

Simulation of the highly non linear properties of bulk superconductors: finite element approach with a backward Euler method and a single time step

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Breaking the gravity....

...with the magnetic levitation



Electrical resistivity of HTS is modelled by a non linear power law

$$E(J) = \rho(J)J = E_c \left(\frac{J}{J_c} \right)^n$$

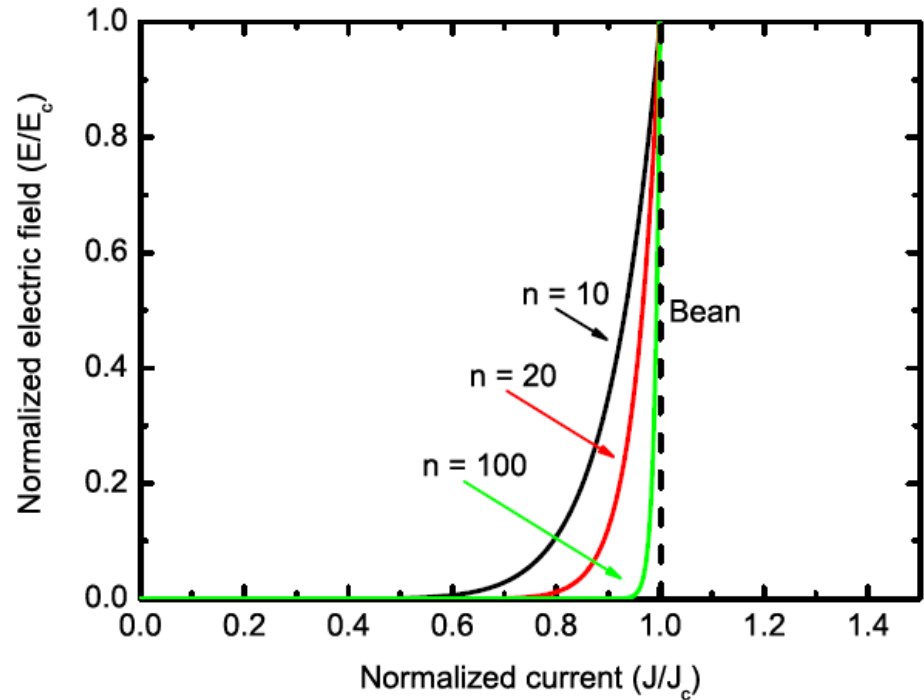
n is the critical exponent
(usually large !)

J_c is the critical current density

Asymptotic behavior

- $n \rightarrow \infty$: **Bean model**

Analytical calculations available for comparison in specific geometries



Finite-element softwares are widely used for simulating HTS-based systems

Advantages

- **Many more geometries** can be treated
- **No extensive writing of numerical codes** is required
- **Treatment of non-linear problems** available in most commercial packages

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Drawbacks

- **Long calculation time** on fine meshing or in 3D geometry
- **Convergence problems** when n is **large**

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Proposed improvements

Single time step method implemented in an open-source solver, GetDP

- **Better control** of the algorithm parameters
- Used for simulating the **penetration of an external magnetic field** that varies linearly with time

Outline

- Finite-element formulation and implementation

A - ϕ formulation

Numerical resolution scheme

- Validation and comparison of the FEM model on simple geometries

A - ϕ formulation - Variables

The Maxwell equations are solved for two independent variables

- the **vector potential \mathbf{A}**
- the **scalar potential ϕ**

defined as

$$\mathbf{B} = \mathbf{B}_{\text{react}} + \mathbf{B}_a = \text{curl } \mathbf{A} + \text{curl } \mathbf{A}_a$$

$$\mathbf{E} = - d\mathbf{A}/dt - d\mathbf{A}_a/dt - \text{grad } \phi$$

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Applied magnetic flux density
(uniform)

A - ϕ formulation - Variables

The Maxwell equations are solved for two independent variables

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The diagram illustrates the definition of the magnetic flux density \mathbf{B} in the A- ϕ formulation. It features a central light-brown rounded rectangle containing two equations in red text. The first equation is $\mathbf{B} = \mathbf{B}_{\text{react}} + \mathbf{B}_a = \text{curl } \mathbf{A} + \text{curl } \mathbf{A}_a$, where $\mathbf{B}_{\text{react}}$ and \mathbf{B}_a are highlighted with grey circular backgrounds. The second equation is $\mathbf{E} = -d\mathbf{A}/dt - d\mathbf{A}_a/dt - \text{grad } \phi$. A blue line originates from the $\mathbf{B}_{\text{react}}$ term, extends upwards, then rightwards, and finally downwards to point at the text "Induced magnetic flux density by the HTS". Another blue line originates from the \mathbf{B}_a term, extends downwards, then rightwards, and finally downwards to point at the text "Applied magnetic flux density (uniform)".

$$\mathbf{B} = \mathbf{B}_{\text{react}} + \mathbf{B}_a = \text{curl } \mathbf{A} + \text{curl } \mathbf{A}_a$$
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Induced magnetic flux density by the HTS

Applied magnetic flux density (uniform)

A - ϕ formulation - Variables

The Maxwell equations are solved for two independent variables

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- the **scalar potential ϕ**

approximated by

Edge function (1st order)

$$\mathbf{A} = \sum_i a_i \mathbf{A}_i$$

where \mathbf{A}_i and ϕ_j are known functions

Node function (1st order)

$$\phi = \sum_j b_j \phi_j$$

- ensures continuity of the tangential component of \mathbf{A}
- ensures continuity of ϕ

A - ϕ formulation – Equations and gauge condition

- The Maxwell equations are reduced to 2 equations

$$\left\{ \begin{array}{l} \nabla \times \nabla \times \mathbf{A} = \mu_0 \sigma(\mathbf{A}, \phi) (-\dot{\mathbf{A}} - \dot{\mathbf{A}}_a - \nabla \phi) \\ \nabla \cdot \{ \sigma(\mathbf{A}, \phi) (-\dot{\mathbf{A}} - \dot{\mathbf{A}}_a - \nabla \phi) \} = 0 \end{array} \right.$$

Ampere's law (rot H = J)

Continuity equations (div J = 0)

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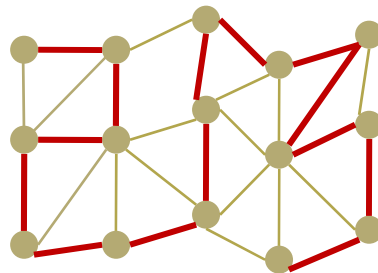
Continuity equations (div J = 0)

- A is not expressed in the Coulomb gauge, gauge condition : $\mathbf{A} \cdot \mathbf{w} = 0$

Set of meshing edges that **connects** all the nodes **without closed contours**



Examples



A - ϕ formulation – Boundary conditions

- The Maxwell equations are reduced to 2 equations

$$\left[\begin{array}{l} \nabla \times \nabla \times \mathbf{A} = \mu_0 \sigma(\mathbf{A}, \phi) (-\dot{\mathbf{A}} - \dot{\mathbf{A}}_a - \nabla \phi) \\ \nabla \cdot \{ \sigma(\mathbf{A}, \phi) (-\dot{\mathbf{A}} - \dot{\mathbf{A}}_a - \nabla \phi) \} = 0 \end{array} \right. \quad \begin{array}{l} \text{Ampere's law (rot } H = J) \\ \text{Continuity equations (div } J = 0) \end{array}$$

- A is not expressed in the Coulomb gauge, gauge condition : $\mathbf{A} \cdot \mathbf{w} = 0$
- Boundary conditions

$$\left. \begin{array}{l} \mathbf{A} = 0 \\ \phi = 0 \end{array} \right\} \text{ at infinity}$$

→ Use of **Jacobian transformation**
for sending the outer surface of a
spherical shell to infinity

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Ampere's law (rot $\mathbf{H} = \mathbf{J}$)

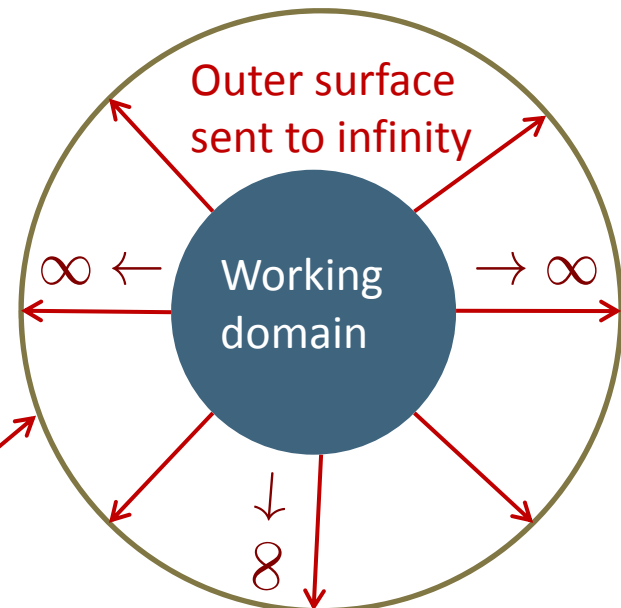
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→ Use of **Jacobian transformation** for sending the outer surface of a spherical shell to infinity

Dirichlet conditions



A - ϕ formulation – External field

- The Maxwell equations are reduced to 2 equations

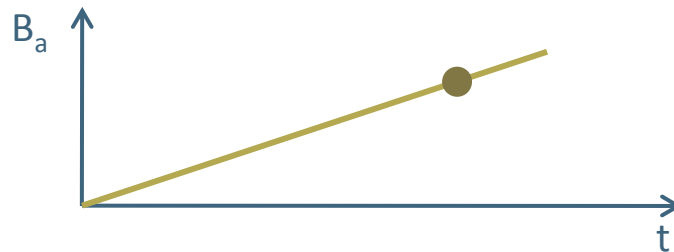
$$\left\{ \begin{array}{l} \nabla \times \nabla \times \mathbf{A} = \mu_0 \sigma(\mathbf{A}, \phi) (-\dot{\mathbf{A}} - \dot{\mathbf{A}}_a - \nabla \phi) \\ \nabla \cdot \{ \sigma(\mathbf{A}, \phi) (-\dot{\mathbf{A}} - \dot{\mathbf{A}}_a - \nabla \phi) \} = 0 \end{array} \right. \quad \begin{array}{l} \text{Ampere's law (rot } H = J) \\ \text{Continuity equations (div } J = 0) \end{array}$$

- A is not expressed in the Coulomb gauge, gauge condition : $\mathbf{A} \cdot \mathbf{w} = 0$

- Boundary conditions $\left. \begin{array}{l} \mathbf{A} = 0 \\ \phi = 0 \end{array} \right\}$ at infinity

Source field \mathbf{A}_a corresponds to a **uniform magnetic flux density** \mathbf{B}_a

The source field is a temporal **ramp** with a **constant sweep rate** (mT/s)



Implicit time-resolution and non-linear Picard iteration

- The equations are solved with a **Galerkin residual minimization method**
- We use the **Backward Euler method** at each time step
- Non linear terms are treated with a **Picard iteration** loop

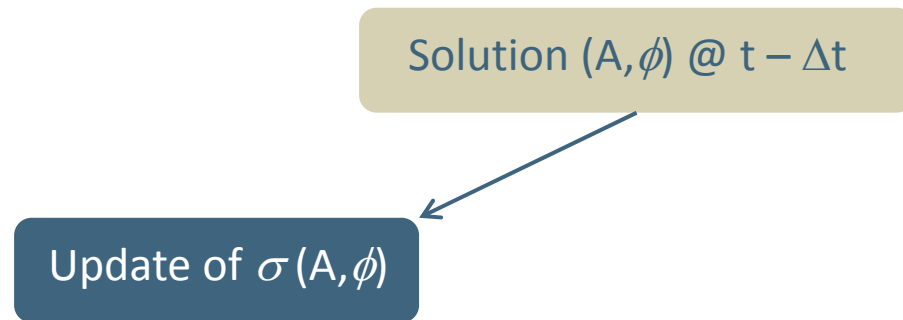
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Solution (A, ϕ) @ $t - \Delta t$

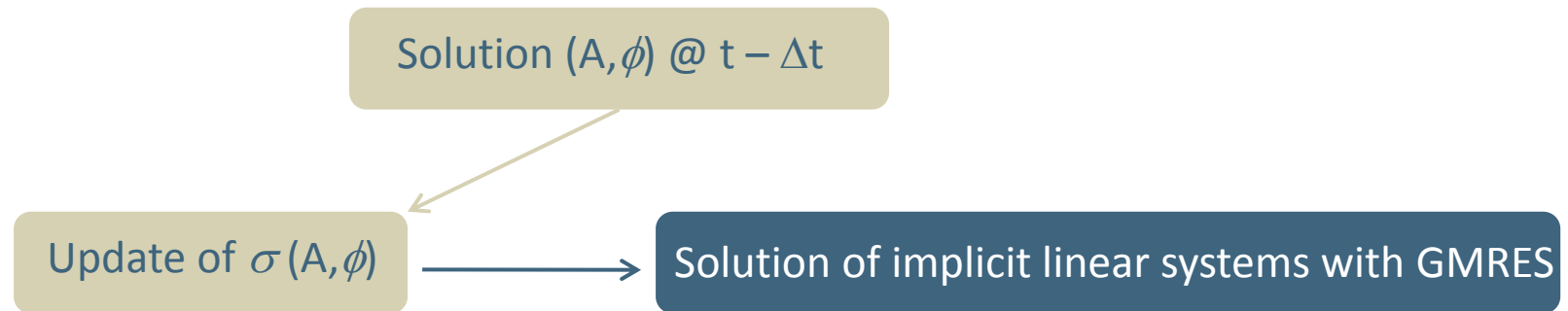
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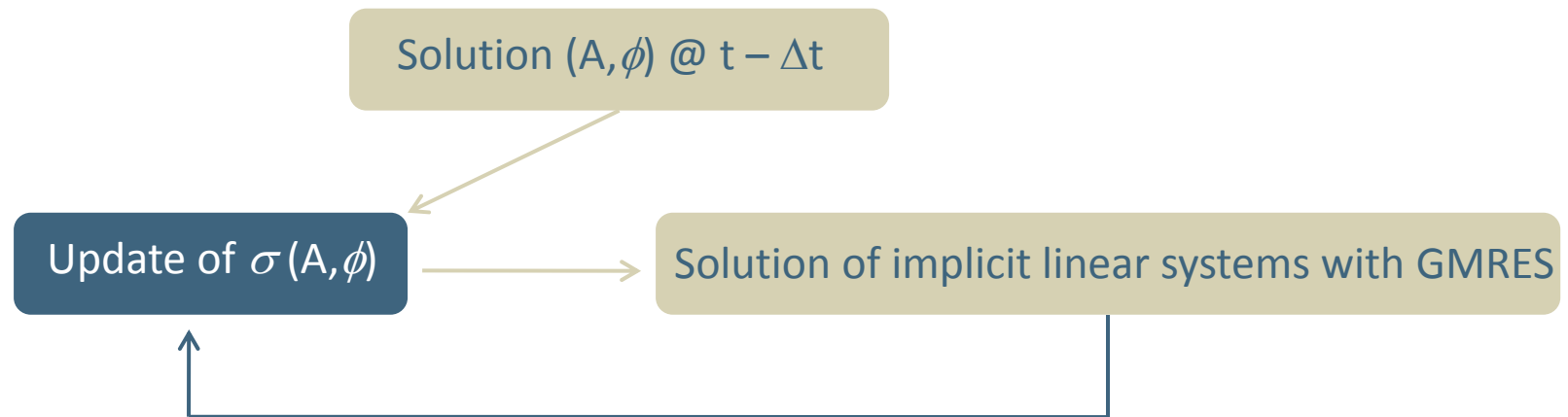
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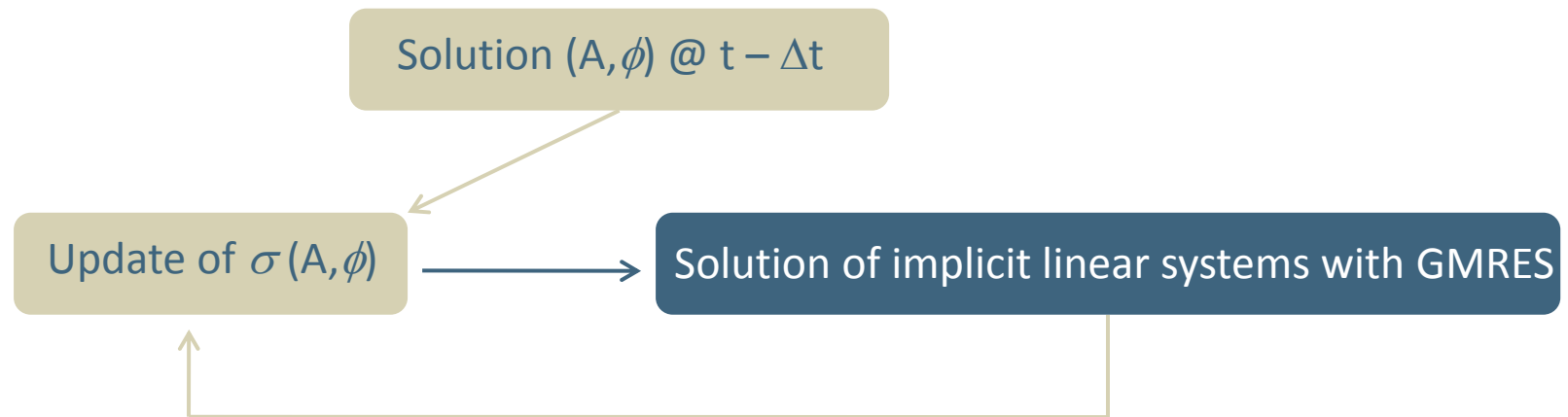
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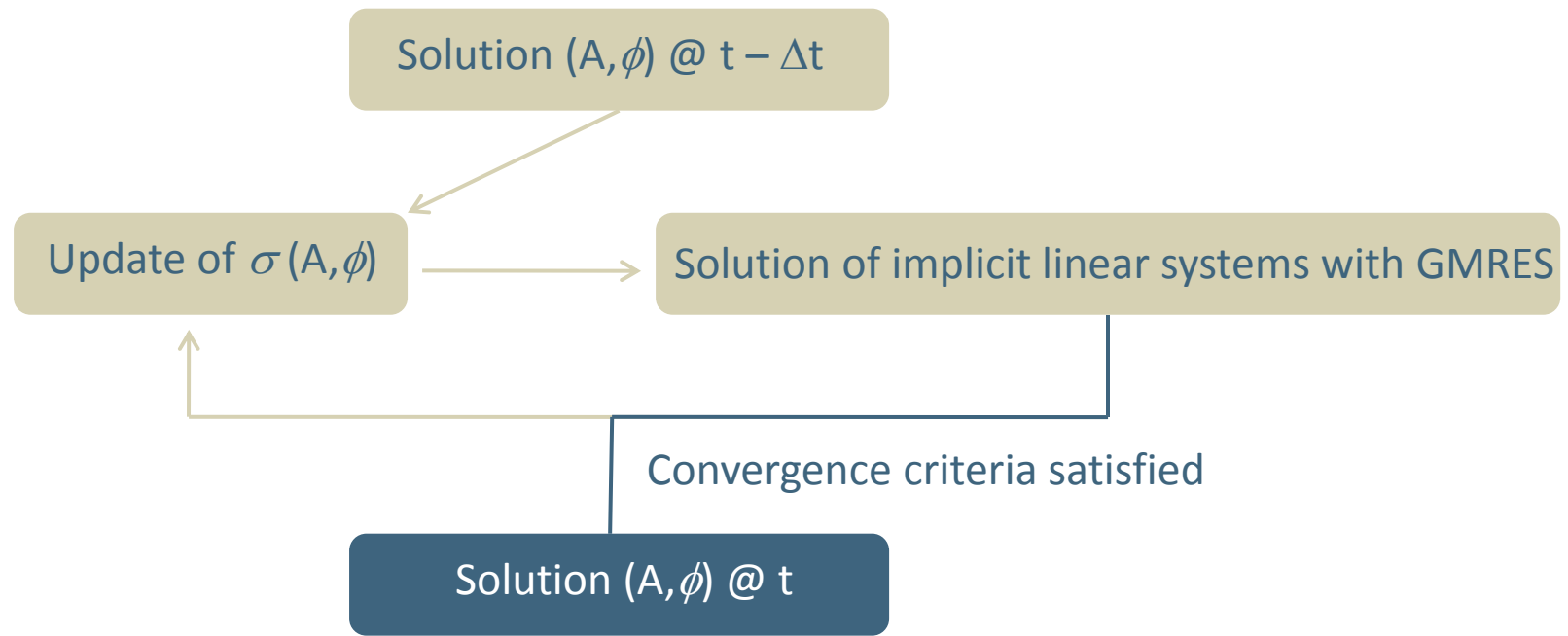
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Outline

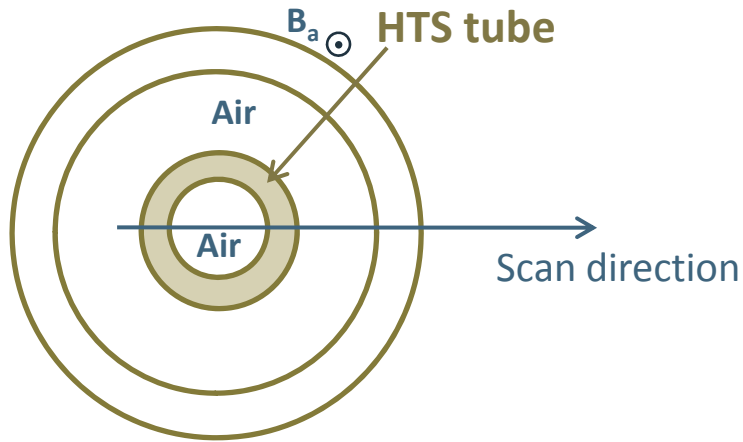
- Finite-element formulation and implementation
- Validation and comparison of the FEM model on simple geometries

Magnetic field penetration in :

- a HTS tube of infinite extension
- a HTS tube of finite height

Bean critical-state is solved with the E(J) model taking $n = 100$

HTS tube of infinite height

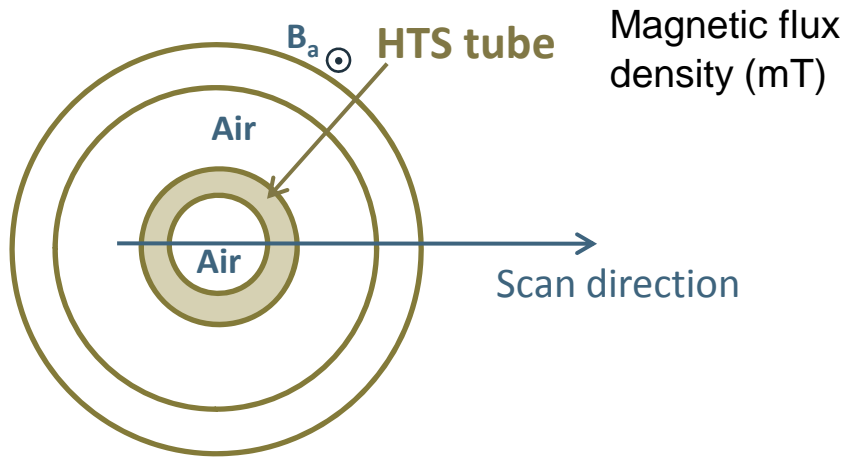


Bean model

linear decay of the magnetic flux density in the wall

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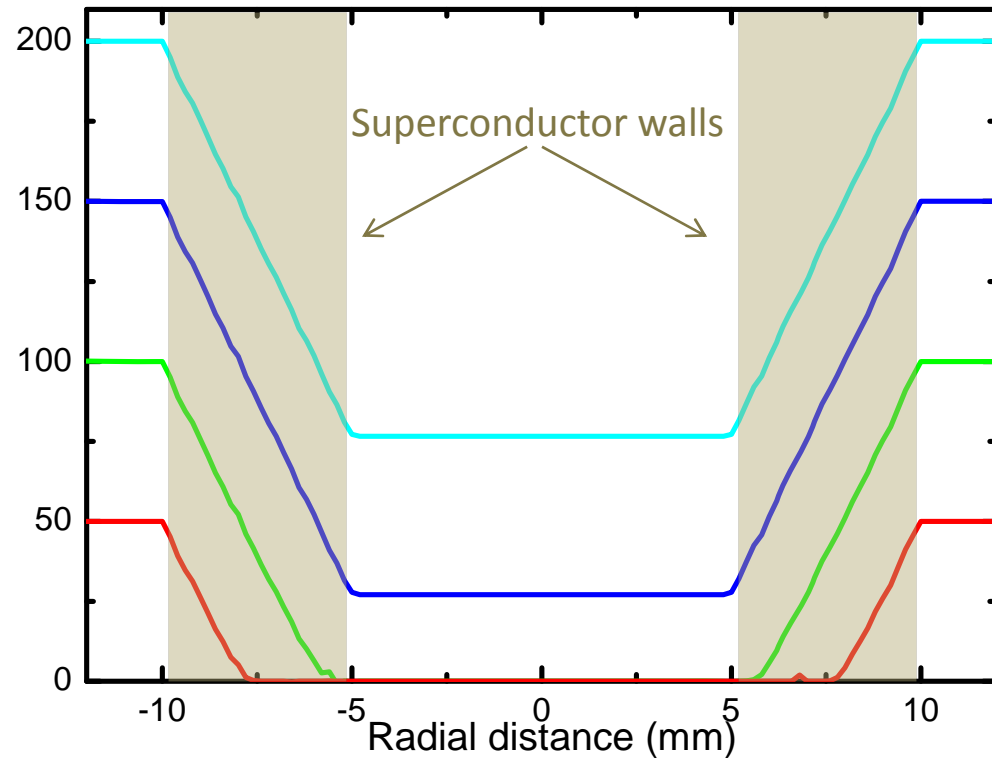
HTS tube of infinite height



Bean model

linear decay of the magnetic flux density in the wall

Magnetic flux penetration



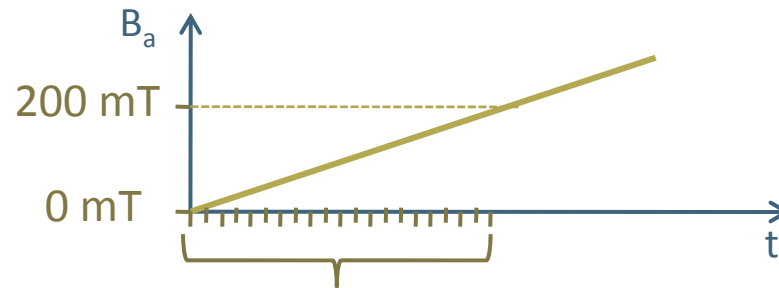
➡ FEM results are consistent with the Bean model

Single time step method produces more accurate results in a smaller calculation time

Choice of the time step

For solving the problem with $B_a = 200 \text{ mT}$ with a sweep rate of 10 mT/s :

1. 20 time steps of 1s



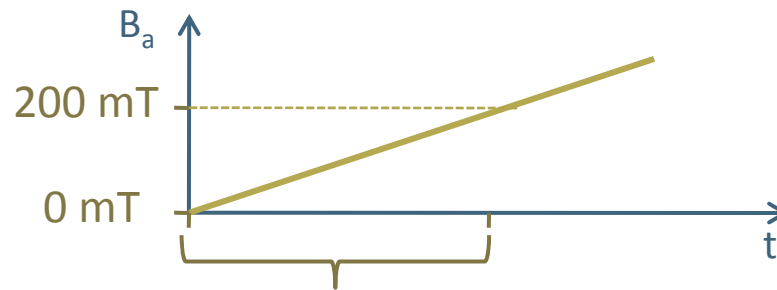
System is solved @ $t = 1\text{s}, 2\text{s}, \dots, 20\text{s}$

Single time step method produces more accurate results in a smaller calculation time

Choice of the time step

For solving the problem with $B_a = 200$ mT with a sweep rate of 10 mT/s:

1. 20 time steps of 1s
2. 1 time step of 20 s → Single time step method



System is solved @ $t = 20$ s

Single time step method produces more accurate results in a smaller calculation time

Magnetic field penetration of an external field increasing from 0 mT to 200 mT (10mT/s)

Comparison of two methods

1. **in a single simulation** with 20 time steps of 1s
2. **in 20 different simulations** with time steps of 1s, 2s, ..., 20s

Single time step method produces more accurate results in a smaller calculation time

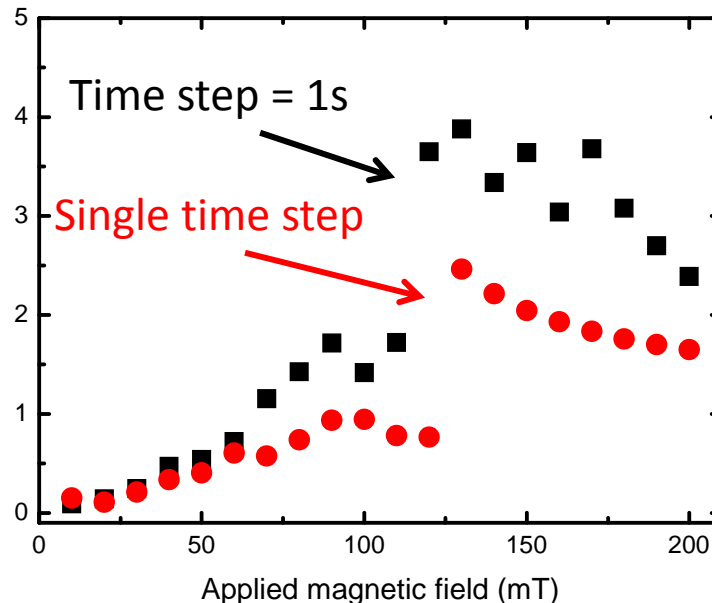
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Analysis of the error

Average deviation
from the Bean model
(mT)



Single time step method produces more accurate results in a smaller calculation time

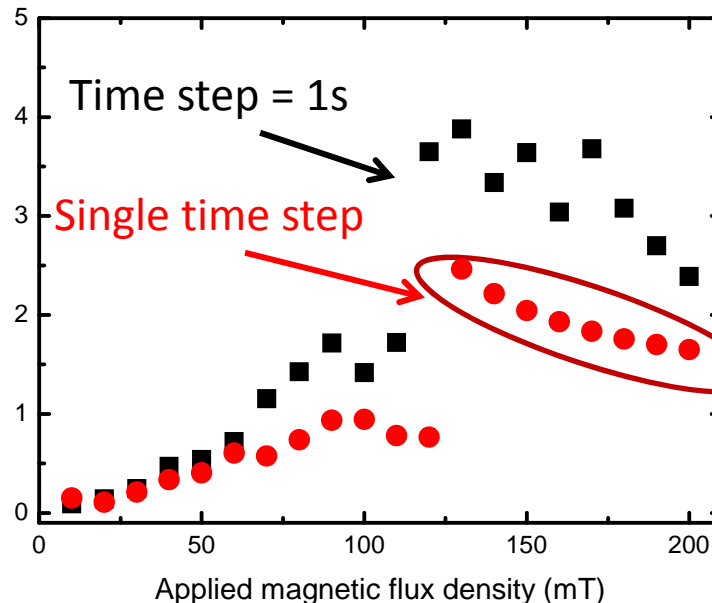
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Calculation time

for solving $B_a = 200$ mT

20 time steps

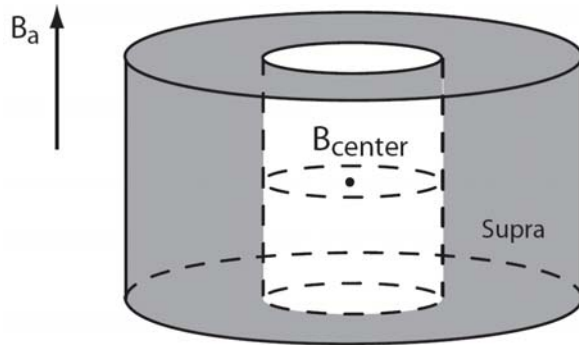
2 days 3/4

Single time-step

20 min

3D FEM simulations on tubes with finite height are consistent with previous observations

HTS tube of finite height



Simulation parameters

$$n = 100$$

$$B_a = 0 - 200 \text{ mT (10 mT/s)}$$

Single time step method

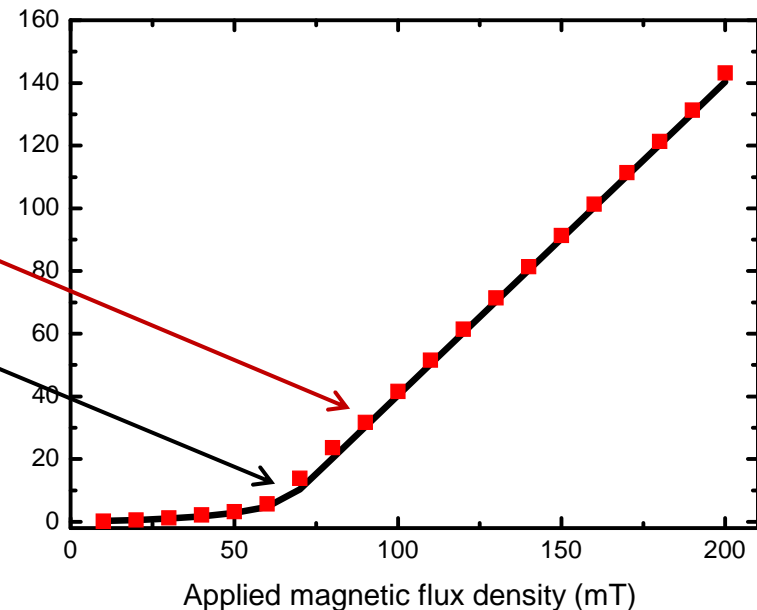
3D geometry

20 FEM simulations
with single time step method

Brandt method (*semi-analytical*)
multiple step method ($\Delta t = 5 \cdot 10^{-4} \text{ s}$)

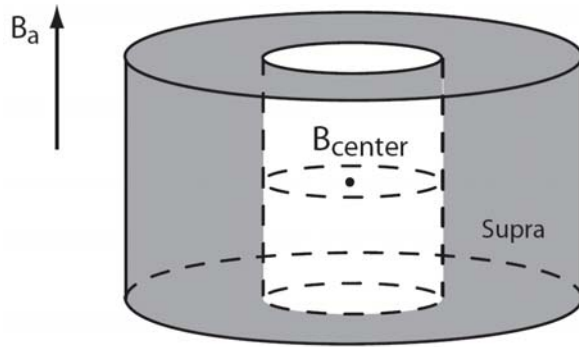
Magnetic flux penetration

Magnetic flux density
in the center of the tube (mT)



Single time step method is more accurate with large critical exponent

HTS tube of finite height



Principle

- Applied magnetic flux density : 200 mT
- FEM single time step compared with Brandt method with the help of B_{center}

$$\text{Error} = \frac{B_{center, \text{Brandt}} - B_{center, \text{FEM}}}{B_{center, \text{Brandt}}}$$

Brandt (*semi-analytical*)

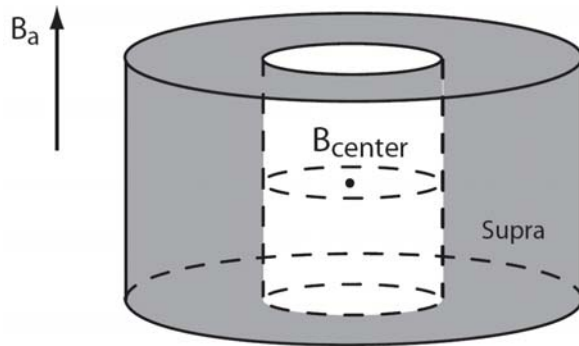
multiple time step method ($\Delta t = 5.10^{-4} \text{s}$)

FEM

single time step method

Single time step method is more accurate with large critical exponent

HTS tube of finite height



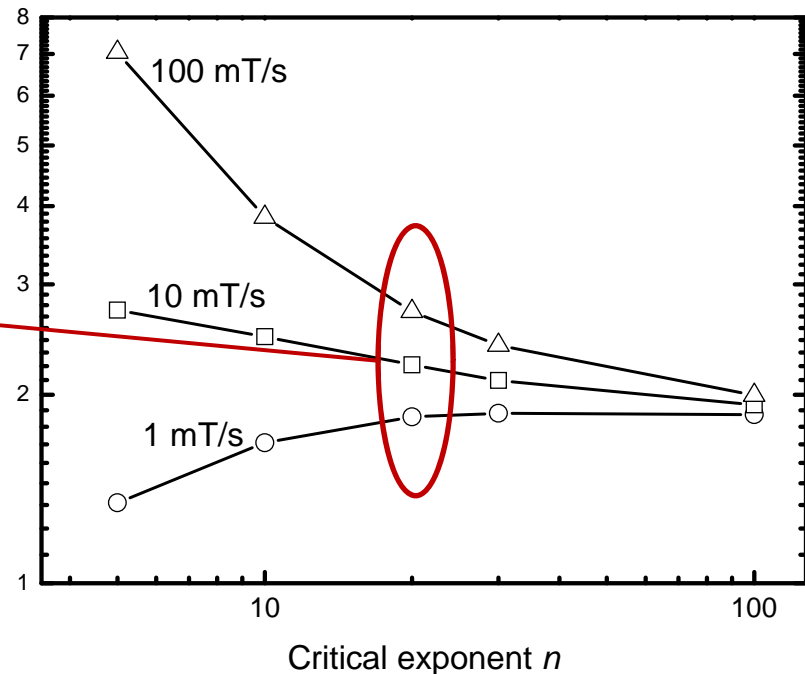
$n \sim 20$: experimental value for the critical exponent

Single time step method is valid

Principle

- Applied magnetic flux density : 200 mT
- FEM single time step compared with Brandt method with the help of B_{center}

Difference between FEM-single step and Brandt for the magnetic flux density @ center of the tube (%)



Conclusion

- Implementation of a finite-element formulation in GetDP with high non linearity
- $A-\phi$ formulation in 3D geometry for calculating the magnetic field penetration
- Single time step method in the case of linearly time varying excitation is fast and accurate

Outline

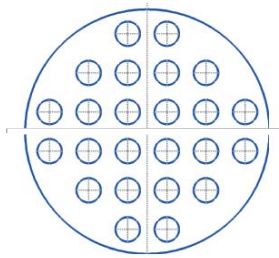
- Finite-element formulation and implementation
- Validation and comparison of the FEM model on simple geometries
- Optimization of the magnetic properties of drilled samples

Influence of the lattice types

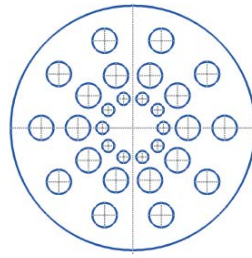
Polar triangular lattice wins ...

HTS tube of infinite height

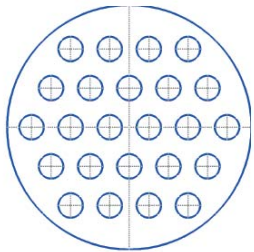
Squared



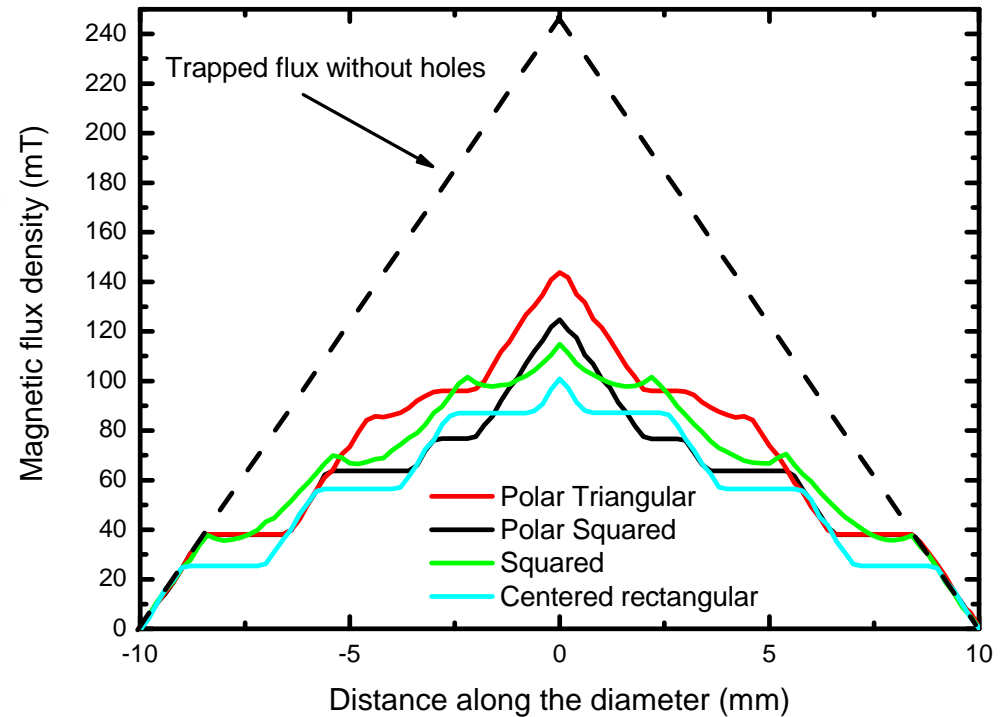
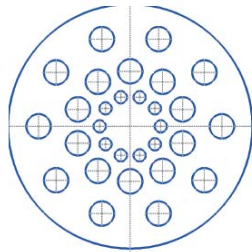
Polar squared



Centered rectangular



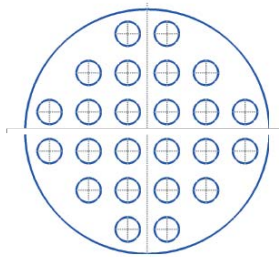
Polar triangular



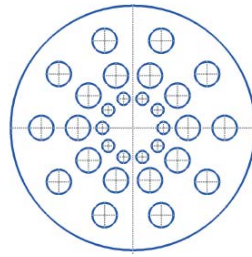
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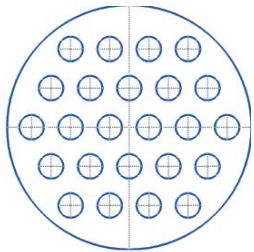
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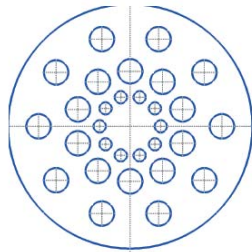
Polar squared



Centered rectangular



Polar triangular



Difference from the
squared lattice (%)

