda_{i} \in[210-$ $\left.230^{\circ}\right], \lambda_{i} \in\left[20-40^{\circ}\right]$, and $\lambda_{i} \simeq 295^{\circ}$. As a consequence (see PSS), it is not possible to determine unambiguously the values of the four parameters $\lambda_{0}, \beta_{0}, a / b$, and $b / c$ because one independent observation is missing. We have therefore chosen to calculate


Fig. 6. Predicted amplitude-aspect pole solutions $\left(\lambda_{0}, \beta_{0}\right.$, and $\left.a / b\right)$ as a function of the semi-axes ratio $b / c$ (see Section IV). The PA pole solutions (see Section III) are indicated with crosses.
the solution of just three parameters ( $\lambda_{0}, \beta_{0}$, $a / b)$ as a function of the fourth one $(b / c)$. While performing the calculations, we have furthermore assigned equal weight to each group of observations pertaining to one of the three distinct oppositions. The results of these calculations are illustrated in Fig. 6 where we have also indicated the PA pole solutions. As most usually (see PSS), we also find that there exist two equally probable sets of pole solutions ( $P_{1}$ and $P_{2}$ ). Due to the nonellipsoidal character of some of the observed photometric lightcurves (cf. those recorded in 1983 and 1986), we conclude that there is an overall good agreement between the predicted PA and AA solutions.

## V. DISCUSSION

Lightcurves of Eugenia have been obtained from six oppositions; 1969, 1978, 1981-1982, 1983, 1984-1985, and 19851986. In four of the oppositions the lightcurves rather typically have two maxima and two minima. The relative amplitudes of the extrema do vary, but that in itself is not unusual. However, in 1983 and 1985-1986 the lightcurve shape changes. There are secondary variations but essentially the
lightcurves have just one maximum and one minimum per rotation cycle (see Weidenschilling et al. 1987). The lightcurves of 1983 and 1985-1986 were both obtained by viewing approximately $30^{\circ}$ from the $106^{\circ}$ north pole solution and its corresponding south pole, respectively. Also, from the PA analysis we conclude that the lightcurves from those two oppositions are switched with respect to each other. That is, the feature on the asteroid causing the maxima in 1983 is causing the minima in 1985-1986, and vice versa. These phenomena were observed by Taylor et al. (1987) in the study of asteroid 532 Herculina. They suggested, as the simplest model, that Herculina might be a spheroidal body with two dark regions separated by approximately $180^{\circ}$ in longitude. They demonstrated that such a model successfully reproduces the Herculina lightcurve amplitudes of each extrema. The Eugenia lightcurve anomalies might be explained by a similar model. Debehogne and Zappalà (1980) also suggest that Eugenia might have an albedo feature. Their 1978 lightcurves display maxima with nearly the same brightness but minima which reach different levels. The albedo conjecture is based on the fact that those 1978 observations were made at less than $5^{\circ}$ phase angle. We urge that the possible albedo structure of Eugenia be tested with simultaneous visual and thermal infrared observations. Radar observations would also aide in the study of Eugenia's shape.

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[^0]:    ${ }^{1}$ Based in part on observations collected at the European Southern Observatory, La Silla, Chile.

[^1]:    ${ }^{2}$ We refer the reader to the work by PospieszalskaSurdej and Surdej (1985), referred to hereafter as PSS, for a comprehensive description of the amplitude-aspect (AA) method as well as for the exact meaning of the parameters $t_{\mathrm{M}}, D, R$, and others used in the remainder of this section.

