



A viscoelastic-viscoplastic-damage constitutive model based on a large strain hyperelastic formulation for amorphous glassy polymers

> V.-D. Nguyen<sup>\*</sup>, X. Morelle, F. Lani, T. Pardoen, C. Bailly and L. Noels ICCM20 19-24 July 2015, Copenhagen, Denmark (\*) vandung.nguyen@ulg.ac.be

## Introduction

- Complex behavior of glassy polymers
  - Rate-, pressure-, temperature-dependent
  - Multi-stage:
    - Viscoelastic at small strains
    - Viscoplastic at large strains
    - Irreversible saturation softening
    - Different failure points in tension and compression



Tensile behavior of epoxy (Fiedler et al. 2001)



Tensile behavior of epoxy (Gerlachet al. 2008)



Compression behavior of RTM6 under monotonic and cyclic loadings (Morelle et al. 2012)

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## Introduction

- Modeling strategy:
- Elastic+Isotropic hardening
- Kinematic hardening

Elastic+Hardening

- Elastic+Hardening+Saturation Softening
- Elastic+Hardening+Saturation Softening+Failure



• Finite deformation decomposition:



- Isotropic softening model:
  - Material degradation is modeled by the reduction of the load carrying surface with a softening parameter D

$$D = 1 - \frac{S_{\text{reduced}}}{S}$$

• Isotropic 3D problems: 
$$\hat{\mathbf{P}} = \frac{\mathbf{P}}{1-D}$$

- Constitutive modeling includes:
  - Undamaged constitutive law :  $\hat{\mathbf{P}}(t) = \hat{\mathbf{P}}(\mathbf{F}(t), \mathbf{Q}(\tau) : 0 \le \tau \le t)$
  - Evolution of the softening variable:  $D(t) = D(\mathbf{F}(t), \mathbf{Q}(\tau) : 0 \le \tau \le t)$

**Q** is a vector, contains internal variables

- Effective stress measures based on an hyperelastic approach:
  - Existence of a viscoelastic potential:  $\Psi = \Psi \left( \mathbf{E}^{ve} 
    ight)$
  - Stress measures:

$$\delta \Psi = \hat{\mathbf{P}} : \delta \mathbf{F} = \hat{\boldsymbol{\kappa}} \cdot \left( \delta \mathbf{F} \cdot \mathbf{F}^{-1} \right) = \hat{\boldsymbol{\tau}} : \delta \mathbf{E}^{ve}$$

- First Piola Kirchhoff stress tensor:  $\hat{\mathbf{P}}$
- Kirchhoff stress tensor:  $\hat{oldsymbol{\kappa}}$
- Co-rotational Kirchhoff stress tensor:  $\hat{oldsymbol{ au}}$
- Relations between the stress measures:

$$\hat{\mathbf{P}} = 2\mathbf{F}^{ve} \cdot \left(\hat{\boldsymbol{\tau}} : \frac{\partial \mathbf{E}^{ve}}{\partial \mathbf{C}^{ve}}\right) \cdot \mathbf{F}^{vp-T}$$

$$\hat{oldsymbol{\kappa}} = \hat{\mathbf{P}} \cdot \mathbf{F}^T$$

$$\hat{\kappa} = \mathbf{R}^{ve} \cdot \hat{\tau} \cdot \mathbf{R}^{veT}$$
   
in terms of the corotational stress

Same form as small

deformation theory



To avoid the problem of loss of solution uniqueness
 → implicit non-local formulation (Peerlings et al. 1996)

$$ar{\gamma}_s - l_s^2 oldsymbol{
abla}_0 \cdot oldsymbol{
abla}_0 ar{\gamma}_s = \gamma_s$$

• Saturation damage evolution:

$$\dot{D}_s = F_s \left( D_s, \mathbf{F}, \bar{\gamma}_s \right) \dot{\bar{\gamma}}_s$$

- $\gamma_s \,, \bar{\gamma}_s$  local& non-local variables associated to the saturation softening stage
  - $l_s$  saturation non-local characteristic length

• Failure stage:



 $0 \leq D_f \rightarrow 1$ when strains increase



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 $l_f$  saturation non-local characteristic length

strain

- Rate dependent behavior:
  - coupled viscoelastic/viscoplastic model for the undamaged constitutive law



## Outline

- Viscoelastic modeling
- Viscoplastic modeling
- Selection of local variables for implicit non-local formulation
- Numerical examples
- Conclusions & perspectives

- Generalized Maxwell model:
  - N+1 springs governed by N+1 bi-logarithmic potentials:

$$\Psi_i = \frac{K_i}{2} \ln^2 \left( J^{ve} \right) + G_i \operatorname{dev} \mathbf{E}^{ve} : \operatorname{dev} \mathbf{E}^{ve}, i = \infty, 1...N$$

• N dashpots governed by N dissipation functions:

$$\mathcal{D}_i, \ i = 1...N$$



• Corotational Kirchhoff stress:

$$\begin{split} \Psi &= \Psi_{\infty} + \sum_{k=1}^{N} \left( \Psi_{k} - \mathcal{D}_{k} \right) \qquad \hat{\tau} = \frac{\partial \Psi}{\partial \mathbf{E}^{ve}} = \hat{\tau}_{\infty} + \sum_{i=1}^{N} \hat{\tau}_{i} \\ \hat{\tau}_{\infty} &= \frac{\partial \Psi_{\infty}}{\mathbf{E}^{ve}} = K_{\infty} \operatorname{tr} \mathbf{E}^{ve} \mathbf{I} + 2G_{\infty} \operatorname{dev} \mathbf{E}^{ve} \\ \hat{\tau}_{i} &= \frac{\partial \Psi_{i}}{\mathbf{E}^{ve}} - \frac{\partial \mathcal{D}_{i}}{\mathbf{E}^{ve}} = K_{i} \operatorname{tr} \mathbf{E}^{ve} \mathbf{I} + 2G_{i} \operatorname{dev} \mathbf{E}^{ve} - \mathbf{q}_{i} \end{split}$$

• Evolution of internal variables  $\mathbf{q}_i, \ i=1...N$ 

$$\operatorname{dev} \dot{\mathbf{q}}_{i} = \frac{\operatorname{dev} \hat{\boldsymbol{\tau}}}{g_{i}} = \frac{2G_{i}}{g_{i}} \operatorname{dev} \mathbf{E}^{ve} - \frac{1}{g_{i}} \operatorname{dev} \mathbf{q}_{i}$$
$$\operatorname{tr} \dot{\mathbf{q}}_{i} = \frac{\operatorname{tr} \hat{\boldsymbol{\tau}}}{k_{i}} = \frac{3K_{i}}{k_{i}} \operatorname{tr} \mathbf{E}^{ve} - \frac{1}{k_{i}} \operatorname{tr} \mathbf{q}_{i}$$

Time constants:  $k_i, g_i, i = 1...N$ 

- Viscoelastic constitutive relation:
  - Deviatoric part:

$$\operatorname{dev} \hat{\boldsymbol{\tau}} = \int_{-\infty}^{t} 2G(t-s) : \frac{\mathrm{d}}{\mathrm{d}s} \operatorname{dev} \mathbf{E}^{ve}(s) \, ds$$
$$G(t) = G_{\infty} + \sum_{i=1}^{N} G_i \exp\left(-\frac{t}{g_i}\right)$$

• Volumetric part:

$$\hat{p}^{cor} = \frac{1}{3} \operatorname{tr} \hat{\boldsymbol{\tau}} = \int_{-\infty}^{t} K(t-s) : \frac{\mathrm{d}}{\mathrm{d}s} \operatorname{tr} \mathbf{E}^{ve}(s) \, ds$$
$$K(t) = K_{\infty} + \sum_{i=1}^{N} K_{i} \exp\left(-\frac{t}{k_{i}}\right)$$

- Yield surface:
  - A generalized version of the Drucker Prager yield surface with an exponent-enhanced octahedral term
  - In terms of the corotational Kirchhoff stress

$$F(\hat{\tau}, \mathbf{b}, \sigma_c, \sigma_t) = \left(\frac{\phi_e}{\sigma_c}\right)^{\alpha} - \frac{m^{\alpha} - 1}{m + 1} \frac{3\phi_p}{\sigma_c} - \frac{m^{\alpha} + m}{m + 1}$$
Current yield surface
$$\begin{cases}
F < 0 \quad \text{Elastic region} \\
F \ge 0 \quad \text{Plastic region}
\end{cases}$$

$$\phi_e = \sqrt{\frac{3}{2}} \text{dev}(\hat{\tau} - \mathbf{b}) : \text{dev}(\hat{\tau} - \mathbf{b})$$

$$\phi_p = \frac{1}{3} \text{tr}(\hat{\tau} - \mathbf{b}) \quad m = \frac{\sigma_t}{\sigma_c}$$

$$\sigma_t : \text{Tensile yield stress}$$

$$\sigma_c : \text{Compression yield stress}$$

$$Initial yield surface$$

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- Yield surface:
  - Influence of the yield exponent  $\, lpha \,$



Non-associated Perzyna-type viscoplastic flow rule:

$$\mathbf{D}^{vp} = \lambda \frac{\partial P}{\partial \hat{\tau}} \qquad \lambda = \frac{1}{\eta} < F >^{\frac{1}{p}} \qquad < F > = \begin{cases} F \text{ if } F \ge 0\\ 0 \text{ if } F < 0 \end{cases}$$

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- Yield function: F
- Plastic flow potential: P
- Viscosity parameters:  $\eta, p$



• Quadratic flow potential:

$$P=\phi_e^2+\beta\phi_p^2$$

$$\phi_e = \sqrt{\frac{3}{2} \operatorname{dev} \left(\hat{\boldsymbol{\tau}} - \mathbf{b}\right)} : \operatorname{dev} \left(\hat{\boldsymbol{\tau}} - \mathbf{b}\right)$$
$$\phi_p = \frac{1}{3} \operatorname{tr} \left(\hat{\boldsymbol{\tau}} - \mathbf{b}\right)$$

• Constant plastic Poisson ratio during plastic flow:

$$\nu_p = \frac{9 - 2\beta}{18 + 2\beta}$$

 $\beta=0 \rightarrow \,$  incompressible plastic flow is recovered

• Equivalent plastic strain:

$$\dot{\gamma} = k\sqrt{\mathbf{D}^{vp} : \mathbf{D}^{vp}} = k\lambda\sqrt{\frac{\partial P}{\partial \hat{\tau}} : \frac{\partial P}{\partial \hat{\tau}}}$$
$$k = \frac{1}{\sqrt{1+2\nu_p^2}} \qquad \beta = 0 \to k = \sqrt{\frac{2}{3}}$$

• Influence of the viscosity exponent in the uniaxial tests:



• Uniaxial compression tests



Compression behavior of RTM6 under monotonic loadings (Morelle et al. 2012)

$$\alpha = 3.5$$
  $p = 0.21$   
 $m = 0.75$   $\eta = 3.10^4$  MPa.s  
 $\beta = 0.3$   $\sigma_c = 100$  MPa



- Local variable for saturation softening stage:  $\dot{\gamma}_s = \dot{\gamma}$
- Uniaxial compression tests without considering failure damage
  - Saturation damage law:  $\dot{D}_s = H_s \bar{\gamma}_s^{n_s} \left( D_{s\infty} D_s \right) \dot{\bar{\gamma}}_s$



Experimental data from Morelle et al. 2012

- Creep tests without considering failure damage :
  - Saturation damage law:  $\dot{D}_s = H_s \bar{\gamma}_s^{n_s} \left( D_{s\infty} D_s \right) \dot{\bar{\gamma}}_s$



Experimental data from Morelle et al. 2012

- Local variable for failure stage: based on a pressure-sensitive failure criterion:
  - Failure criterion:  $F_f = \Phi_f\left(\hat{\boldsymbol{\tau}}, \hat{X}_c, \hat{X}_t\right) r \leq 0$ 
    - Power failure surface:

$$\Phi_{f} = \frac{m_{c} + 1}{m_{c}^{\alpha_{f}} + 1} \left(\frac{\hat{\tau}_{e}}{\hat{X}_{c}}\right)^{\alpha_{f}} + \frac{1 - m_{c}^{\alpha_{f}}}{m_{c}^{\alpha_{f}} + m_{c}} \frac{\operatorname{tr} \hat{\tau}}{\hat{X}_{c}} - 1$$

$$\hat{\tau}_{e} = \sqrt{\frac{3}{2}} \operatorname{dev} \hat{\tau} : \operatorname{dev} \hat{\tau}$$
s:

• Failure stresses:

$$\hat{X}_{c,t} = \hat{X}_{c,t}(\dot{\gamma})$$
 e.g. power law: 
$$\begin{aligned} \hat{X}_c &= \hat{X}_c^0 + A_c \dot{\gamma}^{\beta_c} \\ \hat{X}_t &= \hat{X}_t^0 + A_t \dot{\gamma}^{\beta_t} \end{aligned}$$

• Failure onset evolution:  $\dot{r} \ge 0, \ F_f \le 0, \ \dot{r}F_f = 0$ 

• Local variable for failure stage: 
$$\dot{\gamma}_f = \begin{cases} \dot{\gamma} \text{ if } \dot{r} > 0 \\ 0 \text{ if } \dot{r} = 0 \end{cases}$$

- Stress/strain behavior with failure:
  - Saturation damage law:  $\dot{D}_s = H_s \bar{\gamma}_s^{n_s} \left( D_{s\infty} D_s \right) \dot{\bar{\gamma}}_s$
  - Failure damage law:  $\dot{D}_f = H_f \bar{\gamma}_f^{n_f} (1 D_f) \dot{\bar{\gamma}}_f$



Denmark

- Transition from the continuum damage to a cohesive law (Nguyen V.P et al. 2010)
  - Voided structure made of epoxy resin
  - Uniaxial traction (mode I) using the orthotropic BC





Uniaxial traction test of the carbon fiber/epoxy composite





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Prescribed

displacement

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## **Conclusions & future works**

- The proposed model can capture:
  - Rate-dependent full-range stress/strain behavior of epoxy
  - Rate-dependent stress/strain behavior of composites
  - Traction/separation law can be extracted on RVEs
- Future works
  - Study the failure behavior of fiber/epoxy composites using the computational homogenization method

# Thank you for your attention!