

Master equation for collective spontaneous emission with quantized atomic motion

Phys. Rev. A **93**, 022124 (2016)

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DPG Spring Meeting - 2 March 2016

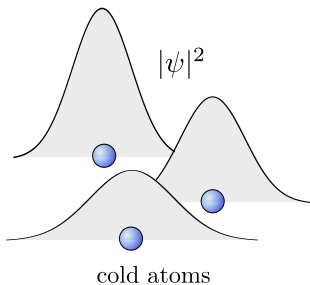


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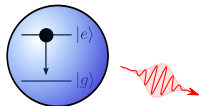
Motivation

Quantization
of the atomic motion

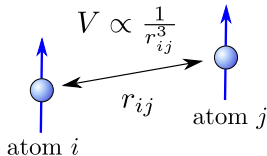


Internal atomic dynamics

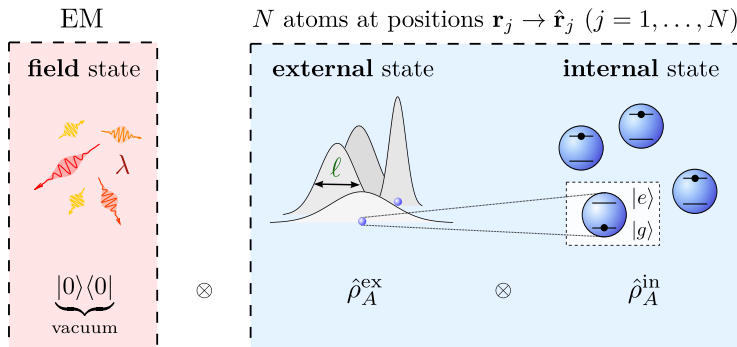
- photon emission



- virtual photon exchange



System



Lamb-Dicke parameter

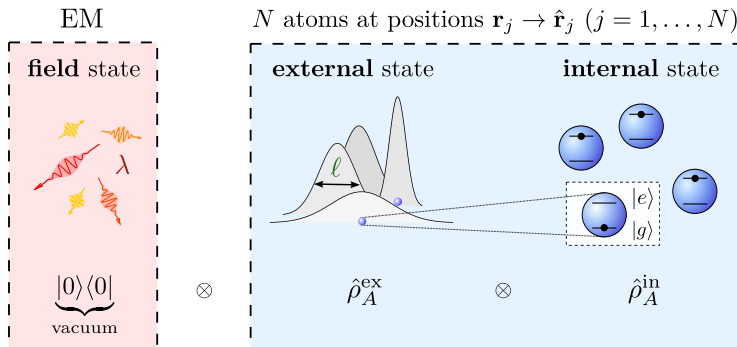
$$\eta = 2\pi \left(\frac{\ell}{\lambda} \right)$$

Lamb-Dicke parameter

ℓ : typical size of the atomic wavepacket

λ : wavelength of the emitted radiation

System



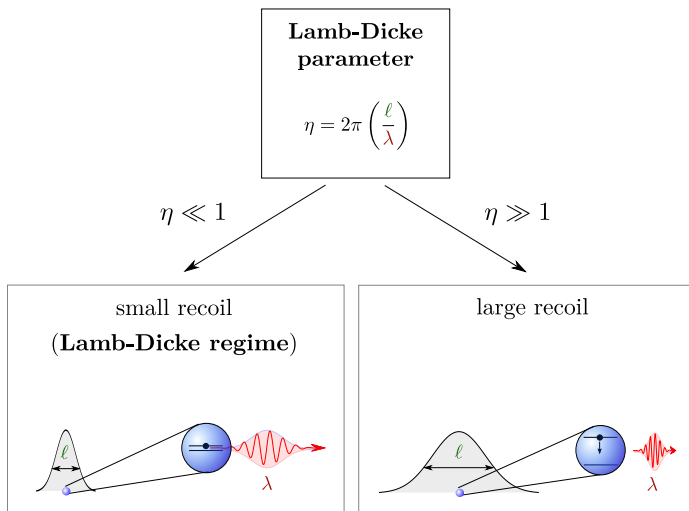
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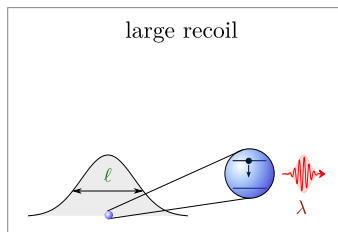
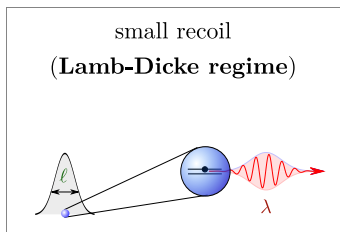
The Lamb-Dicke parameter



Description of internal dynamics

Previous works

- Most of theoretical methods
 - are restricted to the Lamb-Dicke regime
See e. g. PRL **104**, 043003 (2010); PRA **84**, 043825 (2011)
 - account either for recoil or indistinguishability, but not for both
See e. g. PRA **51**, 3959 (1995); PRA **53**, 390 (1996); PRA **59**, 3797 (1999)



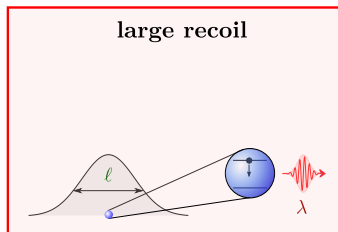
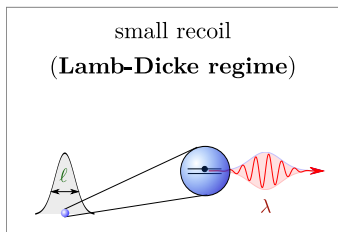
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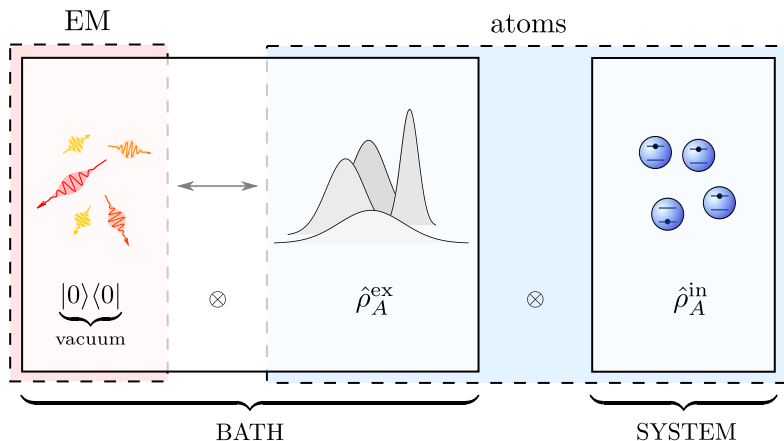
Our contribution

- **Master equation for internal degrees of freedom beyond the Lamb-Dicke regime**



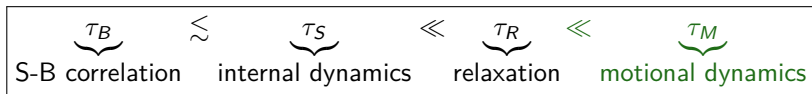
Derivation of a master equation

System-bath decomposition



Derivation of a master equation

Timescales and approximations



- $\tau_B \ll \tau_R$: Born-Markov condition (for field degrees of freedom)
⇒ no memory effects (no photon absorption)
- $\tau_S \ll \tau_R$: RWA condition
⇒ only energy conserving transitions
- $\tau_R \ll \tau_M$: Born condition (for motional degrees of freedom)
⇒ freezing of the intrinsic motional dynamics ($\hat{r}_j(t) \rightarrow \hat{r}_j$)

Master equation

$$\frac{d\hat{\rho}_A^{\text{in}}(t)}{dt} = -\frac{i}{\hbar} \left[\hat{H}_{\text{dd}}, \hat{\rho}_A^{\text{in}}(t) \right] + \mathcal{D} [\hat{\rho}_A^{\text{in}}(t)]$$

conservative part

$$\hat{H}_{\text{dd}} = \sum_{i \neq j}^N \hbar \Delta_{ij} \hat{\sigma}_+^{(i)} \hat{\sigma}_-^{(j)}$$

with dipole-dipole shifts

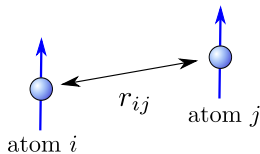
$$\Delta_{ij} = \int_{\mathbb{R}^3} \Delta^{\text{cl}}(\mathbf{r}) \mathcal{F}_{\mathbf{r}}^{-1} [C_{ij}^{\text{ex}}(\mathbf{k})] d\mathbf{r}$$

dissipative part

$$\mathcal{D}[\cdot] = \sum_{i,j}^N \gamma_{ij} \left(\hat{\sigma}_-^{(j)} \cdot \hat{\sigma}_+^{(i)} - \frac{1}{2} \left\{ \hat{\sigma}_+^{(i)} \hat{\sigma}_-^{(j)}, \cdot \right\} \right)$$

with decay rates

$$\gamma_{ij} = \int_{\mathbb{R}^3} \gamma^{\text{cl}}(\mathbf{r}) \mathcal{F}_{\mathbf{r}}^{-1} [C_{ij}^{\text{ex}}(\mathbf{k})] d\mathbf{r}$$



Master equation

dipole-dipole shifts

$$\Delta_{ij} = \int_{\mathbb{R}^3} \Delta^{\text{cl}}(\mathbf{r}) \mathcal{F}_{\mathbf{r}}^{-1} [C_{ij}^{\text{ex}}(\mathbf{k})] d\mathbf{r}$$

classical
expression

decay rates

$$\gamma_{ij} = \int_{\mathbb{R}^3} \gamma^{\text{cl}}(\mathbf{r}) \mathcal{F}_{\mathbf{r}}^{-1} [C_{ij}^{\text{ex}}(\mathbf{k})] d\mathbf{r}$$

classical
expression

Master equation

dipole-dipole shifts

$$\Delta_{ij} = \int_{\mathbb{R}^3} \Delta^{\text{cl}}(\mathbf{r}) \mathcal{F}_{\mathbf{r}}^{-1} [C_{ij}^{\text{ex}}(\mathbf{k})] d\mathbf{r}$$

decay rates

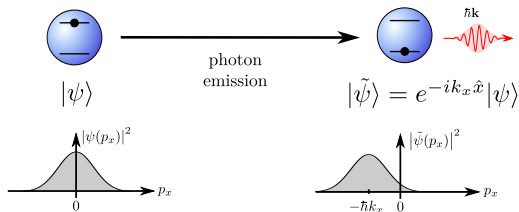
$$\gamma_{ij} = \int_{\mathbb{R}^3} \gamma^{\text{cl}}(\mathbf{r}) \mathcal{F}_{\mathbf{r}}^{-1} [C_{ij}^{\text{ex}}(\mathbf{k})] d\mathbf{r}$$

$$\mathcal{F}_{\mathbf{r}}^{-1} [C_{ij}^{\text{ex}}(\mathbf{k})] = \int_{\mathbb{R}^3} \frac{d\mathbf{k}}{(2\pi)^3} e^{i\mathbf{k}\cdot\mathbf{r}} C_{ij}^{\text{ex}}(\mathbf{k})$$

motional correlation function

$$C_{ij}^{\text{ex}}(\mathbf{k}) = \langle e^{i\mathbf{k}\cdot\hat{\mathbf{r}}_{ij}} \rangle_{\hat{\rho}_A^{\text{ex}}}$$

recoil operator



Internal dynamics

Lamb-Dicke regime : $\eta \rightarrow 0$ (classical fixed atomic positions)

$$\Delta_{ij} \longrightarrow \Delta^{\text{cl}}(\mathbf{r}_{ij}) = \frac{3\gamma_0}{4} \left[-p_{ij} \frac{\cos \xi_{ij}}{\xi_{ij}} + q_{ij} \left(\frac{\sin \xi_{ij}}{\xi_{ij}^2} + \frac{\cos \xi_{ij}}{\xi_{ij}^3} \right) \right]$$

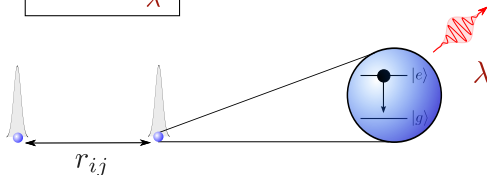
$$\gamma_{ij} \longrightarrow \gamma^{\text{cl}}(\mathbf{r}_{ij}) = \frac{3\gamma_0}{2} \left[p_{ij} \frac{\sin \xi_{ij}}{\xi_{ij}} + q_{ij} \left(\frac{\cos \xi_{ij}}{\xi_{ij}^2} - \frac{\sin \xi_{ij}}{\xi_{ij}^3} \right) \right]$$

$$\xi_{ij} = 2\pi \frac{r_{ij}}{\lambda}$$

- p_{ij}, q_{ij} angular factors which depend on the dipolar transition (π or σ^\pm)
- γ_0 : single-atom spontaneous emission rate

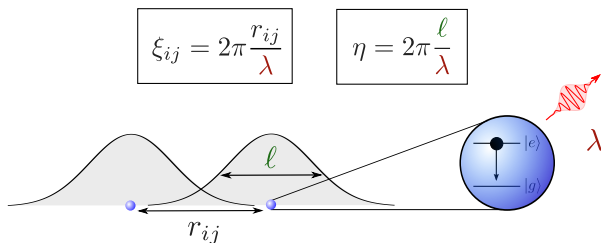
Internal dynamics in Lamb-Dicke regime ($\eta \rightarrow 0$)

$$\xi_{ij} = 2\pi \frac{r_{ij}}{\lambda}$$



Regime	Relevant Phenomena
$1 \ll \xi_{ij}$	$\gamma_{ij} \simeq \gamma_0 \delta_{ij} \Rightarrow$ independant spontaneous emissions
$\xi_{ij} \ll 1$	$\gamma_{ij} \simeq \gamma_0 \Rightarrow$ cooperative effects

Internal dynamics beyond the Lamb-Dicke regime



Regime	Relevant Phenomena
$1 \lesssim \eta \ll \xi_{ij}$	recoil
$\eta \ll 1 \ll \xi_{ij}$	independent spontaneous emissions
$\eta \lesssim \xi_{ij} \ll 1$	cooperative effects
$1 \ll \xi_{ij} \lesssim \eta$	indistinguishability, recoil
$\xi_{ij} \ll 1 \lesssim \eta$	cooperative effects, recoil, indistinguishability
$\xi_{ij} \lesssim \eta \ll 1$	cooperative effects, indistinguishability

Internal dynamics beyond the Lamb-Dicke regime

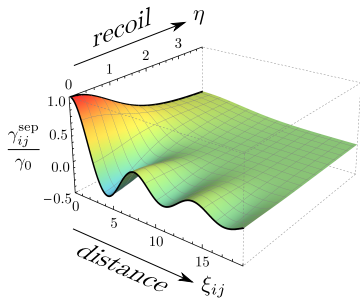
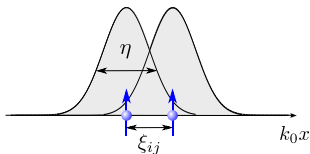
RESULTS :

$$\left. \begin{array}{l} \gamma_{ii} = \gamma_0 \quad \forall i = 1, \dots, N \\ |\gamma_{ij}| \leq \gamma_0 \quad \forall i, j = 1, \dots, N \end{array} \right\} \text{for any motional state}$$

- γ_{ij} and Δ_{ij} depend on the motional state $\hat{\rho}_A^{\text{ex}}$ through $\mathcal{C}_{ij}^{\text{ex}}(\mathbf{k}) = \langle e^{i\mathbf{k} \cdot \hat{\mathbf{r}}_{ij}} \rangle_{\hat{\rho}_A^{\text{ex}}}$
- Analytical expressions of γ_{ij} found for
 - ▶ Gaussian states
 - ▶ Fock states
 - ▶ thermal states
- Numerical computations of Δ_{ij} for Gaussian states

Motional Gaussian states - decay rates γ_{ij}

Distinguishable atoms

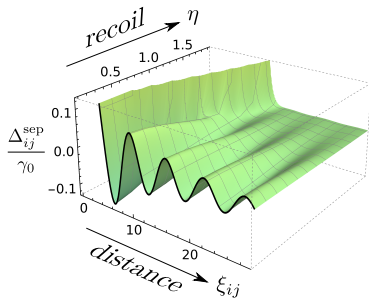
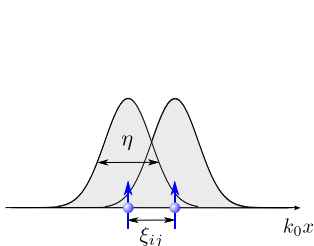


Analytical expression **valid for any** ξ_{ij} and η

$$\gamma_{ij}^{\text{sep}}(\xi_{ij}, \eta) = \frac{3\gamma_0}{16\eta^5} \left(\frac{\sqrt{\pi}}{6} e^{-\frac{\xi_{ij}^2}{4\eta^2}} [16\eta^4 - q_{ij} (4\eta^4 + 3\xi_{ij}^2 - 6\eta^2)] \operatorname{Re} \left\{ \operatorname{erf} \left(\eta + \frac{i\xi_{ij}}{2\eta} \right) \right\} \right. \\ \left. - q_{ij}\eta e^{-\eta_0^2} [2\eta^2 \cos \xi_{ij} - \xi_{ij} \sin \xi_{ij}] \right)$$

Motional Gaussian states - dipole-dipole shifts Δ_{ij}

Distinguishable atoms

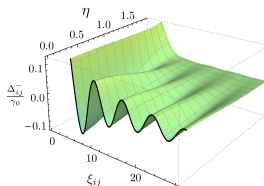
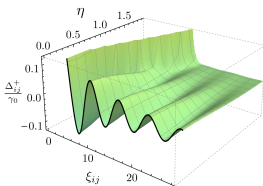
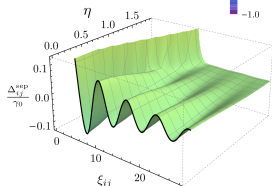
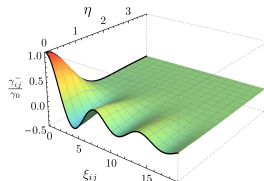
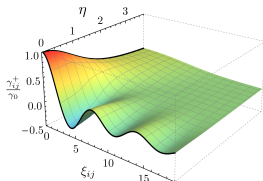
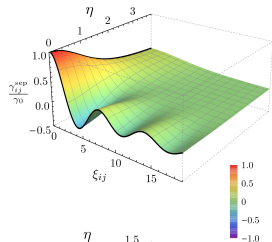


$$\Delta_{ij} = \left[\int_{-\infty}^{-\epsilon} + \int_{\epsilon}^{\infty} \right] \Delta^{\text{cl}}(x) \mathcal{F}_x^{-1} [C_{ij}^{\text{ex}}(k_x)] dx$$

- Δ_{ij} diverges when $\mathcal{F}_{x=0}^{-1} [C_{ij}^{\text{ex}}(k_x)] \neq 0$
- A cutoff ϵ must be introduced ($\epsilon \simeq a_0 = \text{Bohr radius}$)

Motional Gaussian states - γ_{ij} and Δ_{ij}

Distinguishable and indistinguishable atoms



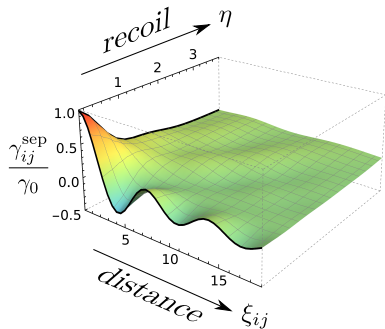
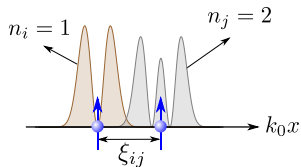
separable
motional state

symmetric
motional state

antisymmetric
motional state

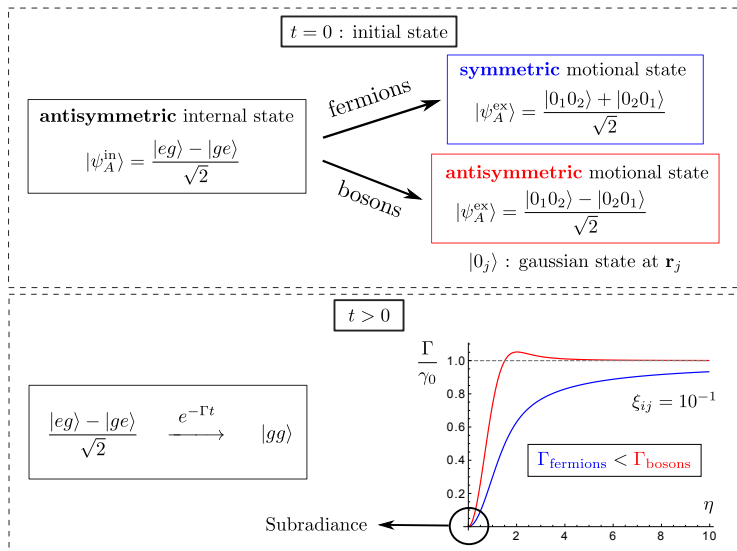
Motional Fock states - decay rates γ_{ij}

Distinguishable atoms



Application

Spatial Pauli-blocking of spontaneous emission : $N = 2$ indistinguishable atoms



Conclusion and applications

Conclusion

Derivation of a general **master equation** for internal dynamics including quantized atomic motion

- General expressions of **decay rates** and **dipole-dipole shifts** valid for any motional state
- Analytical expressions obtained for relevant motional states

Applications and outlook

Quantum programmed internal dynamics through motional state engineering

- spatial Pauli-blocked spontaneous emission
- control of cooperative processes (super and subradiance)
- tuning of the dipole-dipole couplings between atoms (Rydberg atoms)

Thanks for your attention !

More info : Phys. Rev. A **93**, 022124 (2016)