# THE DISSOCIATIVE ELECTROIONIZATION OF AMMONIA AND AMMONIA- $d_3$ . II. THE H<sup>+</sup> AND H<sub>2</sub><sup>+</sup> DISSOCIATION CHANNELS

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#### Abstract

The dissociative electroionization of NH<sub>3</sub> (ND<sub>3</sub>), in the H<sup>+</sup> (D<sup>+</sup>) and H<sub>2</sub><sup>+</sup> (D<sub>2</sub><sup>+</sup>) dissociation channels, is investigated in the 18-50 eV electron energy range. The H<sup>+</sup> (D<sup>+</sup>) translational energy distribution spectrum shows several components and that of H<sub>2</sub><sup>+</sup> (D<sub>2</sub><sup>+</sup>) is continuous. The thermal energy component of the latter distribution is entirely ascribed to H<sub>2</sub><sup>+</sup> /H<sub>2</sub>. For both ions the kinetic energy versus appearance energy diagram is obtained and discussed. The H<sup>+</sup> ion at the lowest onset is probably formed at the expense of NH<sub>3</sub><sup>+</sup> ( $\tilde{A}^2E$ ) state. In the 22.5 eV electron energy range both H<sup>+</sup> and H<sub>2</sub><sup>+</sup> are formed by dissociative autoionization. Evidence is also brought for H<sup>+</sup> and H<sub>2</sub><sup>+</sup> formation at 22.5 and 26.5 eV through spontaneous decomposition of NH<sub>2</sub><sup>+</sup>. Above 35 eV dissociation of doubly ionized states by Coulomb repulsion is mainly involved.

### 1. Introduction

Recently, the results on the dissociative ionization of  $NH_3$  in the  $NH_2^+$  and  $NH^+$  dissociation channels have been reported [1]. Translational energy distribution measurements, together with appearance energy determinations allowed us to discuss in detail most of the dissociation processes.

The dissociation of  $NH_3$  into  $H^+$  and  $H_2^+$  has been investigated by a few groups. Appearance energies were measured but very briefly discussed without kinetic energy measurements. The earliest work of Mann et al. [2] and the most recent publication of Mark et al. [3] included both ions. Only the former discussed dissociation processes, while the latter only mentioned onset energies.

Using energy-loss electron spectroscopy and electron-ion coincidence technique, Wight et al. [4] and Brion et al. [5] determined the fragmentation pattern of  $NH_3$ , including the  $H^+$  dissociation channels.

More recently Carnahan et al. [6] and Kurawaki and Ogawa [7] published their results on high Rydberg H atoms and excited H (n=2, 3, 4) atoms produced by electron impact on NH<sub>3</sub>. Time-of-flight distribution and Doppler profile analysis of Balmer emission lines were used to derive kinetic energy distributions. Dissociative excitation processes, discussed in detail by the authors, can be linked to dissociative ionization work through the "core-ion" model [8,9].

This paper deals with the dissociative electroionization of  $NH_3$  in the  $H^+$  and  $H_2^+$  dissociation channels, including detailed ion translational energy analysis. The results on  $D^+$  and  $D_2^+$  from  $ND_3$  have been added. The 15-50 eV electron energy range is covered in order to include the  $NH_3$  double ionization energy range. Preliminary results on this particular aspect have already been published earlier [1,10].

## 2. Experimental

The experimental technique, as well as its most recent modifications used in the present work, have extensively been described elsewhere [11,12]. Briefly, the ions produced in a Nier-type ion source by the impact of energy-controlled electrons are allowed to drift out of the ion chamber, are focused on the ion source exit hole, energy analyzed by a retarding lens and mass selected in a quadrupole filter. The detected ion current is continuously scanned as a function of either the electron energy at fixed retarding potential  $V_{\rm R}$  or the retarding potential at fixed electron energy  $E_{\rm e}$ . Both signals are electronically differentiated. The experiment is almost computer controlled.

The experimental conditions are identical with those described previously [1]. Ammonia (Air Liquide, 99.95%) and ammonia- $d_3$  (Merk, Sharp & Dome 99.7%) samples, distilled under vacuum, are introduced at  $10^{-7}$  Torr as measured in the vacuum chamber.

The maximum of the  $NH_3^+$  ( $ND_3^+$ ) ion energy distribution is used as zero-energy calibration point of the translational energy scale. The appearance energy of  $NH_2^+$  ( $ND_2^+$ ), i.e 15.72 ± 0.1 eV is used to calibrate the electron energy scale.

For each fragment ion, at each retarding potential or electron energy setting, the first derivative of the ionization efficiency curve or of the retarding potential curve is scanned 100-800 times. This procedure is repeated at least five times. The energies quoted in the next sections are average values of these five independent measurements. The quoted errors and drawn error bars are given by the standard deviation. In the diagrams, linear regressions are fitted to the experimental data.

## **3. Experimental results**

As an introductory remark, it has to be pointed out that beside  $NH_3$ , two different proton sources could be suspected, (i) Notwithstanding its low intensity after prolongated back out,  $H_2O$  is present in the background mass spectrum and produces  $H^+$  with a large cross section. (ii) As will be shown below  $H_2^+$  is probably formed through ionization of  $H_2$  produced by pyrolysis of  $NH_3$  on the electron emitting filament. Furthermore  $H_2^+$  is observed in the  $H_2O$  mass spectrum, though with very low intensity.

The simultaneous analysis of  $D^+/ND_3$  and  $H^+/NH_3$  would remove, at least to a large extent, the former proton source. On the other hand, the kinetic energy distribution and the appearance energies of  $H^+/H_2$  are available [13] and have to be kept in mind for the interpretation of the experimental results.

The same comments hold for  $H_2^+/NH_3$ . However, most confusing will be the thermal or nearly thermal  $H_2^+(D_2^+)$  ions.

#### *3.1.* The $H^+$ ( $D^+$ ) dissociation channel

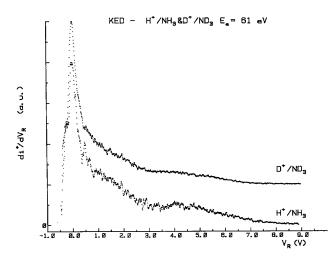
*Typical first differential retarding potential curves of*  $H^+/NH_3$  and  $D^+/ND_3$ , as observed at 61 eV electron energy, are displayed in fig. 1. The identical shape of these curves indicates the absence of any contaminant, at least at low energy. Beside, a dominant low-energy peak, a shoulder at about 0.25 eV, slope changes around 0.6 eV and 1.5 eV, and finally a maximum near 4.0 eV are observed for  $H^+$ . These values slightly shift toward lower energies for  $D^+$ . The translational energy spectrum spreads up to 9 eV. Figs. 2 and 3 show the same distribution for  $D^+$  from 76-31 eV and 26.5-20.1 eV respectively. The maximum of the  $ND_3^+$  thermal ion energy distribution (see fig. 3) and the  $D^+$  low-energy peak only coincide for electron energies close to the lowest onset up to 23 eV. In the electron energy range 23-60 eV the same maximum regularly shifts upward to 45 meV. The same phenomenon is observed for  $H^+$  and is represented by the diagram in fig. 4. This means that the thermal  $H^+$  ions are only formed near the lowest onset. Above 23 eV only energetic protons are produced and dominate the translational energy spectrum.

At this point a comparison has to be made between the present results and those obtained for high Rydberg H(Rydb.) atoms and excited  $H^*(n)$  atoms [6,7]. Concerning the position of the abovementioned features, the present H<sup>+</sup> and D<sup>+</sup> translational energy spectra agree fairly well with the H\* (*n*=2) [6] H\* (*n*=3,4) [7] kinetic energy spectra. However, the relative intensities are quite different, e.g. the distribution peaking at 4.0 eV dominates the H\* kinetic energy spectrum at 70 eV electron energy. Surprisingly even less agreement is found with the translational energy spectrum of H(Rydb.) atoms [6], where a maximum at about 5.5 eV is the most intense feature. This could mean that at least for a number of processes leading to H(Rydb.), the "core-ion" model no longer holds.

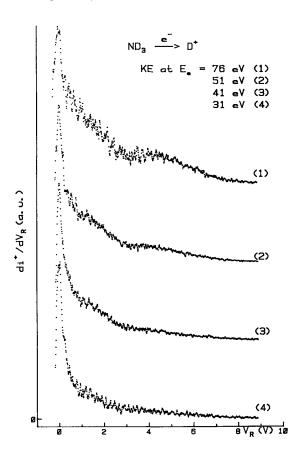
The first differentiated ionization efficiency curves of  $H^+/NH_3$ , as observed for different retarding potential settings  $V_R$ , are shown in fig. 5. Vertical bars locate the average position of the successive appearance energies. The same measurements have been performed on D<sup>+</sup>/ND<sub>3</sub>. For comparison fig. 6 clearly shows both ions having almost the same first differentiated ionization efficiency curve. Six appearance energies have been measured, i.e.  $19.5 \pm 0.2$ ,  $22.6 \pm 0.3$ ,  $26.2 \pm 0.1$ ,  $34.9 \pm 0.4$ ,  $36.2 \pm 0.3$  and  $39.2 \pm 0.2$  eV. A critical energy is also measured at  $45.7 \pm 0.3$  eV and could only be observed for  $0.0 < V_R < 0.9$  V.

The kinetic energy versus appearance energy plot resulting from these measurements is displayed in fig. 7. The same plot for  $D^+/ND_3$ , restricted to the low energy range, is reproduced in fig. 8 together with the results on  $H^+/NH_3$  for comparison. Though both diagrams almost coincide within experimental error, systematic differences are observed which probably should be ascribed to isotope effects.

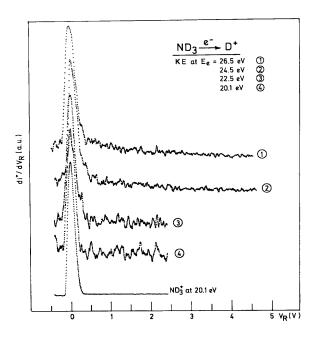
*Fig. 1.* The kinetic energy distribution curves of  $H^+/NH_3$  and  $D^+/ND_3$  observed at 61 eV electron energy.



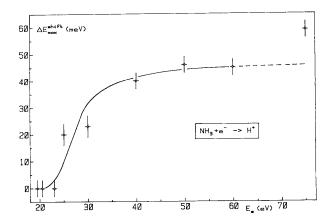
*Fig. 2.* Kinetic energy distribution spectra of  $D^+/ND_3$  recorded between 76 and 31 eV electron energy.



*Fig. 3.* Kinetic energy distribution spectra of  $D^+/ND_3$  recorded between 26.5 and 20.1 eV electron energy. The  $ND_3^+$  distribution is included for comparison.



*Fig. 4.* Electron energy dependence of the position of the maximum of the  $H^+$  low-energy peak.



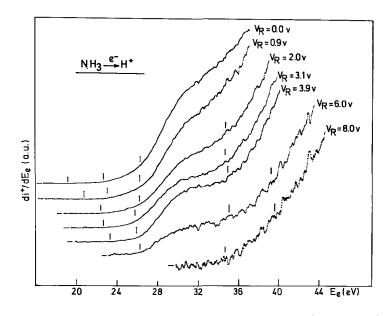
3.2. The  $H_2^+$  ( $D_2^+$ ) dissociation channel

Despite its low intensity (about 0.05% of total ionization), the  $H_2^+$  ion has been investigated because of its potential importance for the interpretation of NH<sup>+</sup> and N<sup>+</sup> formation at high energies.

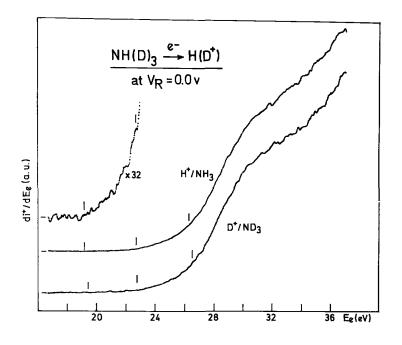
*The*  $D_2^+/ND_3$  *translational energy distribution spectrum,* as obtained by the first derivative of the retarding potential curve, is shown in fig. 9 for different impact energies of the electrons. The  $H_2^+/NH_3$  ion energy distribution is strictly identical. Both distributions exhibit an intense thermal energy peak and a long and structureless tail, extending to about 5 eV, which disappears below 26 eV electron energy.

The first differentiated ionization efficiency curves  $ofD_2^+/ND_3$ , as observed for different values of the retarding potential, are displayed in fig. 10. In fig. 11 the same curve for  $H_2^+/NH_3$  is shown for comparison. For both ions, different appearance energies, marked by a vertical bar, are measured respectively at  $15.5 \pm 0.2$ ,  $22.3 \pm 0.2$ ,  $26.6 \pm 0.3$  and  $35.7 \pm 0.4$  eV. The kinetic energy versus appearance energy diagram for both  $H_2^+/NH_3$  and  $D_2^{+/}ND_3$  is represented in fig. 12.

*Fig. 5.* First differentiated ionization efficiency curves of  $H^+/NH_3$  as observed for indicated retarding potential  $(V_R)$  settings. Vertical bars locate the average appearance energies.



*Fig. 6.* Comparison of the first differentiated ionization efficiency curves of  $H^+/NH_3$  and  $D^+/ND_3$ , both recorded at  $V_R=0.0$  V. Vertical bars locate the average appearance energies at low electron energies.



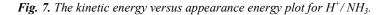
#### 4. Discussion

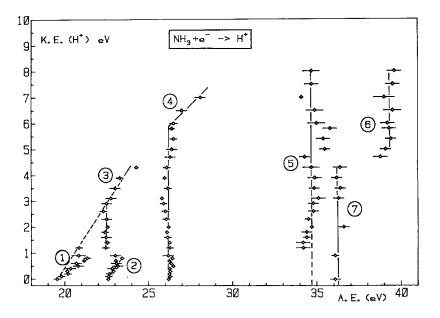
For clarity in the following discussion, the results obtained in the present work are summarized in table 1. All the thermodynamical and spectroscopic data used to calculate critical energies are gathered in table 2.

# 4.1. The $H^+$ ( $D^+$ ) dissociation channel

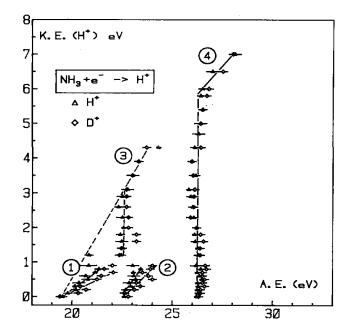
The lowest onset determined for the  $H^+$  production from  $NH_3^+$  is 19.5 ±0.2 eV. This result disagrees

drastically with previously published electron impact results, i.e.  $23.3\pm0.5$  eV [2] and  $23.0\pm0.2$  eV [3]. The discrepancy is at least of 3.5 eV. However, the present result is consistent with the measurements for D<sup>+</sup>/ND<sub>3</sub> for which an appearance energy of  $19.4\pm0.2$  is obtained (see fig. 8).

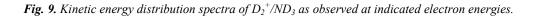


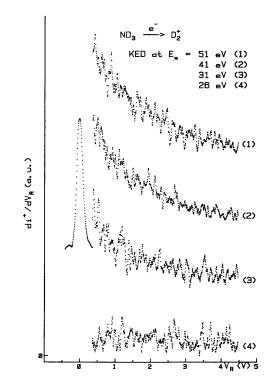


*Fig. 8.* The kinetic energy versus appearance energy plot for  $H^+/NH_3$  (- $\Delta$ -) and  $D^+/ND_3$  (- $\diamond$ -) between 18 and 33 eV electron energy.

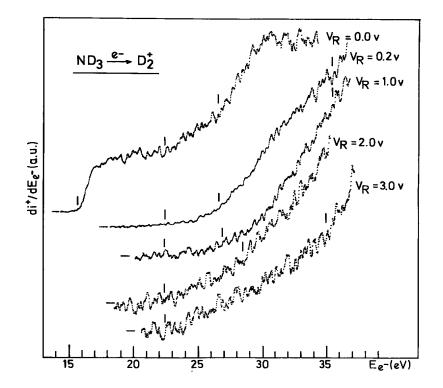


The kinetic energy versus appearance energy diagram related to this threshold is the straight line (1) in fig. 7. In spite of the scattering of the onset measurements for 0-0.9 V retarding potential, a least-squares fit gives a correlation coefficient of 0.96 for a straight line extrapolating to 19.67 eV with a slope of 0.49. For the same measurements in  $D^+/ND_3$ , the straight line extrapolates to 19.62 eV with a slope of 0.29 and a correlation coefficient of 0.95.

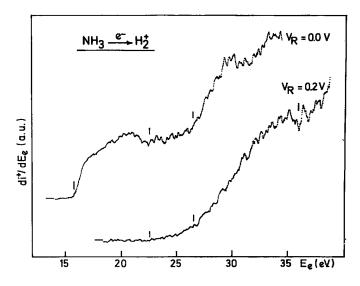




*Fig. 10. First differentiated ionization efficiency curves of*  $D_2^+/ND_3$  *recorded at indicated retarding potential settings*  $V_R$ . *Vertical bars locate the average appearance energies.* 



**Fig. 11**. First differentiated ionization efficiency curves of  $H_2^+/NH_3$  recorded at  $V_R=0.0$  V and 0.2 V shown for comparison with  $D_2^+/ND_3$ .



**Table 1.** The kinetic energy distributions (KED) and appearance energies (AE) (eV) for  $H^+(D^+)$  and  $H_2^+(D_2^+)$  from  $NH_3$  and  $ND_3$  as obtained in this work

	KED	AE
$\mathrm{H}^{+}\left(\mathrm{D}^{+}\right)$	0.0-45 meV	$19.5 \pm 0.2$
	shift	$22.6 \pm 0.3$
	≈0.25	$26.2 \pm 0.1$
	≈0.6	$34.9\pm0.4$
	≈1.5	$36.2 \pm 0.3$
	$\approx 4.0$	$39.2 \pm 0.2$
		$45.7 \pm 0.4$
$H_2^+(D_2^+)$	$(0.0)^{a}$	$(15.5\pm0.2)^{a)}$
		$22.3 \pm 0.2$
		$26.6\pm\!0.3$
		$35.7 \pm 0.4$

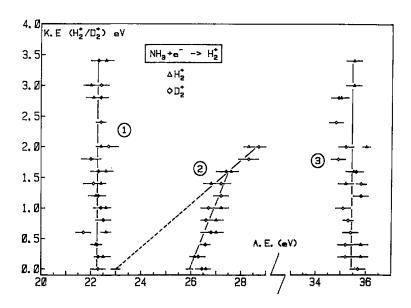
 $^{a)}$  In parentheses data ascribed to  ${\rm H_2^+}\,({\rm D_2^+}\,)$  from H\_2 (D\_2).

Table 2. Disocation energies (D), ionization energies (IE) and excitation energies (EE) (eV) used in this work <sup>a)</sup>

$D(NH_2-H)=4.51 \pm 0.09^{b}$ $D(NH-H) = 3.90 \pm 0.09^{b}$ $D(N-H) = 3.69^{c}$ $D(H-H) = 4.476^{c}$	$EE(NH_{2}^{+}(\widetilde{X}^{3}B_{1})) = 0.0^{d})$ $EE(NH_{2}^{+}(\widetilde{a}^{1}A_{1})) = 0.99^{d})$ $EE(NH_{2}^{+}(\widetilde{b}^{1}B_{1})) = 2.23^{d})$ $EE(NH_{2}(\widetilde{X}^{2}B_{1})) = 0.0^{f}$
$\begin{split} \mathrm{IE}(\mathrm{NH}_2) &= 11.46 \pm 0.01^{\mathrm{d}})\\ \mathrm{IE}(\mathrm{NH}) &= 13.49 \pm 0.01^{\mathrm{d}})\\ \mathrm{IE}(\mathrm{H}_2) &= 15.496^{\mathrm{c}})\\ \mathrm{IE}(\mathrm{H}) &= 13.595^{\mathrm{e}})\\ \mathrm{IE}(\mathrm{N}) &= 14.549^{\mathrm{e}}) \end{split}$	$EE(NH_2(\tilde{A}^{2}A_{1})) = 1.27^{f_{1}}$ $EE(NH(X^{3}\Sigma^{-}))=0.0^{g_{1}}$ $EE(NH(a^{1}\Delta)) = 1.558^{g_{1}}$ $EE(NH(b^{1}\Sigma^{+}))=2.629^{g_{1}}$ $EE(NH(A^{3}\Pi)) = 3.69^{g_{1}}$ $EE(NH(c^{1}\Pi)) = 5.423^{g_{1}}$
EE(N( <sup>4</sup> S))=0.0 <sup>e)</sup> EE(N( <sup>2</sup> D))=2.38 <sup>e)</sup>	$\begin{split} & \text{EE}(\text{NH}^{+}(\text{X}\ ^{2}\Pi)){=}0.0^{\text{h}} \\ & \text{EE}(\text{NH}^{+}(\text{a}\ ^{4}\Sigma)) {=}0.044^{\text{h}} \\ & \text{EE}(\text{NH}^{+}(\text{A}\ ^{2}\Sigma)) {=}2.674^{\text{h}} \\ & \text{EE}(\text{NH}^{+}(\text{B}\ ^{2}\Delta)){=}2.847^{\text{h}} \\ & \text{EE}(\text{NH}^{+}(\text{C}\ ^{2}\Sigma^{+})) {=}4.409^{\text{h}} \end{split}$

<sup>&</sup>lt;sup>a)</sup> 1 eV=23.060 kcal mol<sup>-1</sup> = 8065.73 cm<sup>-1</sup>. <sup>b)</sup> Ref. [14]. <sup>c)</sup> Ref. [15]. <sup>d)</sup> Ref. [16]. <sup>e)</sup> Ref. [17]. <sup>f)</sup> Ref. [18]. <sup>g)</sup> Ref. [19]. <sup>h)</sup> Ref. [20].

*Fig. 12.* The kinetic energy versus appearance energy plot for  $H_2^{+/}NH_3$  (- $\Delta$ -) and  $D_2^{+/}ND_3$  (- $\diamond$ -) between 20 and 36 eV electron energy.



The slope of 0.49 could suggest that the observed  $H^+$  would be produced by dissociative ionization of  $H_2$ . As will be mentioned below,  $H_2$  is in fact produced very likely by thermolysis of  $NH_3$  on the electron emitting filament. In this case the linear dependence of the kinetic energy upon the electron energy would indeed be a straight line with a slope of 0.5. However, the same phenomenon being observed for  $ND_3$ , the same slope of 0.5 has to be expected. Experimentally a slope of 0.29 is observed. Furthermore, the appearance energy of  $H^+/H_2$  is well below 19.5eV, i.e. 18.1 eV [13].

The lowest energy required for the production of H<sup>+</sup> from NH<sub>3</sub> through

$$NH_3+e^- \rightarrow NH_2 (\widetilde{X}^2B_1)+H^+({}^1S_g)+2e^-,$$

is calculated at 18.105 eV, based on the data listed in table 2. As shown in fig. 6, in the first differentiated ionization efficiency curve of  $H^+$  and  $D^+$  no signal could be detected below 19.5 eV. Therefore the energy difference between measured and calculated onset has to be ascribed to the excitation of NH<sub>2</sub> only, i.e. 19.5 — 18.1 = 1.4 eV. This energy difference is close to the excitation energy of NH<sub>2</sub> in its <sup>2</sup>A<sub>1</sub> state found at 1.27 eV by optical spectroscopy (see table 2). The most probable mechanism to be invoked for the threshold at 19.5 eV is

$$NH_3+e^- \rightarrow NH_3^+ (\tilde{A}^2E) \rightarrow NH_2 (\tilde{A}^2A_1)+H^+ + 2e^-,$$

and lies in the energy range of the  $NH_3^+$  ( $\tilde{A}^2E$ ) state. From previous work, this state is known to be split by Jahn-Teller interaction into  ${}^2A' + {}^2A''$  states [21] and to give rise to  $NH^+$  and  $NH_2^+$  by predissociation [21, 1]. The photoelectron band related to the  $\tilde{A}^2E$  state has a shape suggesting that this state is populated well above 16.9 eV [22], the lowest appearance energy of  $NH^+$  [1]. It seems to be likely that at 19.5 eV dissociative ionization is observed through the direct population of the  $\tilde{A}^2E$  ( ${}^2A' + {}^2A''$ ) dissociation continuum. However, at the present time the fact that the lowest dissociation process calculated at 18.105 eV is not observed, remains unexplained.

The expected slope for the dissociation process involving H<sup>+</sup> and NH<sub>2</sub> is  $m_{\text{NH2}}/m_{\text{NH3}} = 16/17=0.94$ , while the observed slope is s=0.49. This indicates that the excess energy with respect to the dissociation limit at 19.5 eV is shared to a large extent (about 50%) between translational and internal (essentially vibrational) energy of NH<sub>2</sub>. This analysis is confirmed by the measurements on D<sup>+</sup>/ND<sub>3</sub>. A slope  $s_i=0.29$  is observed instead of the expected value of  $m_{\text{ND2}}/m_{\text{ND3}} = 18/20=0.9$ .

Both processes involving vibrational energy content of NH<sub>2</sub> and ND<sub>2</sub>, the ratio of the two experimental slopes  $s_i/s$  has to be related to the isotope effect on the vibrational frequencies of both species. This isotope effect is usually measured by  $\rho = (v^i/v)^2$ . For a triatomic molecule of the type XY<sub>2</sub>, the relationships between  $\rho$  and the

molecular masses are given by Herzberg [23]. When the antisymmetric frequency is chosen,  $\rho_3 = (v_3^i/v_3)^2 = 0.56$ , whereas for the symmetric vibrations  $\rho_{I,2}(v_{1,2}^i l/v_{1,2})^2 = 0.28$ . A fairly good agreement is found between  $\rho_3$  and  $s_i/s = 0.59$ , in spite of the poor correlaton coefficient of the least-squares fit. This might be an indication that essentially  $v_3$  is active in this process.

As mentioned above, two investigations of dissociative excitation were performed on H atom formation from NH<sub>3</sub>. Both studies include the kinetic energy distributions and threshold energy measurements. These results can be "extrapolated" to dissociative ionization work. Furthermore, Böse and Sroka [24] examined various H (Lyman) emissions produced by electron impact on NH<sub>3</sub> without kinetic energy measurements.

Carnahan et al. [6] measured the first onset at  $17.8\pm 0.8$  and  $23.2\pm 0.7$  eV for H(*n*=2) and H (Rydb.) atoms respectively. For both processes the minimum translational energy carried by the H atom is about 0.1 eV. Kurawaki and Ogawa [7] observed the first onset at  $21.0\pm 0.8$  and  $22.5\pm 0.8$  eV for H(*n*=3) and H(*n*=4) atoms carrying respectively 1.0 eV and 1.0 eV kinetic energy. None of these energies correlate with a dissociative ionization appearance energy of  $19.5\pm 2$  eV. However, Böse and Sroka [24] measured an onset at  $16.1\pm 0.4$  eV for H (Lym.- $\alpha$ ). When the energy difference of IE(H) — EE(H, *n*=2) = 3.4 eV is added to this value, a threshold of  $19.5\pm 0.4$  eV is obtained for H<sup>+</sup>/NH<sub>3</sub>.

In fig. 7 two onsets are observed around 22.6 eV, i.e.  $22.6 \pm 0.3$  eV, in the kinetic energy range 0-0.8 eV (see straight line (2)), and  $22.4 \pm 0.1$  eV (see vertical line (3)). This latter process is observed over the 2.8-4.3 eV kinetic energy range. Fig. 8 shows the result of the same measurements for D<sup>+</sup>/ND<sub>3</sub>.

At  $22.6 \pm 0.3$  eV enough energy is available to produce the proton through

 $\begin{array}{c} NH_3 + e^- \\ \longrightarrow NH_3^*(autoionization) + e^- \\ & \longmapsto NH(X \, {}^{3}\Sigma^-) + H({}^{2}S_g) + H^+ \end{array}$ 

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+KE(0eV)+e^{-},
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for which the lowest onset is calculated at  $22.01 \pm 0.2$  eV from the data listed in table 2. The higher excitation of NH<sub>2</sub> although possible cannot be considered because of lack of data. If the energy difference of  $0.5 \pm 0.3$  eV in excess between experimental and calculated threshold is significant, it would be ascribed to internal energy of the NH radical, since no translational energy is available at the onset (see diagram (2) in fig. 7). From the following discussion, most likely, the excess energy above the onset is partitioned between vibrational and translational energy.

The straight line (2) starting at 22.6 ± 0.3 eV is the result of a least-squares fit, has a slope of 0.89 and extrapolates to 22.51 eV. The correlation coefficient is rather poor, i.e. 0.93. The expected slope of H<sup>+</sup>, produced through the abovementioned process, is roughly given by  $m_{\rm NH}/2m_{\rm NH3} = 0.44$ . Notwithstanding the poor correlation, the discrepancy between experimental and expected slope is large. Furthermore, when the same fit is applied to the corresponding data of D<sup>+</sup>/ND<sub>3</sub> (see fig. 8, line (2)), a straight line with a slope of 0.54 and extrapolating to 22.47 is obtained. The correlation is also 0.93. The slopes are significantly different: the eight onsets measured for D<sup>+</sup> are systematically higher than the corresponding data for H<sup>+</sup>. Moreover, as evidenced earlier in this discussion, for a diatomic radical  $\rho = (v^i/v)^2 = \mu/\mu^i = 0.94/1.75 = 0.54$ , whereas  $s_i/s \approx 0.5/0.9 \approx 0.55$ . These arguments tend to support a partition of the total excess energy between vibrational and translational energy.

To account for the experimental results a two-step proton production mechanism has to be invoked, i.e.

$$NH(D)_{3} + e^{-} \rightarrow NH(D)_{3}^{*}$$
  
$$\rightarrow NH(D)_{2}^{+} + H(D) + 2e^{-},$$
  
$$NH(D)_{2}^{+} \rightarrow NH(D)(X^{3}\Sigma^{-}) + H(D)^{+},$$

where the proton carries 15/16=0.94 of the total translational energy, close to the observed slope. This agreement would suggest total energy conversion into translational energy. Expecting the same reaction sequence in ND<sub>3</sub>, the theoretical slope for D<sup>+</sup> would be 0.88 while the observed one is 0.54. This discrepancy can only be ascribed

to the excess energy partitioning between translational and vibrational energy. This latter would account for the abovementioned isotope effect.

At about the same energy, i.e.  $22.4 \pm 0.1$  eV, a vertical line followed by a straight line is observed. The least-squares fit of the data points (3) in the diagram of fig. 7 gives a straight line with a slope of 0.93 and extrapolates to 19.31 eV. The correlation coefficient is 0.97. The corresponding data for D<sup>+</sup>/ND<sub>3</sub> (see fig. 8) are in very good agreement. Compared to the expected slope of 0.94, straight line (3) is related to a dissociative ionization process where the excess energy is entirely converted into translational energy of the fragments. The slope of 0.93 and the extrapolated appearance energy AE<sub>KE=0</sub>= 19.3 eV clearly indicate that H<sup>+</sup> is produced by the process

 $\begin{array}{l} NH_3 + e^- \\ \longrightarrow NH_3^*(autoionization) + e^- \\ \longmapsto NH_2(\tilde{A}^2A_1) + H^+ + e^- \\ + KE(2.9 - 4.3 \ eV) \ . \end{array}$ 

The excess energy with respect to the dissociation limit is entirely converted into translational energy carried away by the fragments.

As already mentioned elsewhere [1], at the energy of 22.6 eV several dissociation processes have been observed, leading to  $NH_2^+$  and  $NH^+$ . However, the photoelectron spectrum of  $NH_3$  clearly shows no cross section for ionization of  $NH_3$  at this energy [25]. As suggested earlier, the intermediate state of  $NH_3$  producing  $NH_2^+$ ,  $NH^+$  and  $H^+$  has to be a neutral super-excited state which could have a Rydberg character. In this case this state could be a member of a series converging to the  $NH_3^+$  ( $\tilde{B}^2A_1$ ) state whose adiabatic ionization energy is 24.0 eV [25]. The photoelectron band at this energy is structureless, the  $\tilde{B}^2A_1$  state could be considered as unstable at least in the Franck-Condon region. The  $NH_3^+$  state giving rise to the abovementioned processes would have the same property.

The comparison with previous electron impact work is more feasible in the energy range of 22.5 eV. Mann et al. [2] measured the first onset of  $H^+$  at 22.5 ±0.5 eV and interpreted this threshold by the formation of NH and H. Mark et al. [3] observed the first appearance energy of the proton at 23.0 ± 0.2 eV, without interpretation.

Considering the dissociative excitation experiments involving  $H^*(n)$  and H(Rydb.) atoms, most experiments agree fairly well by measuring an onset near 22.5 eV, i.e.  $22.5\pm0.8$  eV for H(n=4) carrying 1-3 eV kinetic energy [7] and  $23.7\pm0.7$  eV for H (Rydb.) carrying 0.4-4 eV kinetic energy [6]. Bose and Sroka [24] determined a second onset for  $H(Lym.-\alpha)$  at  $21.8 \pm 0.8$  eV. The authors agree to assign to these energies the same process

# $NH_3 + e^- \rightarrow H^* + H + NH + e^-$ ,

but only Kurawaki and Ogawa [7] were able to distinguish two H(n=4) contributions at  $22.5 \pm 0.8$  eV, i.e. that with 1 eV kinetic energy is ascribed to the abovementioned process, whereas a 3 eV kinetic energy component is assigned to

$$NH_3 + e^- \rightarrow NH_2 + H(n=4) + e^-$$
,

in agreement with the present work. Carnahan et al. [6] ascribed the H(2s) and H(Rydb.), at 21.8 and 23.2 eV respectively, to the same process involving, however, 0.4-10 eV kinetic energy.

A vertical line spread over 6 eV translational energy carried by  $H^+$  is observed at  $26.2 \pm 0.1 \text{ eV}$ . A straight line (4) in fig. 7, defined by only three data points, extends to about 7 eV with a slope of 0.55. This value has to be compared with either (i) 0.94 if the proton is produced through the dissociation of the  $NH_2^+$  into NH and  $H^+$  or (ii) 0.44 when  $H^+$  is directly produced by dissociative ionization of NH<sub>3</sub> together with NH and H. When the former process takes place, the discrepancy between observed and calculated slopes could be ascribed to the excess energy partition between translational and internal energy of the polyatomic fragment. In this case an isotope effect ought to be observed. The results on  $D^+/ND_3$  at the same energy (see fig. 8) do, however, not differ significantly. When the latter process is involved, both slope values are in fairly good agreement and no

drastic isotope effect on the slope, i.e. 0.4 instead of 0.44, is expected.

However, taking into account the minimum translational energy of 6 eV, i.e. 13.6 eV in terms of total kinetic energy, and the onset at 26.2 eV, a minimum appearance energy  $AE_{KE=0}=12.6$  eV is calculated. This energy does not fit any dissociative ionization limit of  $NH_3^+$ .

Altogether these arguments strongly suggest that the proton is produced at 26.2 eV by dissociative ionization of  $H_2$ . It will be shown in section 4.2 that free  $H_2$  is produced by thermolysis of  $NH_3$  on the electron emitting filament.

A further reason to suspect the  $H^+/H_2$  production is that Köllman [13] measured an onset at 26.1 ±0.5 eV, in very good agreement with the present value. However, the translational energy range covered by these protons is 4.5-8.0 eV; for 6 eV kinetic energy the onset is about 29 eV.

Despite these evidences, there still remain arguments for the direct  $H^+$  production from NH<sub>3</sub> at 26.2 eV. (i) The isotope effect is not observed because of the restricted number of data points in diagram (4) in fig. 8. (ii) Starting from 6 eV kinetic energy, an appearance energy  $AE_{\kappa E=0}(H^+) = 19.7$  eV, very close to 19.5 eV (the first onset ascribed to  $H^+/NH_3$ ), is obtained, (iii) The energy of 26.2 eV corresponds to the vertical ionization energy of the NH<sub>3</sub><sup>+</sup>( $\tilde{B}^2A_1$ ) state [25] which is known to give rise to fragments with large amounts of kinetic energy [1].

In the high electron energy range, i.e. above 34 eV, four appearance energies are measured:  $34.9 \pm 0.4$ ,  $36.2 \pm 0.3$ ,  $39.2 \pm 0.2$  and  $45.7 \pm 0.4$  eV. The same results are obtained for D<sup>+</sup>/ND<sub>3</sub>. The kinetic energy versus appearance energy diagrams related to these critical energies are shown in fig. 7 by vertical lines (5), (7) and (6) respectively. The last onset being observed for  $0.0 < V_R < 0.9$  V has not been included in this figure.

As mentioned earlier [1], dissociation processes of doubly ionized NH<sub>3</sub> states are mainly involved in this energy range. Ionization energies corresponding to  $NH_3^{2+}$  formation were measured at  $34.9 \pm 0.2$  and  $39.4 \pm 0.5$  eV [10]. By PIPICO a doubly ionized state is observed at  $44.5 \pm 0.5$  eV [26]. When doubly ionized states dissociate into singly ionized fragments, the appearance energy for both species must be equal. Some of these processes have already been discussed in the frame of the  $NH_2^+$  ( $ND_2^+$ ) and  $NH^+$  formation from  $NH_3$  [1].

The  $H^+$  ( $D^+$ ) ions appearing at 34.9  $\pm$  0.4 eV, spread over 8 eV translational energy, are probably related to the  $NH_2^+$  ( $ND_2^+$ ) ions observed at 35.4  $\pm$  0.3 eV. This onset was ascribed to

$$ND_3 + e^- \longrightarrow ND_3^{2+}({}^{1}A_1, (3a_1^{-2})) + 3e^-$$
$$\longmapsto ND_2^{+}({}^{1}A_1) + D^+,$$

for which an onset energy is calculated at 30.5 eV (see data listed in table 2 ). The energy difference of about 35.4 - 30.5 = 5 eV between observed and calculated onsets is ascribed to translational energy carried by both ND<sub>2</sub><sup>+</sup> and D<sup>+</sup>, i.e. 0.5 and 4.5 eV respectively. As for ND<sub>2</sub><sup>+</sup> a translational energy of 0.7 eV is measured [1], D<sup>+</sup> is observed up to 8 eV. Above 5 eV kinetic energy, protons could be produced with other fragments, i.e. H<sup>+</sup>, NH<sup>+</sup> and N<sup>+</sup>. For NH<sup>+</sup> no threshold is observed in the vicinity of 35 eV [1]. Moreover, these protons being of low intensity, H<sup>+</sup> ions produced by the dissociation of H<sub>2</sub><sup>+</sup> (<sup>2</sup>Σ<sub>u</sub><sup>+</sup>) could become important [13]. As mentioned earlier, free H<sub>2</sub>/NH<sub>3</sub> is known to be formed. The large scattering of the data above 5 eV does not allow a least-squares fit.

No comparison can be made with previous electron impact data. Only Mark et al. [3] mention an onset at  $30.6 \pm 0.5$  eV. No critical energy is found in the energy range 26-34 eV in the present work. However, H(*n*=4) and H(Rydb.) atoms have been observed with excitation thresholds at  $33.3 \pm 1.2$  eV [7] and  $33.8 \pm 1.0$  eV [6], respectively. The former fragments carry 4-5 eV kinetic energy whereas the latter are produced with 3.5-8.2 eV kinetic energy. These processes could be related to the protons observed at  $34.9 \pm 0.4$  eV with kinetic energies higher than 5 eV.

By photoionization [26],  $H^+$  and  $NH_2^+$  coincidences are measured at 35.4 ±0.5 eV in good agreement with both  $D^+$  and  $ND_2^+$  onsets observed by electron impact in this work.

The appearance energy at 39.21 0.2 eV is observed for  $5.0 < V_R < 8.0$  V retarding potential. The simultaneous observation of a double-ionization energy of NH<sub>3</sub> at 39.4 ± 0.5 eV would suggest that this state dissociates [10]. Though 8 eV protons could be formed through

no onset is observed at this energy for  $NH_2^+$  nor  $NH_2^+$ [1]. An alternative explanation of this critical energy would be the dissociation of a highly excited, singly ionized  $NH_3^+$  state. In this case the  $NH_3^{2+}$  ionic state would be formed through autoionization which would compete with dissociation. Kurawaki and Ogawa [7] determined an excitation onset for H(n=4) atoms at 38.4 ± 1.5 eV carrying 8-12 eV kinetic energy. Through the "core-ion" model, the corresponding dissociative ionization process would lie at about  $39.3\pm 1.5$  eV accounting for IE(H)-EE(H,n=4) = 0.85 eV. However, no H(Rydb.) atoms have been detected in this energy range by Carnahan et al. [6].

Between 34.9 and 39.2 eV an onset at  $36.2 \pm 0.3$  eV is less clearly observed in the first differentiated ionization efficiency curve of H<sup>+</sup>, but could be observed up to 4.3 eV kinetic energy. As it has already been pointed out [1], appearance energies were measured at  $36.6 \pm 0.3$  and  $36.2 \pm 0.4$  eV for ND<sub>2</sub><sup>+</sup>/ND<sub>3</sub> and NH<sup>+</sup>/NH<sub>3</sub>, respectively. It has been shown that the NH<sup>+</sup> formation occurs through

The corresponding threshold for  $ND_2^+$  has been interpreted by

ND<sub>3</sub>+e<sup>-</sup>→ND<sub>3</sub><sup>2+</sup>(<sup>3</sup>E)  
ND<sub>2</sub><sup>+</sup>(
$$\tilde{X}$$
 <sup>3</sup>B<sub>1</sub>)+D<sup>+</sup>+3e<sup>-</sup>.

The onset for this process is calculated at 29.5 eV, using the data listed in table 2. The amount of total translational energy carried away by the fragments in this process would be 6.7 eV, i.e. 6 eV by  $D^+$  and 0.7 eV by  $ND_2^+$ . Experimentally  $ND_2^+$  is observed up to 0.9 eV whereas the  $D^+$  ion could be measured up to 4.3 eV kinetic energy.

As mentioned in section 3, the appearance energy at  $45.7 \pm 0.4 \text{ eV}$  could only be measured with a reasonable accuracy up to 1.0 eV kinetic energy. At this energy the H<sup>+</sup> contribution very rapidly decreases with increasing retarding potential and the signal/noise ratio becomes worse.

For ND<sub>2</sub><sup>+</sup> as well as NH<sup>+</sup> a threshold energy is measured at 46.5  $\pm$  0.4 and 46.4  $\pm$  0.5 eV; the ions carrying at least 0.7 and 0.9 eV translational energy respectively [1]. Probably these critical energies are related to one another. In addition PIPICO measurements were performed in this energy range and a component, unresolved with respect to m/e= 14 and m/e= 15, has an onset at 44.5±0.5 eV. On the other hand, H(n=4) [7] H(Rydb.) atoms [6] were observed up to 41.7±1.7 eV with translational energies of 8-12 eV.

At this energy, the proton is very likely produced through both

$$NH_3 + e^- \longrightarrow NH_3^{2+}$$

$$\longrightarrow NH^+ (A^2\Sigma^-) + H(^2S_g) + H^+ + 3e^-,$$

$$NH_2^+(?) + H^+ + 3e^-,$$

even when for experimental reasons  $H^+$  (D<sup>+</sup>) could not be observed at  $V_R > 1.0$  V. On the basis of the data listed in table 2 and the observed 7.6 eV total translational energy carried by NH<sup>+</sup> [1], the energy required for the abovementioned NH<sup>+</sup>-producing process is 45.8 eV.

In the  $NH_2^+$  dissociation channel, the measurements were performed on  $ND_3$ . On  $ND_2^+$  only 0.7 eV kinetic energy is measured, i.e. 7 eV total translational energy. The D<sup>+</sup> ion is observed up to 1 eV kinetic energy. The highest dissociation limit for  $ND_2^+$  formation, which could be calculated from the data listed in table 2, lies at 32.7 eV. By adding 7 eV total translational energy of the fragments, an energy of 39.7 eV is obtained. The energy difference of about 7 eV between the observed onset and the foregoing energy balance could only be

ascribed to the excitation of  $ND_2^+$ .

## 4.2. The $H_2^+$ ( $D_2^+$ ) dissociation channel

The lowest appearance energy measured for both  $H_2^+$  and  $D_2^+$  (see figs. 10 and 11) lies at  $15.5 \pm 0.2$  eV. This value is in very good agreement with that published by Mann et al. [2], who observed only one onset at  $15.5 \pm 0.5$  eV. Mark et al. [3] mentioned a threshold at  $14.8 \pm 0.2$  eV.

At this energy the first differentiated ionization efficiency curve rises very steeply and this part of the curve is suppressed by applying 0.2 V retarding potential. The  $H_2^+$  ( $D_2^+$ ) ions produced at 15.5 eV are thermal and are the only contribution to the thermal peak observed in the kinetic energy distribution displayed in fig. 9. Moreover, the ionization energy of molecular  $H_2$  is 15.496 eV (see table 2), in very good agreement with the onset energy for  $H_2^+$  ( $D_2^+$ ) in this work. Together these arguments strongly suggest that free  $H_2$  is produced from NH<sub>3</sub> in the ion source (likely by thermolysis on the electron emitting filament). The same phenomenon is observed for ND<sub>3</sub>.

The first onset ascribed to  $H_2^+$  ( $D_2^+$ ) formation through dissociative ionization of NH<sub>3</sub> (ND<sub>3</sub>) lies at 22.3  $\pm$  0.2 eV. Its dependence upon the retarding potential is given by a vertical line (1) in fig. 12. Within experimental error, the data related to  $H_2^+$  and  $D_2^+$  behave in the same way.

The lowest threshold which could be calculated for  $H_2^+$  production from NH<sub>3</sub> is at 19.45 ± 0.18 eV by using the data listed in table 2. This energy corresponds to the reaction

$$\begin{split} NH_3 + e^- & \longrightarrow NH_3^*(autoionization) \\ & \longmapsto H_2^+ \left( X\,^2\Sigma_g^+ \right) + NH(X\,^3\Sigma^-) + 2e^- \,. \end{split}$$

Experimentally both  $H_2^+$  and  $D_2^+$  carry away 3.4 eV kinetic energy. The appearance energy  $AE(H_2^+)_{KE=0}=22.3-3.4X17/15 = 18.5$  eV. This value is lower than that for the predicted dissociation mechanism mentioned above. The discrepancy between calculated onset and latter value is about 1.0 ±0.4 eV. In the discussion of the results on the  $NH_2^+$  and  $NH^+$  ion formation [1], discrepancies of 0.24 ±0.13 and 0.5 ± 0.3 eV have already been mentioned. They were ascribed to an overestimate of one or both dissociation energies *D* (NH<sub>2</sub>-H) and (NH-H).

Considering all the fragment ions studied so far, appearance energies are measured at about 22.6 eV for  $NH_2^+$  [1],  $NH^+$  [1],  $H^+$  (see table 1) and  $H_2^+$ . This critical energy lies where no ionization cross section is measured for  $NH_3$  [25]. All abovementioned ions are formed with large amounts of kinetic energy. The neutral  $NH_3^*$  state is very likely a Rydberg state, member of a series converging to the  $NH_3^+$  ( $\tilde{B}^2A_1$ ) state. This latter being almost dissociative, the  $NH_3^*$  state is very probably unstable in all the dissociation channels. Therefore dissociative autoionization will be the most probable mechanism of  $H_2^+$  formation at 22.3 eV.

From 26.6  $\pm$ 0.3 eV the first derivative of the ionization efficiency curve of  $H_2^+$  ( $D_2^+$ ) rises more steeply. The kinetic energy versus appearance energy plot is shown in fig. 12 by diagram (2). Owing to the scattering of the measurements, the data related to  $H_2^+$  and  $D_2^+$  are not distinguishable. From 0.2 eV translational energy, the data more or less clearly deviate from a vertical line at about 26.6 eV.

Taking into account all the measurements related to both  $H_2^+$  and  $D_2^+$ , the best correlation is found when gathering the data into two groups. Two fits are found by this way: (i) a first straight line between 0.2 and 1.4 eV kinetic energy, with a slope of 0.9 extrapolating to 26 eV with a correlation of 0.8 and (ii) a second straight line spread over 1.4-2.0 eV with a slope of 0.35 extrapolating to 23.0 eV with a correlation coefficient of 0.95.

The slope expected for the kinetic energy versus appearance energy straight line depends on the dissociation mechanism involved,

$$NH_{3}+e^{-} \longrightarrow NH+H_{2}^{+}+2e^{-}, \qquad (a)$$

$$\longrightarrow N+H+H_{2}^{+}+2e^{-}, \qquad (b)$$

$$\longrightarrow NH_{2}^{+}+H \longrightarrow N+H_{2}^{+}+2e^{-}, \qquad (c)$$

i.e. roughly 0.88, 0.27 and 0.87 for reactions (a), (b) and (c) respectively.

The slowest rising straight line has to be related to reaction (b): (i) the experimental slope of 0.35 compares favourably with the expected value of 0.27 and (ii) the extrapolated threshold of 23 eV is close to the calculated onset of 23.1 eV. The steeply rising straight line could be ascribed to either reaction (a) or (c) or both. Using the data of table 2, the corresponding appearance energies are 19.51 and 23.12 eV. Introducing the excitation energy of NH, the highest dissociation limit calculated for reaction (a) is 24.9 eV. For the two-step reaction (c), when both the parent ion NH<sub>2</sub><sup>+</sup> and the product N contain excitation energy, i.e. are formed in the  $\tilde{a}$  <sup>1</sup>A<sub>1</sub> and <sup>2</sup>D state respectively (see table 2), an onset of 26.46 eV can be calculated. This value is comparable to the extrapolation at 26 eV.

These processes have to occur through the  $NH_3^+$  ( $\tilde{B}^2A_1$ ) state, whose vertical ionization energy lies at about 26 eV and which very probably is unstable in the Franck-Condon region. The dissociation of  $NH_3^+$  at the expense of the same state has been observed for  $NH^+$  and  $H^+$  but not for  $NH_2^+$ . The latter ion would be formed, but spontaneously dissociates through reaction (c).

In the high electron energy range only one onset is measured for  $H_2^+$  ( $D_2^+$ ) at 35.7 ± 0.4 eV. This threshold is observed up to 3.4 eV translational energy (see vertical line (3) in fig. 12). In this energy range double ionization of NH<sub>3</sub> occurs [10] as well as dissociation of doubly ionized states [1,26]. In this case the two singly ionized species will have the same appearance energy and will share the same total translational energy.

Energetically the lowest dissociation process leading to  $H_2^+$  from a doubly ionized state is

$$\begin{split} NH_3 + e^- &\to NH_3^{2+} \\ &\to H_2^+ (X\,^2\Sigma_g^+) + NH^+ (X\,^2\Pi) + 3e^- \,, \end{split}$$

for which an onset is calculated at 32.8 eV by using the data listed in table 2.

For the NH<sup>+</sup> ion, the threshold measured at  $36.2 \pm 0.4 \text{ eV}$  [1] agrees within experimental error with the onset measured for H<sub>2</sub><sup>+</sup>. Furthermore, on NH<sup>+</sup> a kinetic energy of 0.4 eV is observed, i.e. 0.4x17/2=3.4 eV total translational energy. For H<sub>2</sub><sup>+</sup> a total kinetic energy of 3.4x17/15 = 3.85 eV is estimated. From both measurements the energy level of the dissociation limit is given by the energy difference 35.7(36.2)-3.85(3.4) = 31.85(32.8) eV.

Taking into account the accuracy on the data, the fairly good correlation between the measurements on NH<sup>+</sup> and H<sub>2</sub><sup>+</sup> as well as the closeness of calculated and experimentally derived dissociation limits allow us to ascribe the abovementioned dissociation mechanism to the appearance energy of H<sub>2</sub><sup>+</sup> at 35.7 eV. However, no conclusion can be drawn whether the NH<sub>3</sub><sup>2+</sup> state involved is a <sup>1</sup>A<sub>1</sub> or a <sup>3</sup>E state.

#### 5. Conclusion

The dissociative ionization of  $NH_3$  in both the  $H^+$  and  $H_2^+$  dissociation channels has been investigated by recording the translational energy distribution of both species at different electron energies and by measuring their appearance energies at increasing retarding potential. The simultaneous examination of  $D^+$  and  $D_2^+$ formation from ND<sub>3</sub> has proved to be helpful in the discussion of the experimental results.

At the lowest onset, only  $H^+$  is directly produced by dissociative ionization of NH<sub>3</sub>, very likely through the NH<sub>3</sub><sup>+</sup> ( $\tilde{A}^2E$ ) state. Vibronic energy of NH<sub>2</sub> has to be involved in this process. This is confirmed by an isotope effect due to ND<sub>2</sub> on the slope of the kinetic ' energy versus appearance energy diagram. The H<sub>2</sub><sup>+</sup> ion is produced by direct ionization of H<sub>2</sub> formed by thermolysis of NH<sub>3</sub> on the electron emitting cathode.

In the energy range 20-26 eV, both species are produced by dissociative autoionization at about 22.5 eV and direct dissociative ionization at 26 eV. However, some evidences are brought to show that a part of the  $H^+$ 

ion current, formed at 22.6 eV, is produced by the spontaneous dissociation of  $NH_2^+$ . The same mechanism is invoked for a part of the  $H_2^+$  ion current observed at 26.2 eV. The  $H^+$  ions appearing at this latter energy could be produced by dissociative ionization of  $H_2$ .

In the high electron energy range of 35-50 eV, mainly direct dissociation of doubly ionized states of  $NH_3$  is involved. Most of the processes invoked to account for the experimental data have already been discussed for the appearance of  $NH_2^+$  and  $NH^+$  from  $NH_3$  in a previous work [1]. Though a double ionization energy of  $NH_3$  was measured at 39.5 eV and an  $H^+$  formation threshold at about the same energy, very likely a dissociation of a highly excited, singly ionized  $NH_3$  state is involved.

#### Acknowledgement

We acknowledge the Fonds de la Recherche Fondamentale Collective (FRFC), the Université de Liège and the Action de Recherches Concertées (ARC) of the Belgian Government for financial support. One of us (ML) wishes to thank the Belgian Government for a PREST grant.

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Published in: Chemical Physics (1988), vol. 125, pp. 425-437. Status: Postprint (Author's Version)

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