# The massive star binary fraction in young open clusters - II. NGC 6611 (Eagle Nebula) 

H. Sana, ${ }^{1,2 \star}$ E. Gosset ${ }^{3} \dagger$ and C. J. Evans ${ }^{4}$<br>${ }^{1}$ European Southern Observatory, Alonso de Cordova 1307, Casilla 19001, Santiago 19, Chile<br>${ }^{2}$ Sterrenkundig Instituut Anton Pannekoek, Universiteit van Amsterdam, Postbus 94249, 1098 XH Amsterdam, the Netherlands<br>${ }^{3}$ Astrophysical Institute, Liège University, Bât. B5c, Allée du 6 Août 17, B-4000 Liège, Belgium<br>${ }^{4}$ UK Astronomy Technology Centre, Royal Observatory Edinburgh, Blackford Hill, Edinburgh EH9 3HJ

Accepted 2009 August 14. Received 2009 August 13; in original form 2009 May 31


#### Abstract

Based on a set of over 100 medium- to high-resolution optical spectra collected from 2003 to 2009, we investigate the properties of the O-type star population in NGC 6611 in the core of the Eagle Nebula (M16). Using a much more extended data set than previously available, we revise the spectral classification and multiplicity status of the nine O-type stars in our sample. We confirm two suspected binaries and derive the first SB2 orbital solutions for two systems. We further report that two other objects are displaying a composite spectrum, suggesting possible long-period binaries. Our analysis is supported by a set of Monte Carlo simulations, allowing us to estimate the detection biases of our campaign and showing that the latter do not affect our conclusions. The absolute minimal binary fraction in our sample is $f_{\min }=0.44$ but could be as high as 0.67 if all the binary candidates are confirmed. As in NGC 6231 (see Paper I), up to 75 per cent of the O star population in NGC 6611 are found in an $\mathrm{O}+\mathrm{OB}$ system, thus implicitly excluding random pairing from a classical IMF as a process to describe the companion association in massive binaries. No statistical difference could be further identified in the binary fraction, mass-ratio and period distributions between NGC 6231 and NGC 6611, despite the difference in age and environment of the two clusters.


Key words: binaries: close - binaries: spectroscopic - stars: early-type - open clusters and associations: individual: NGC 6611.

## 1 INTRODUCTION

At a distance of $\sim 2 \mathrm{kpc}$, NGC 6611 at the core of the Eagle Nebula (M16) is one of the most famous star-formation regions of the Southern sky. Popularized by the dramatic Hubble Space Telescope (HST) image of Hester et al. (1996), the striking morphology of the nebula is the result of the feedback of the massive stars in its core on the original natal cloud. In the last 6 Myr, M16 seems to have developed a continuous star-formation history, resulting in a diverse stellar population that encompasses high- and intermediate-mass stars, low-mass pre-main sequence stars, young embedded objects, maser sources and Herbig-Haro objects. As such, the Eagle Nebula is also one of the best places to study the positive and negative influence of young massive stars to trigger new formation events. Oliveira (2008) has recently reviewed star formation in the Eagle Nebula, so that we only provide here some of the key facts that are of direct interest for the present work.

[^0]Several authors have identified variable extinction across the cluster and an abnormal reddening law. These properties have strongly affected the distance determination and a wide range of values, from 1.7 to 3.2 kpc , can be found in the literature. Yet, all the recent works seem to converge towards the closest distances. In the following, we adopt $d=1.8 \pm 0.1 \mathrm{kpc}$ as found by Dufton et al. (2006) from the spectroscopic parallaxes of 24 OB stars.

The cluster age is $2-3$ Myr typically, although the age spread is probably significant and extends from 1 to 6 Myr. The most recent age determinations (Bonatto, Santos \& Bica 2006; Martayan et al. 2008) seem to favour an even younger age around 1.5 Myr . Regarding the cluster mass, Bonatto et al. (2006) argued that the OB-type stars alone amount for a minimal mass of $\sim 1.6 \times 10^{3} \mathrm{M}_{\odot}$, although Wolff et al. (2007) suggested that the total mass could be as large as $\sim 2.5 \times 10^{4} \mathrm{M}_{\odot}$.

As the second paper in our series, this work mainly focuses on the multiplicity properties of the O-type star population in NGC 6611. Yet, given that those objects are the main source of radiative and kinetic energy of the Eagle Nebula, the present study has implications beyond the sole properties of the O-type stars in the cluster. It forms a prerequisite to any detailed understanding
of the energetic balance of the mechanisms at work in the nebula.

With 13 O-type stars and about 50 B0 to B5 stars, the young open cluster NGC 6611 at the core of the Eagle Nebula hosts a rich earlytype population. Compared to NGC 6231 (Sana et al. 2008, hereafter Paper I), the O star population in NGC 6611 is definitely younger and contains earlier spectral types. While several authors have investigated the multiplicity of these objects, the picture is still not complete. Bosch, Morrell \& Niemelä (1999) obtained multi-epoch low- ( $R \approx 4000$ ) and medium- to high-resolution ( $R \approx 9000-$ 15000 ) spectroscopy of 10 of the brightest OB stars and reported three definite binaries and three multiple candidates. Subsequent medium- to high-resolution ( $R \approx 7000-29000$ ) spectroscopy by Evans et al. (2005) and Martayan et al. (2008) confirmed two of the known O-type binaries, further suggesting an additional SB1 candidate not observed by Bosch et al. (1999). All in all, the minimum binary fraction $f_{\min }$ is $3 / 13 \approx 0.23$. It could, however, be much higher given that some objects lack multi-epoch monitoring and that four other binary candidates have been identified.
On the very large separation side, the adaptive optics survey of Duchêne et al. (2001) shed a complementary light on the multiplicity properties of the cluster. Focusing in the separation range from 200 to 3000 au ( $\approx 0.1-1.5 \mathrm{arcsec}$ ) around 60 OB-type cluster members, they identified low-mass visual companions in $18 \pm 6$ per cent of the cases, two of which are around O-type stars.

Finally, Gvaramadze \& Bomans (2008) have recently identified three O-type stars that have most likely been ejected from NGC 6611. Given the number of known O stars in the cluster (either isolated or within a multiple object), those ejected stars would amount to about 20 per cent of the initial O-type star population in NGC 6611. Gvaramadze \& Bomans (2008) further suggested that the three detected runaway stars might only represent about one-fifth of the ejected O stars. This would imply that NGC 6611 already displays a significant dynamical evolution. We emphasize that the present work only addresses the multiplicity properties of the current O-type star population in the cluster.

This work is organized as follows. Section 2 describes our observing campaign and the data reduction process. Section 3 revises the properties of the individual O-type stars in the cluster. Section 4 presents an analysis of the observational biases using Monte Carlo simulations. Finally, Section 5 discusses the present results, and Section 6 summarizes our conclusions.

## 2 OBSERVATIONS AND DATA HANDLING

The core of our data set is formed by 66 high-resolution feros spectra obtained from 2004 May to 2006 June at the ESO/MPG 2.2-m telescope at La Silla (PI: Sana). The feros spectrograph provides us with the complete optical spectrum (3800-9200 $\AA$ ) at once, with a spectral resolving power of 48000 . The feros properties are identical to those given in Paper I, and the data reduction process is described in detail in Sana (2009). Exposure times ranged between 10 and 45 min depending on the object magnitude, yielding a typical signal-to-noise ratio (SNR) of 200 as measured in the continuum close to $5000 \AA$.

A second part of our data set is composed of 10 FLAMES-GIRAFFE spectra, five feros spectra and six flames-uves spectra from the VLT-FLAMES survey of massive stars (PI: Smartt). Those data and the data reduction have been previously described by Evans et al. (2005). Additionally, we retrieved one feros spectrum of HD 168075 and one of HD 168076 (PI: Bouret) from the ESO archive. We also retrieved three UVES long-slit echelle spectra

Table 1. Brief comparison of the properties of the various spectrographs used to assemble our data set. The last column provides the number of epochs $n$ obtained with each spectrograph.

| Spectrograph | $\lambda \lambda$ range $(\AA)$ | Resolving power | $n$ |
| :--- | :--- | :--- | :---: |
| FEROS $_{\text {FLAMES-GIRAFFE }}$ a | $3800-9200$ | 48000 | 78 |
| FLAMES-UVES | $4140-6210$ | $20850-29600$ | 10 |
| UVES | $3280-4560$ | 47000 | 6 |
| ISIS | $3800-5100,6200-6800$ | 70000 | 3 |

${ }^{a}$ Each of the GIRAFFE epochs is actually composed of several observations using different wavelength settings. See Evans et al. (2005) for more information.
of HD 168076 (PI: Andre). The two FEROS spectra were reduced as described above while the uves data were reduced using the ESO CPL-based pipeline and normalized by fitting low-order polynomials to the continuum of the individual orders.

Finally, five objects were re-observed with Feros in March 2009 (PI: Evans) to improve the detection likelihood of long-period systems. These additional data were reduced as described above and have a typical SNR of $250-300$ at $5000 \AA$. Table 1 gives a brief overview of the wavelength coverage and spectral resolution of the different spectra in our data set.
Using the reduced spectra, Doppler shifts and equivalent widths (EWs) were measured by simultaneously fitting one to three Gaussians (depending on whether an SB1, SB2 or SB3 signature was visible) to a series of line profiles (Table 2). Effective rest wavelengths were taken from Conti, Leep \& Lorre (1977) and from Underhill (1994) for lines below and above $4800 \AA$, respectively, to compute the radial velocities (RVs). Table 3 provides the journal of the observations and lists, for each spectral line, the measured RVs.
The achieved RV accuracy is strongly dependent on the width of the lines, thus on the projected rotation rate of the star. Using a purely empirical approach, Fig. 17 in this work and fig. 10 in Paper I provide us with the following guidelines. For slow rotators ( $v \sin i<50 \mathrm{~km} \mathrm{~s}^{-1}$ ), an rms accuracy of $1 \mathrm{~km} \mathrm{~s}^{-1}$ or better can be achieved. For more typical rotation rates ( $v \sin i \sim 100-200 \mathrm{~km} \mathrm{~s}^{-1}$ ), typical dispersions in the measured velocities are between 2 and $5 \mathrm{~km} \mathrm{~s}^{-1}$. For fast rotators ( $v \sin i>300 \mathrm{~km} \mathrm{~s}^{-1}$ ), the precision of the measurements drops significantly, with rms dispersions close to $15 \mathrm{~km} \mathrm{~s}^{-1}$.
As in Paper I, the spectral classification is based on the quantitative criteria of Conti \& Alschuler (1971), Conti (1973), Mathys (1988) and Mathys (1989), that rely on the EWs of given diagnostic lines. We adopt the following notations: $\log W^{\prime}=\log W(\lambda 4471)-$ $\log W(\lambda 4542), \log W^{\prime \prime}=\log W(\lambda 4089)-\log W(\lambda 4144)$ and $\log W^{\prime \prime \prime}=\log W(\lambda 4388)+\log W(\lambda 4686)$, where the EWs are expressed in mA. The first criterion, $\log W^{\prime}$, can be used across the entire O-type range. $\log W^{\prime \prime}$ and $\log W^{\prime \prime \prime}$ are restricted to stars strictly later that O6 and O8, respectively, while $\log W_{\lambda 4686}$ can be used for O8 stars and earlier. The measurements and the corresponding spectral types are given in Table 4 and discussed in Section 3. Measured EWs for binary components only provide a lower limit on the real strength of the lines. This affects the latter two criteria, so that, in these cases, the values quoted in Table 4 for $\log W^{\prime \prime \prime}$ and $\log W_{\lambda 4686}$ are only lower limits.

## 3 O-TYPE STARS IN NGC 6611

This section briefly discusses the data associated with each Otype object in our sample, with a particular emphasis on the

Table 2. Number of RV measurements obtained for each object with respect to the considered spectral lines. '-' means that the corresponding line has not been measured.

| Object | $\begin{array}{r} \mathrm{He}_{\mathrm{I}} \\ \lambda 4026 \end{array}$ | $\begin{array}{r} \text { Si IV } \\ \lambda 4089 \end{array}$ | $\begin{array}{r} \mathrm{He} \text { I } \\ \lambda 4144 \end{array}$ | $\begin{array}{r} \mathrm{He} \text { I } \\ \lambda 4388 \end{array}$ | $\begin{array}{r} \mathrm{He}_{\mathrm{I}} \\ \lambda 4471 \end{array}$ | $\begin{array}{r} \mathrm{Mg}_{\mathrm{II}} \\ \lambda 4481 \end{array}$ | $\begin{array}{r} \mathrm{He}_{\text {II }} \\ \lambda 4542 \end{array}$ | $\begin{array}{r} \mathrm{He}_{\text {II }} \\ \lambda 4686 \end{array}$ | $\begin{array}{r} \mathrm{He} \text { I } \\ \lambda 4922 \end{array}$ | $\begin{array}{r} \mathrm{He}_{\text {II }} \\ \lambda 5412 \end{array}$ | $\begin{array}{r} \mathrm{He}_{\mathrm{I}} \\ \lambda 5876 \end{array}$ | $\begin{array}{r} \mathrm{He}_{\mathrm{I}} \\ \lambda 7065 \end{array}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{BD}-13^{\circ} 4923$ | 11 | 11 | 10 | 12 | 12 | 10 | 12 | 13 | 9 | 10 | 10 | 9 |
| $\mathrm{BD}-13^{\circ} 4927$ | 7 | 7 | - | 8 | 8 | 8 | 8 | 8 | 6 | 8 | 6 | 6 |
| $\mathrm{BD}-13^{\circ} 4928$ | 7 | 7 | 7 | 8 | 8 | - | 8 | 8 | 6 | 6 | 6 | 6 |
| $\mathrm{BD}-13^{\circ} 4929$ | 11 | 11 | 11 | 11 | 12 | - | 12 | 12 | - | 10 | 10 | - |
| $\mathrm{BD}-13^{\circ} 4930$ | 7 | 7 | 7 | 8 | 8 | 8 | 8 | 8 | 6 | 6 | 6 | 5 |
| HD 168075 | 9 | 9 | 9 | 9 | 11 | 11 | 11 | 11 | 11 | 11 | 9 | 8 |
| HD 168076 | 9 | - | - | - | 7 | - | 7 | 7 | - | 7 | 7 | - |
| HD 168137 | 14 | 14 | 14 | 14 | 14 | 14 | 14 | 14 | 13 | 14 | 14 | 12 |
| HD 168183 | 13 | 13 | 13 | 13 | 13 | - | 13 | 13 | 13 | 13 | 13 | 13 |

spectral properties and on the multiplicity status. As mentioned earlier, NGC 6611 hosts 13 O-type stars. Table 5 provides the cross-identification between HD/BD numbers and the numeration of Walker (1961), Evans et al. (2005, hereafter ESL05) and Guarcello et al. (2007, hereafter GPM07). It also lists the stellar parameters obtained by Dufton et al. (2006) from atmosphere model fitting. The four faintest O-stars were not re-observed as part of our monitoring program, so that only one or two spectra are available for each object. Those four stars are thus omitted in the rest of this paper.

This section is organized as follows. We first present the gravitationally bound systems (Section 3.1). In Section 3.2, we discuss objects whose spectra most likely present multiple signatures, although our current data set does not allow us to conclude whether the different components are gravitationally bound or arise from spurious alignment. Finally, Section 3.3 summarizes the properties of the remaining, presumably single, O-type stars.

### 3.1 Gravitationally bound systems

### 3.1.1 BD-13영ㅇ (W175)

Classified as $\mathrm{O} 5 \mathrm{~V}\left(\left(\mathrm{f}^{*}\right)\right)$ by Hillenbrand et al. (1993), BD-13 4923 was first suspected to be an $\mathrm{O}+\mathrm{O}$ binary with a period larger than 8 d by Bosch et al. (1999). Based on the fact that the $\mathrm{He}_{\mathrm{I}}$ and $\mathrm{He}_{\text {I }}$ lines were moving in the opposite direction, these authors suggested that the primary was an early O star and the secondary, a late one, a fact later confirmed by Evans et al. (2005). Duchêne et al. (2001) further reported a low-mass visual companion at $0.67 \mathrm{arcsec}(\approx 1200 \mathrm{au})$, with an estimated mass between 2 and $4 \mathrm{M}_{\odot}$.

We collected 10 additional spectra over 3 yr . At least four of them clearly show the signatures of the two components both in the He I and $\mathrm{He}_{\text {II }}$ lines (Fig. 1), allowing us to put additional constraints on the orbital and physical properties of the components. Using the well-separated spectra only, we refined the spectral classification to $\mathrm{O} 4 \mathrm{~V}((\mathrm{f}))+\mathrm{O} 7.5 \mathrm{~V}$ with, respectively, the O 5 and O 7 spectral subtype at $1 \sigma$. The third, more distant, component is likely a late-B/early-A star. It is thus expected to contribute to less than 1 per cent of the total continuum level and will not affect the RV and EW measurements of the two O stars. It has thus been neglected in the present analysis.

Because we only have a limited data set, we also made use of previously discarded data. First, one of our feros spectra was obtained late during morning twilight and is contaminated by the solar spectrum. From $\mathrm{He}_{\text {II }} \lambda 4686$ and redward, the solar and stellar spectra can however be easily separated, allowing us to recover most of the prominent stellar lines in that region. Secondly, the WHT-IsIS spectrum from Evans et al. (2005) exhibits a clear wave-
length calibration shift in the $\lambda 4800 \AA$ setting. To solve this issue, we have cross-calibrated the isIs wavelength solution with respect to our own data, by using the narrow DIBs at $\lambda \lambda 4763,4780$ and 4964 (Hobbs et al. 2008). This allowed us to achieve a corrected wavelength calibration with an rms dispersion better than $0.1 \AA$ ( $\sim 5 \mathrm{~km} \mathrm{~s}^{-1}$ ), sufficient thus for our purpose.

In the following, we focus explicitly on the $\mathrm{He}_{\text {II }} \lambda 4686$ line which is the only strong SB2 line common to all our spectra. To improve the disentangling of heavily blended profiles, we adopted the line shapes as determined on widely separated spectra obtained during the same run, thus fitting the Doppler shifts only. Although it is not clear whether our data cover any of the extrema of the RV curve, we attempted to constrain the orbital period using the Fourier method of Heck, Manfroid \& Mersch (1985), as revisited by Gosset et al. (2001). The obtained periodogram shows two dominant peaks at $v \approx 0.075$ and $0.078 \mathrm{~d}^{-1}$ (Fig. 2). Including the $\mathrm{He}_{\text {II }}$ measurements of Bosch et al. (1999) further allows us to determine that the signal at $v \approx 0.075 \mathrm{~d}^{-1}$ is most likely to correspond to the true periodicity of the orbit.

Adopting the corresponding period $P \approx 13.2679 \mathrm{~d}$ as a first guess, we used the Liège Orbital Solution Package (LOSP; Sana \& Gosset 2009) to compute a preliminary SB2 orbital solution. Table 6 and Fig. 3 respectively give the best-fitting orbital parameters and the corresponding RV curves. We note that the period uncertainties given in Table 6 only refer to the fitted models. As such, they do not account for the presence of multiple peaks in the periodogram. Both solutions, computed with and without including literature data, are in good agreement although Fig. 3 reveals some systematic deviations between the best fit solution and the RV points of Bosch et al. (1999). Yet, the latter measurements might suffer from significant systematic errors because the SB 2 signature was not separated. With $e \approx 0.3$, the obtained solution is significantly eccentric. The derived mass ratio is in perfect agreement with the estimated spectral types of the components. Given the computed minimal masses, one would predict an orbital inclination about $30^{\circ}$ to $40^{\circ}$, so that no eclipses are expected.

### 3.1.2 HD 168075 (W197)

So far classified as an O6-O7 V((f)) star (Evans et al. 2005), HD 168075 was first suggested to be a binary by Conti et al. (1977). Later on, Bosch et al. (1999) reported no short-term variations but mentioned significant changes in data taken 2-yr apart. These authors further reported the detection of $\mathrm{Si}_{\text {III, }}, \mathrm{N}_{\text {III }}$ and $\mathrm{O}_{\text {II }}$ absorption lines normally not present in a typical O7 spectrum.
 available in the electronic version of the article - see Supporting Information.

| $\begin{array}{r} \text { HJD } \\ -2400000 \end{array}$ | $\begin{array}{r} \mathrm{He}_{\mathrm{I}} \lambda 4026 \\ 4026.072 \end{array}$ | Si iv $\lambda 4089$ 4088.863 | $\begin{array}{r} \mathrm{He}_{\mathrm{I}} \lambda 4144 \\ 4143.759 \end{array}$ | $\begin{array}{r} \mathrm{He}_{\mathrm{I}} \lambda 4388 \\ 4387.928 \end{array}$ | $\begin{array}{r} \mathrm{He}_{\mathrm{I}} \lambda 4471 \\ 4471.512 \end{array}$ | $\begin{array}{r} \mathrm{Mg}_{\text {II }} \lambda 4481 \\ 4481.228 \end{array}$ | $\begin{array}{r} \text { Не }_{\text {II }} \lambda 4542 \\ 4541.590 \end{array}$ | $\begin{array}{r} \text { He II } \lambda 4686 \\ 4685.682 \end{array}$ | $\begin{array}{r} \mathrm{He}_{\mathrm{I}} \lambda 4922 \\ 4921.929 \end{array}$ | $\begin{array}{r} \text { Не пI } \lambda 5412 \\ 5411.520 \end{array}$ | $\begin{array}{r} \mathrm{He}_{\mathrm{I}} \lambda 5876 \\ 5875.620 \end{array}$ | $\begin{array}{r} \text { He I } \lambda 7065 \\ 7065.190 \end{array}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathrm{BD}-13^{\circ} 4923$ prim |  |  |  |  |  |  |  |  |  |  |  |  |
| $52836.70812^{d}$ | - | - | 82.4 | - | 62.7 | 55.1 | 68.6 | 91.7 | - | - | - | - |
| $52839.75137^{d}$ | - | - | - | - | 60.5 | 52.6 | 75.2 | 92.2 | - | - | - | - |
| $53134.95474{ }^{e}$ | - | - | - | - | - | - | - | 11.2 | - | 11.3 | 42.0 | - |
| $53509.84538{ }^{\text {a }}$ | -79.4 | - | - | - | -63.7 | - | -64.4 | -56.3 | - | -63.8 | -70.5 | - |
| $53510.87032^{a}$ | - | - | - | - | - | - | -4.1 | -11.5 | - | -0.2 | 24.6 | - |
| $53511.76887^{a}$ | - | - | - | - | - | - | 18.2 | 25.0 | - | 18.4 | -2.9 | - |
| $53512.75277^{a}$ | - | - | - | - | 47.3 | 44.6 | 35.6 | 49.9 | - | 37.9 | 54.9 | - |
| $53860.83796^{a}$ | 80.7 | - | - | - | 58.8 | - | 69.0 | 81.9 | - | 70.3 | 62.2 | - |
| $53861.72041^{a}$ | - | - | 31.8 | - | 43.4 | 56.1 | 39.1 | 78.7 | 49.7 | 41.3 | 56.2 | - |
| $53862.72524^{a}$ | - | - | 22.0 | - | - | 46.2 | 27.5 | 59.5 | - | 27.2 | 42.7 | - |
| $53863.84924^{a}$ | - | - | - | - | - | - | 16.3 | 42.2 | - | 14.2 | 6.3 | - |
| $53864.75611^{a}$ | - | - | - | - | - | -43.6 | -3.5 | 24.9 | -52.7 | -2.1 | 17.8 | - |
| $\mathrm{BD}-13^{\circ} 4923 \mathrm{sec}$ |  |  |  |  |  |  |  |  |  |  |  |  |
| $52836.70812^{d}$ | -21.1 | -85.5 | -83.0 | -79.8 | -86.0 | -88.3 | -126.1 | -80.5 | - | - | - | - |
| $52839.75137^{d}$ | - | - | - | -81.5 | -86.3 | -89.1 | -101.5 | -76.6 | - | - | - | - |
| $53134.95474{ }^{e}$ | - | - | - | - | - | - | - | - | - | - | -50.6 | - |
| $53509.84538^{\text {a }}$ | 161.0 | 157.8 | 149.1 | 56.0 | 152.3 | 142.0 | 164.0 | 158.3 | 149.8 | 174.8 | 153.9 | 53.2 |
| $53510.87032^{a}$ | 4.3 | 63.5 | 64.2 | 60.9 | 15.1 | - | - | 49.9 | 55.2 | - | - | -9.6 |
| $53511.76887^{a}$ | -9.1 | -21.5 | -27.6 | -23.1 | -12.2 | -14.3 | - | -11.5 | -19.6 | - | - | -9.6 |
| $53512.75277^{a}$ | -27.0 | -69.3 | -73.9 | -68.5 | -74.2 | -81.1 | - | -82.5 | -70.8 | - | -71.7 | -71.5 |
| $53860.83796^{a}$ | -85.0 | -93.0 | - | -97.7 | -101.5 | -102.1 | -139.1 | -95.8 | -96.7 | -115.1 | -96.1 | -99.3 |
| $53861.72041^{a}$ | -19.9 | -79.2 | -80.9 | -74.3 | -88.9 | -84.4 | - | -88.3 | -81.1 | - | -77.1 | -79.5 |
| $53862.72524^{a}$ | -13.5 | -49.1 | -50.3 | -46.0 | -32.8 | -50.9 | - | -67.2 | -44.4 | - | -46.4 | -44.4 |
| $53863.84924^{a}$ | -6.7 | 3.6 | 5.5 | -2.6 | 2.9 | -1.7 | - | -32.6 | 3.4 | - | - | 6.6 |
| $53864.75611^{a}$ | -0.1 | 62.5 | 56.3 | 58.0 | 16.4 | 46.7 | - | - | 57.9 | - | - | 54.1 |

$\overline{{ }^{a} \text { ESO2 } 2 \mathrm{~m}+\mathrm{FEROS} ;{ }^{b} \text { VLT + UVES; }{ }^{c} \text { VLT + FLAMES-UVES; }{ }^{d} \text { VLT + FLAMES-GIRAFFE; }{ }^{e} \text { WHT + ISIS }, ~ . ~}$

Table 4. Diagnostic line ratios and corresponding spectral types for the O-type objects studied. $n / a$ indicates criteria that either do not apply to the considered spectral type or for which the lines could not be securely measured and/or separated. We refer to text for discussion of the individual cases and to Table 9 for the finally adopted spectral types. The errorbars give the $1 \sigma$ dispersions on the measurements.

| Object | Component | $\log W^{\prime}$ | $\log W^{\prime \prime}$ | $\log W_{\lambda 4686}$ | $\log W^{\prime \prime \prime}$ | Sp. Type |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| BD $-13^{\circ} 4923$ | Primary | $-0.612 \pm 0.130$ | $n / a$ | $>2.751 \pm 0.071$ | $n / a$ | O4 V((f)) |
| BD $-13^{\circ} 4923$ | Secondary | $0.125 \pm 0.063$ | $n / a$ | $>2.147 \pm 0.182$ | $n / a$ | 07.5 V |
| BD $-13^{\circ} 4927$ | $n / a$ | $-0.020 \pm 0.014$ | $n / a$ | $2.348 \pm 0.041$ | $n / a$ | O7 III(f) |
| BD $-13^{\circ} 4928$ | $n / a$ | $0.541 \pm 0.057$ | $0.147 \pm 0.151$ | $n / a$ | $5.531 \pm 0.109$ | 09.5 V |
| BD $-13^{\circ} 4929$ | Primary | $-0.044 \pm 0.086$ | $n / a$ | $>2.637 \pm 0.063$ | $n / a$ | O7 V |
| BD $-13^{\circ} 4930$ | $n / a$ | $0.272 \pm 0.066$ | $-0.299 \pm 0.021$ | $n / a$ | $5.333 \pm 0.009$ | 08.5 V |
| HD 168075 | Primary | $-0.096 \pm 0.030$ | $n / a$ | $>2.837 \pm 0.012$ | $n / a$ | O6.5 V((f)) |
| HD 168076 | Composite | $-0.721 \pm 0.061$ | $n / a$ | n/a | $n / a$ | O4 V((f)) |
| HD 168137 | Primary | $-0.075 \pm 0.039$ | $n / a$ | $>2.644 \pm 0.033$ | $n / a$ | 06.5 V |
| HD 168137 | Secondary | $0.202 \pm 0.068$ | $n / a$ | $>2.598 \pm 0.046$ | $n / a$ | O8 V |
| HD 168183 | Primary | $0.608 \pm 0.026$ | $0.034 \pm 0.031$ | $n / a$ | $>5.270 \pm 0.021$ | O9.5 III |

Table 5. Cross-identification and physical properties of the O-type stars in NGC 6611. The magnitudes are from Guarcello et al. (2007), except for the two brightest objects for which data from Hillenbrand et al. (1993) are quoted. The stellar parameters come from Dufton et al. (2006). The uncertainties are the $1 \sigma$ errorbars quoted by the authors.

| HD | BD | Walker | ESL05 | GPM07 | V | $B-V$ | $\begin{aligned} & T_{\text {eff }} \\ & (\mathrm{kK}) \end{aligned}$ | $\begin{aligned} & \log g \\ & \left(\mathrm{~cm} \mathrm{~s}^{-2}\right) \end{aligned}$ | $\begin{aligned} & \log L \\ & \left(\mathrm{~L}_{\odot}\right) \end{aligned}$ | $\begin{aligned} & v \sin i \\ & \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $168183{ }^{\text {a }}$ | $-13^{\circ} 4991$ | W412 | 6611-001 | 24374 | 8.18 | 0.34 | 32.0 | 3.60 | 5.39 | 142 |
| $168076{ }^{\text {a }}$ | $-13^{\circ} 4926$ | W205 | 6611-002 | 27436 | 8.18 | 0.43 | 41.5 | 3.90 | 5.86 | 102 |
| $168075^{\text {a }}$ | $-13^{\circ} 4925$ | W197 | 6611-003 | 18360 | $8.752 \pm 0.037$ | $0.543 \pm 0.107$ | 40.0 | 3.90 | 5.63 | 87 |
| $168137{ }^{a}$ | $-13^{\circ} 4932$ | W401 | 6611-004 | 13706 | $8.942 \pm 0.006$ | $0.407 \pm 0.025$ | 37.0 | 4.00 | 4.85 | 76 |
| - | $-13^{\circ} 4930^{b}$ | W367 | 6611-006 | 13867 | $9.368 \pm 0.023$ | $0.343 \pm 0.079$ | 34.0 | 4.10 | 4.84 | 43 |
|  |  |  |  |  |  |  | 31.3 | 4.00 | 4.81 | 20 |
| - | $-13^{\circ} 4927$ | W246 | 6611-008 | 6835 | $11.114 \pm 0.541^{c}$ | $0.669 \pm 0.607$ | 36.0 | 3.50 | 5.75 | 100 |
| - | $-13^{\circ} 4929^{a}$ | W314 | 6611-011 | 19208 | $9.803 \pm 0.203$ | $0.706 \pm 0.208$ | 36.0 | 4.20 | 5.24 | 66 |
| - | $-13^{\circ} 4923$ | W175 | 6611-014 | 5890 | $10.007 \pm 0.011$ | $0.979 \pm 0.019$ | - | - | - | - |
| - | $-13^{\circ} 4928$ | W280 | 6611-015 | 1806 | $10.044 \pm 0.039$ | $0.560 \pm 0.048$ | 32.5 | 3.90 | 4.77 | 410 |
| - | - | W166 | 6611-017 | 18715 | $10.296 \pm 0.003$ | $0.728 \pm 0.010$ | 36.0 | 3.95 | 5.02 | 95 |
| - | - | W161 | 6611-029 | 5818 | $11.215 \pm 0.007$ | $1.241 \pm 0.008$ | 36.0 | 3.85 | 5.37 | 135 |
| - | - | W584 | 6611-045 | 5510 | $12.052 \pm 0.008$ | $1.235 \pm 0.010$ | 35.0 | 4.00 | 5.04 | 25 |
| - | - | W222 | 6611-080 | 3820 | $12.968 \pm 0.008$ | $1.519 \pm 0.009$ | 40.0 | 4.00 | 5.33 | 95 |

${ }^{a}$ For those objects, the composite nature of their spectrum was not known, thus not taken into account while deriving the stellar parameters. The latter might thus be indicative only.
${ }^{b}$ The entries on the second row quote the results obtained by Hunter et al. (2007), who re-analysed the BD $-13^{\circ} 4930$ spectrum and obtained slightly different stellar parameters.
${ }^{c}$ Hillenbrand et al. (1993) quote $V=9.46$, suggesting that $\mathrm{BD}-13^{\circ} 4923$ is brighter than reported by Guarcello et al. (2007), although marginally in agreement within the (rather large) errorbars from the latter.

Our data sample, formed by about 10 spectra collected over three years, confirms that this object is a long period binary. As found by Bosch et al., we clearly detect $\mathrm{Si}_{\text {III, }} \mathrm{N}_{\text {III }}$ and $\mathrm{O}_{\text {II }}$ lines in absorption, whose motion is anticorrelated with the $\mathrm{He}_{\text {I }}$ and $\mathrm{He}_{\text {II }}$ line motion, thus confirming their association with the secondary star. As observed in Fig. 4, the Не clear profile changes, likely due to the blending of the primary and secondary signatures, while the Не II $\lambda 4542$ line profile seems to remain constant in shape. This strongly suggests the secondary to be a B0-B1 star. Using the He i $\lambda 4471$ over He II $\lambda 4542$ EW ratio, we adopt an $\mathrm{O} 6.5 \mathrm{~V}(\mathrm{f})$ ) spectral type for the primary, with a O 7 subtype well within $1 \sigma$. We finally note that the $\mathrm{He}_{\text {I }} \lambda 4471$ line EW might have been overestimated because of the secondary blend, so that the primary could be slightly hotter than the adopted spectral subtype.

Our data alone do not allow us to definitely constrain the orbital properties of HD 168075 . Using a much larger data set, Barbá et al. (2009) recently proposed a preliminary orbital solution, with a period of 43.6 d and a slightly eccentric orbit $e=0.17$. They inde-
pendently confirmed the $\mathrm{O} 6.5 \mathrm{~V}((\mathrm{f}))$ classification for the primary and proposed that the secondary is a B0.2 III star.

### 3.1.3 HD 168137 (W401)

HD 168137 is known as an O 8.5 V isolated star. Between 2004 May and 2006 May, we acquired 10 feros spectra and, in 2009 March, we obtained one additional feros observation. Our prime data set is completed by one GIRAFFE spectrum and one UVES spectrum from Evans et al. (2005) obtained in 2003. While our 2004 May, 2006 and 2009 observations are very similar to the spectra of Evans et al. (2005), the May 2005 data reveal a clear SB2 signature (Fig. 5). No further RV variations could be detected within the 2005 campaign. Similarly, no variations are seen between the 2006 May and June spectra, suggesting a period larger than 30 d . Using the wellseparated spectra only, we revise the spectral classification to O7 $\mathrm{V}+\mathrm{O} 8 \mathrm{~V}$, with the O 6.5 and 07.5 subclasses at $1 \sigma$, respectively, for the primary and secondary components. A long-term monitoring is


Figure 1. $\mathrm{BD}-13^{\circ} 4923$ : Не $\quad \lambda \lambda 4471,5876$ and $\mathrm{He}_{\text {II }} \lambda \lambda 4542,4686$ line profiles at different epochs. The heliocentric Julian dates at mid-exposure are given at right-hand side in format HJD-2 450000 . The spectra have been shifted along the $y$-axis for clarity. The vertical dashed lines indicate the adopted rest-wavelength for each spectral line.


Figure 2. BD $-13^{\circ} 4923$ : periodograms (upper row) and spectral windows (lower row) computed using the data from this work alone (left-hand column) or combined with the Bosch et al. measurements (right-hand column).
required to bring further constraints on the orbital properties of this system.

### 3.1.4 HD 168183 (W412)

At $\sim 13$ arcmin South-East of the centre of the cluster, HD 168183 is a known SB1 eclipsing binary with a period close to 4 d (Hipparcos Catalogue 1997; Bosch et al. 1999). The object has been classified O9.5 I by Hillenbrand et al. (1993) and Bosch et al. (1999), and B0 III by Evans et al. (2005). With a significant He II $\lambda 4542$ line in absorption, HD 168183 cannot be a B star. Indeed, the EW ratio

Table 6. $\mathrm{BD}-13^{\circ} 4923$ : best-fit orbital solutions using the data set from the present work alone or combined with the Bosch et al. RV measurements (primary only). $T$ (in HJD-2 450000 ) is the time of periastron passage and is adopted as phase $\phi=0.0$ in Fig. 3. Quoted uncertainties correspond to $1 \sigma$ errorbars.

| Parameter | This work | Combined |
| :--- | :--- | :--- |
| $P(\mathrm{~d})$ | $13.2722 \pm 0.0027$ | $13.2677 \pm 0.0007$ |
| $e$ | $0.285 \pm 0.030$ | $0.302 \pm 0.045$ |
| $\omega\left({ }^{\circ}\right)$ | $199.3 \pm 9.5$ | $188.5 \pm 9.8$ |
| $T$ | $2995.491 \pm 10.144$ | $3005.016 \pm 0.376$ |
| $\gamma_{1}\left(\mathrm{~km} \mathrm{~s}^{-1}\right)$ | $30.7 \pm 4.3$ | $28.7 \pm 4.6$ |
| $\gamma_{2}\left(\mathrm{~km} \mathrm{~s}^{-1}\right)$ | $3.4 \pm 5.3$ | $7.5 \pm 6.0$ |
| $K_{1}\left(\mathrm{~km} \mathrm{~s}^{-1}\right)$ | $81.0 \pm 5.3$ | $82.6 \pm 7.0$ |
| $K_{2}\left(\mathrm{~km} \mathrm{~s}^{-1}\right)$ | $140.2 \pm 9.2$ | $142.8 \pm 12.2$ |
| $q=M_{2} / M_{1}$ | $0.578 \pm 0.043$ | $0.578 \pm 0.043$ |
| $M_{1} \sin ^{3} i\left(\mathrm{M}_{\odot}\right)$ | $8.3 \pm 1.5$ | $8.6 \pm 2.1$ |
| $M_{2} \sin ^{3} i\left(\mathrm{M}_{\odot}\right)$ | $4.8 \pm 0.8$ | $5.0 \pm 1.2$ |
| $\mathrm{rms}\left(\mathrm{km} \mathrm{s}^{-1}\right)$ | 7.9 | 12.9 |



Figure 3. BD $-13^{\circ} 4923$ : RV-curves corresponding to the orbital solutions of Table 6. Left-hand column: using data from this work. Right-hand column: combined with the Bosch et al. measurements. Filled and open symbols indicate the primary and secondary RV measurements from this work while crosses give the primary RVs from Bosch et al.


Figure 4. HD 168075: Не ${ }_{\mathrm{I}} \lambda \lambda 4471,5876$ and $\mathrm{He}_{\text {II }} \lambda \lambda 4542$, 4686 line profiles at different epochs. Note the changing asymmetry of the He I lines.
of Не I $\lambda 4471$ to Не ${ }_{\text {II }} \lambda 4542$ points to an O9.5 spectral type. With He i $\lambda 4144$ only slightly shallower than Si iv $\lambda 4089$, plus He if $\lambda 4686$ and $\mathrm{H} \alpha$ in absorption, the properties of the spectrum of HD 168183 are definitely not those of a supergiant. Given the brightness of HD 168183 , we finally adopt a O9.5 III spectral classification for the primary.
Thanks to their high SNR, our data reveal the secondary signature for the first time. Though faint, the secondary spectrum is indeed clearly visible in the Не $~ \lambda 5876$ line, as a broad and very shallow component, peaking at $\approx 2$ per cent of the continuum level (Fig. 6).


Figure 5. HD 168137: Не ${ }_{I} \lambda \lambda 4471,5876$ and $\mathrm{He}_{\text {II }} \lambda \lambda 4542,4686$ line profiles at different epochs. Note the clear SB2 signature at HJD-2 450000 ~ 3509.8 and 3512.9.


Figure 6. HD 168183 : Не II $\lambda 4686$ and $\mathrm{He}_{\text {II }} \lambda 5876$ line profiles at different epochs. Note the faint SB2 signature at HJD-2 $450000 \sim 3512.7$ and 3861.7.

Because of the large difference in the line EWs (a factor ~20 for He I $\lambda 5876$ ), indicating a significant difference in flux, and because we could not detect the secondary signature in the Нe it lines, we suggest that the secondary is a mid- B type star.

As for BD $-13^{\circ} 4923$, we used the Fourier technique of Heck et al. (1985) and Gosset et al. (2001) to estimate the orbital period of the system. Resulting periodograms are shown in Fig. 7 and present a clear peak at $v \approx 0.25 \mathrm{~d}^{-1}$. Using LOSP, we first computed the primary orbital solution. Significantly better residuals are obtained using a limited eccentricity compared to a circular solution. We


Figure 7. HD 168183: periodograms (upper row) and spectral windows (lower row) computed using the He I data from this work alone (left-hand column) or combined with the Bosch et al. measurements (right-hand column).
derived a RV-curve semi-amplitude $K_{1}$ and a mass function $f_{\text {mass }}$ larger than those proposed by Bosch et al. (1999). Combining our data with theirs yields essentially the same best-fitting solution (Table 7).

Assuming that the primary orbital solution is correct and that both components share the same systemic velocity, it is in principle possible to constrain the secondary orbit with a single RV measurement. For the purpose of this exercise, we adopt the Не $\boldsymbol{\lambda} \lambda 5876$ orbital solution of Table 7, propagating the uncertainties on the primary parameters at the first order by means of the theory of error propagation (Bevington 1969). Using the two observations, where the SB2 signature is well seen, we estimate $K_{2}=279.1 \pm$ $10.9 \mathrm{~km} \mathrm{~s}^{-1}$. This yields $q=M_{2} / M_{1}=0.301 \pm 0.013, M_{1} \sin ^{3} i=$ $15.3 \pm 1.9 \mathrm{M}_{\odot}$ and $M_{2} \sin ^{3} i=4.6 \pm 0.4 \mathrm{M}_{\odot}$. The corresponding RV curves are displayed in Fig. 8. While this result supposes that the adopted time series alias corresponds to the true period, we note that the obtained mass ratio is in good agreement with the component spectral types as previously discussed. Under such assumption, the orbital inclination would be $i \approx 65^{\circ}$, compatible with the presence of eclipses as detected by hipparcos.

### 3.2 Composite systems

## 

Classified B0 V by Hillenbrand et al. (1993) and Bosch et al. (1999), the star's spectral type was revised to O 9 V by Evans et al. (2005). The latter authors further noted slight profile asymmetries, suggesting that $\mathrm{BD}-13^{\circ} 4929$ is an SB2 binary system. Our new data allow us to shed new light on the nature of BD $-13^{\circ} 4929$. From Fig. 9, it is obvious that $\mathrm{BD}-13^{\circ} 4929$ is an SB3 system, formed by a shortperiod binary with nearly identical early B companions displaying broad and shallow lines, and a third O-type star. Because the O

Table 7. HD 168183: orbital solutions obtained using different data sets: the RVs of the He i $\lambda 5876$ line ( Col .2 ), the average of the $\mathrm{He}_{\text {I }}$ measurements ( Col .3 ) and the average of the $\mathrm{He}_{\text {I }}$ measurements combined with the Bosch et al. data (Col. 4). The usual notations have been used. $T$ (in HJD-2 450000 ) is the time of periastron passage and is adopted as phase $\phi=0.0 \mathrm{in}$ Fig. 8. Quoted uncertainties correspond to $1 \sigma$ errorbars.

|  | He $\quad \lambda 5876$ line | He I line | Combined |
| :--- | :--- | :--- | :--- |
| $P(\mathrm{~d})$ | $4.01554 \pm 0.00017$ | $4.01556 \pm 0.00015$ | $4.015575 \pm 0.000022$ |
| $e$ | $0.039 \pm 0.015$ | $0.052 \pm 0.014$ | $0.059 \pm 0.015$ |
| $\gamma\left(\mathrm{~km} \mathrm{~s}^{-1}\right)$ | $19.2 \pm 1.0$ | $19.3 \pm 0.9$ | $18.5 \pm 0.8$ |
| $K\left(\mathrm{~km} \mathrm{~s}^{-1}\right)$ | $83.9 \pm 1.5$ | $78.5 \pm 1.3$ | $77.8 \pm 1.3$ |
| $T$ | $2998.718 \pm 0.388$ | $2998.786 \pm 0.244$ | $2998.881 \pm 0.195$ |
| $\omega\left({ }^{\circ}\right)$ | $13.9 \pm 35.8$ | $20.1 \pm 22.7$ | $28.9 \pm 17.5$ |
| $f_{\text {mass }}$ | $0.245 \pm 0.013$ | $0.200 \pm 0.010$ | $0.195 \pm 0.009$ |
| rms $\left(\mathrm{km} \mathrm{s}^{-1}\right)$ | 3.3 | 2.8 | 3.7 |



Figure 8. HD 168183: RV-curves corresponding to orbital solutions of Table 7, based on the Не г $\lambda 5876$ line data set (left) and on the average of the $\mathrm{He}_{\mathrm{I}}$ line RVs combined with Bosch et al. measurements (right). Filled and open circles show primary and secondary RV measurements from this work while filled triangles show the Bosch et al. data.
star has much sharper line profiles, it is easily disentangled from the close $\mathrm{B}+\mathrm{B}$ pair using multi-Gaussian profile fitting. Based on the $\mathrm{He}_{\text {I }} \lambda 4471$ to $\mathrm{He}_{\text {II }} \lambda 4542$ line ratio, we assign the O star an O 7 spectral subtype, with the 07.5 type at $1 \sigma$. With no sign of emission lines plus the fact that the observed (diluted) EW of the He II $\lambda 4686$ is $\log W_{\lambda 4686}=2.64 \pm 0.06(\mathrm{~m} \AA)$, we confirm the dwarf luminosity class. The signatures of the fainter companions are not visible in the $\mathrm{He}_{\text {II }} \lambda 4542$ line, but show a faint contribution to Не II $\lambda 4686$. We thus adopt a B0.5 spectral subtype for both companions of the O star. Given the object magnitude, we consider that the $B+B$ pair is also formed by two dwarfs. While in reality the two B stars are not exactly identical as indicated by slight differences in their EWs (Fig. 9), more data would be needed to refine the classification further.

The $B+B$ binary has a probable period close to $4 d$, but our data do not allow us to put quantitative constraints on the orbit of the tight system. Similarly, the O-type component does not show any significant RV variations and we cannot decide whether the O star is gravitationally linked to the $\mathrm{B}+\mathrm{B}$ pair or not. We however note that the chance of spurious alignment of an O and a $B$ star in NGC 6611 is very small. $\mathrm{BD}-13^{\circ} 4929$ is thus an excellent candidate of hierarchical triple system formed by a close $\mathrm{B} 0.5 \mathrm{~V}+\mathrm{B} 0.5 \mathrm{~V}$ pair plus a wider O 7 V companion.

### 3.2.2 HD 168076 (W205)

Close to the cluster centre, HD 168076 is the brightest optical source in NGC 6611. Reported as an SB candidate by Conti et al. (1977) based on two discrepant RV points, further investigations by Bosch


Figure 9. $\mathrm{BD}-13^{\circ} 4929$ : Не ${ }_{\mathrm{I}} \lambda \lambda 4471,5876$ and $\mathrm{He}_{\text {II }} \lambda \lambda 4542,4686$ line profiles at different epochs. The close $\mathrm{B}+\mathrm{B}$ pair signature is clearly visible in the He I lines at HJD-2 450000 from 3862.8 to 3863.7.
et al. (1999) could not confirm its binary nature. We collected only a few spectra in 2006 that show no variability, neither on short timescales nor, by comparison with similar quality data from Evans et al. (2005) and from 2009 May, on longer time-scales (Fig. 11). The EW ratio of the $\mathrm{He}_{\mathrm{I}} \lambda 4471$ and $\mathrm{He}_{\text {II }} \lambda 4542$ lines and the presence of He II $\lambda 4686$ strongly in absorption point towards an $\mathrm{O} 4 \mathrm{~V}((\mathrm{f}))$ classification, as previously adopted by various authors.
Bosch et al. (1999) however reported a relatively bright companion at 0.15 arcsec from HD 168076. Given the aperture on the sky of the FEROS fibre $(\oslash=2.0 \mathrm{arcsec})$, the spectrograph thus provides us with a blended spectrum of the two stars and the classification criteria used so far are not applicable. In a second approach, we used the O2-O4 spectroscopic atlas of Walborn et al. (2002) to refine the spectral classification of HD 168076. Because of the presence of Nv in absorption and of Si iv in emission with moderate strength, but no N iv $\lambda 4058$ (Fig. 10), the spectrum of HD 168076 looks very much like the one of HD 93128 and we thus adopt an $\mathrm{O} 3.5 \mathrm{~V}((\mathrm{f}+))$ spectral classification for the earliest object of the pair.


Figure 10. HD 168076: spectrum in the range 4000-4700 Å. Note the Si iv $\lambda \lambda 4089-4116$ and $\mathrm{N}_{\text {III }} \lambda \lambda 4634-4640-4642$ lines in emission and the N v $\lambda \lambda 4604-4620$ absorption.


Figure 11. HD 168076: Не ${ }_{\text {I }} \lambda 4471$ and Не II $\lambda \lambda 4542,4686$ line profiles at different epochs.

Dufton et al. (2006) reported masses of 38 and 16 for both objects, which is likely to have been underestimated given that a typical O3.5 V star is about $52 \mathrm{M}_{\odot}$ (Martins, Schaerer \& Hillier 2005). Assuming that the quoted mass ratio $(q=0.42)$ is correct, the secondary component could be an 07.5 V or an O 9 III . The expected luminosity ratio would correspond to 5 and 3.5 respectively, which is sufficient in both cases to significantly affect the apparent $\mathrm{He}_{\mathrm{I}}$ to $\mathrm{He}_{\text {II }}$ line ratio in the composite spectrum. Because it is unlikely that the secondary evolved more quickly than the primary and because previous mass-transfer episode appears very unlikely in such a young cluster, we finally adopt the 07.5 V classification for the secondary. Under these assumptions, and given the measured angular separation from Bosch et al. (1999), we estimate a minimal orbital period of several hundreds of years.

### 3.3 Presumably single stars

### 3.3.1 BD $-13^{\circ} 4927$ (W246)

Reported as O7 II(f) by Hillenbrand et al. (1993), no significant RV variation was detected by Bosch et al. (1999). We collected five FEROS spectra from 2004 to 2006, plus one spectrum in 2009, and we further included GIRAFFE data from 2007 in our analysis. With $W_{\lambda 4471} \approx W_{\lambda 4542}$, this object is definitely an O7 star. Small RV variations ( $\sim 10 \mathrm{~km} \mathrm{~s}^{-1}$ peak-to-peak) are detected, although at the limit of significance, and can easily be mimicked by low-amplitude wind


Figure 12. $\mathrm{BD}-13^{\circ} 4927$ : Не ${ }_{\mathrm{I}} \lambda 4471$, Не ${ }_{\text {II }} \lambda \lambda 4542,4686$ and $\mathrm{H} \alpha$ line profiles at different epochs. Note the variable emission that partially fills the Не II $\lambda 4686$ and $\mathrm{H} \alpha$ lines.


Figure 13. $\mathrm{BD}-13^{\circ} 4928$ : Не ${ }_{\mathrm{I}} \lambda \lambda 4471,5876$ and $\mathrm{He}_{\text {II }} \lambda \lambda 4542,4686$ line profiles at different epochs. Note the shallow He iI $\lambda 4542$ line.
effects. Because $\log W_{\lambda 4686}<2.75$ (m $\AA$ ), and $\mathrm{He}_{\text {II }} \lambda 4686$ and $\mathrm{H} \alpha$ are partly filled with emission (Fig. 12), the adopted spectroscopic criteria indicate a giant luminosity class. Yet, with $\mathrm{N}_{\text {III }} \lambda \lambda 4634-40-$ 42 in emission, the spectrum of $\mathrm{BD}-13^{\circ} 4927$ is much more alike HD 151515 (O7 II(f)) than HD 93522 (O7 III((f))). Combined with the fact that BD-13 4927 displays a magnitude midway between O7 III and O7 I typical objects (Martins et al. 2005), we confirm the $\mathrm{O} 7 \mathrm{II}(\mathrm{f})$ spectral classification. From our campaign, the star is likely to be single.

### 3.3.2 BD-13 4928 (W280)

Located close to the cluster centre, $\mathrm{BD}-13^{\circ} 4928$ was reported as a single O9.5 V star by Bosch et al. (1999). We obtained seven additional feros spectra that revealed rotationally broadened lines (Fig. 13). Because of this, very few continuum windows were
available in the $\lambda \lambda 4050-4700 \AA$ region of the spectrum. In this range, we used the spectra obtained just before or just after those of $\mathrm{BD}-13^{\circ} 4928$ to empirically correct for the instrument response curve. Using this approach, we routinely achieved a precision better than half a percent of the continuum in the response curve correction, only leaving a residual slope in the $\mathrm{BD}-13^{\circ} 4928$ spectrum. The latter slope, due to the difference of effective temperature between the two objects, was finally removed by fitting a low-order polynomial to the continuum of the pre-normalized spectra.

Minor line profile variations seem to be present but can originate from the object's rapid rotation. Because the lines are rotationally broadened, their profiles deviate significantly from a Gaussian shape. In the present case, we used a rotation profile to measure the Doppler shifts, fitting simultaneously the line centre, its amplitude and the star's projected rotational velocity $v \sin i$. The RVs obtained show a rms dispersion close to $10 \mathrm{~km} \mathrm{~s}^{-1}$ for rotational velocities in the range $360-390 \mathrm{~km} \mathrm{~s}^{-1}$. Based on the $W^{\prime}$ and $W^{\prime \prime \prime}$ criteria, we confirm the 09.5 V classification.

## 

So far classified as O9.5 V (Hillenbrand et al. 1993), Bosch et al. (1999) noted that the star had a peculiar velocity compared to the cluster systemic velocity, thus questioning the membership of BD-13 4930 to NGC 6611 .
We acquired six feros spectra, complemented by two giraffe spectra and one feros spectrum from Evans et al. (2005). With an internal RV dispersion of about $1 \mathrm{~km} \mathrm{~s}^{-1}$ only, we could not detect any variability in our data. $\mathrm{BD}-13^{\circ} 4930$ is definitely a slow rotator. Most of the lines display a steep Lorentzian profile although some lines, such as He i $\lambda 5876$ definitely show a Gaussian profile. The $\mathrm{He}_{\text {II }}$ lines further display a broader profile than the $\mathrm{He}_{\text {I }}$ lines. Using rotationally broadened profiles for comparison, we estimated that the core of the $\mathrm{He}_{\mathrm{I}}$ lines is not compatible with a projected rotational velocity larger than $40 \mathrm{~km} \mathrm{~s}^{-1}$. This is in line with earliest measurements by Dufton et al. (2006) and Hunter et al. (2007) (see Table 5). The Не ${ }_{\text {II }} \lambda 4686$ line core indicates $v \sin i \approx 50 \mathrm{~km} \mathrm{~s}^{-1}$; the Нe ${ }_{\text {II }} \lambda 5412$ line core $v \sin i \approx 60 \mathrm{~km} \mathrm{~s}^{-1}$, while the He ${ }_{\text {II }} \lambda 4542$ core yields $v \sin i \gtrsim 100 \mathrm{~km} \mathrm{~s}^{-1}$. While we use a purely rotational profile for comparison, we note that the broadening of the $\mathrm{He}_{\text {II }}$ lines might rather be due to the Stark broadening.
While clearly displaying the $\mathrm{He}_{\text {I }} \lambda \lambda 4542,4686$ and 5412 lines typical of an O-type star (Fig. 14), BD $-13^{\circ} 4930$ also displays a large number of metallic sharp lines typical of late-O/early-B stars (Fig. 15). Considering only the $\mathrm{He}_{\text {I }}$ and $\mathrm{He}_{\text {II }}$ lines, our spectral classification criteria suggest an 08.5 V spectral type, with the O 9 V subtype well within $1 \sigma$. FEROS spectra of two spectral standards from Walborn \& Fitzpatrick (1990), HD 93028 (O9 V) and $\tau$ Sco (B0.2 V ), are available in the feros archive from ESO. We retrieved both of them and reduced them as described in Section 2. Fig. 15 provides a direct comparison of the three spectra. Clearly, BD $-13^{\circ} 4930$ displays properties at mid-course between HD 93028 and $\tau$ Sco. We thus adopt O 9.5 V as our final classification.

## 4 OBSERVATIONAL BIASES

RV techniques alone have proved to be an efficient method to detect high-mass spectroscopic binaries. Yet, the parameter space to probe is very large. It covers 3-4 orders of magnitude in period, two orders of magnitude in mass (both for the primary and secondary) and the full range of eccentricities from circular orbits to eccentricities of


Figure 14. $\mathrm{BD}-13^{\circ} 4930$ : Не $\mathrm{I} \lambda \lambda 4471,5876$ and Не ${ }_{\text {II }} \lambda \lambda 4200,4542,4686$ line profiles at different epochs.
0.9 or above. Naturally, RV techniques are not equally sensitive in all the regions of the parameter space. This not only results from the amplitude of the RV signal but also from the very different time sampling needed to uncover a short-period system with a high mass-ratio versus a long period, eccentric object with a lower mass secondary component. As a consequence, depending on the probed location in the parameter space, a given time sampling does not allow us to unveil multiple systems with the same probability.
To better estimate the observational biases of our campaign, we have run a series of Monte Carlo simulations using two different approaches. In the following simulations, we adopted a $\Delta R V$ threshold for the primary star of $20 \mathrm{~km} \mathrm{~s}^{-1}$ to consider a binary as detected.

Estimating the detection biases is unfortunately an ill-posed problem as the probability of detection in a given part of the parameter space and for a given observational sampling, needs to be convolved with the distribution of the orbital parameters which, for massive stars, are still poorly constrained. Still, reasonable estimates of the detection biases can be obtained using assumptions that are based on the results of Paper I. The period distribution is chosen to be biuniform in $\log$ scale, between $\log P=0.3$ and 1.0 (d), and between 1.0 and 3.5 (d), with 60 per cent of the systems in the former interval. The eccentricity distribution is taken to be uniform between 0.0 and 0.9 , with the additional constraint that the separation at periastron passage should be larger than $20 \mathrm{R}_{\odot}$ in order to avoid contact. The mass-ratio distribution is chosen to be uniform between 0.2 and 1.0.
Adopting the actual observing sampling corresponding to each object, we randomly drew 10000 orbital configurations, adopting a random orientation of the orbital plane in space and a random time of periastron passage. Table 8 reports the detection probabilities ( $P_{\text {detect }}$ ) for short-, intermediate- and long-period binaries, as well as for the full period range. It indicates that our campaign is very sensitive in the short-period regime ( $P_{\text {detect }}>0.99$ ) while its performances remain very good up to $1 \mathrm{yr}\left(P_{\text {detect }} \approx 0.8-0.9\right)$. Beyond that, the detection probability tends to drop to $0.5-0.6$ typically. Yet, given that only 15 per cent of the spectroscopic binary population is expected to have periods between 1 and 10 yr , our lower sensitivity in this range has a limited impact on the total detection probability as shown by the last column of Table 8.
As a second step, to check the consistency of the estimated individual observational biases in the context of a sample of several stars, we randomly drew nine orbital configurations for nine O stars whose masses match the observed primary masses in NGC 6611.


Figure 15. BD $-13^{\circ} 4930$ spectrum in the range $4000-4700 \AA$ compared to the spectra of HD 98028 (O9 V) and $\tau$ Sco (B0.2 V).

Table 8. Binary detection probability for the time sampling associated with different objects (Col. 1) and for various period ranges (Cols. 2 to 5).

| Time sampling | Short <br> $(2-10 \mathrm{~d})$ | Intermed. <br> $(10-365 \mathrm{~d})$ | Long <br> $(365-3000 \mathrm{~d})$ | All <br> $(2-3000 \mathrm{~d})$ |
| :--- | :---: | :---: | :---: | :---: |
| $\mathrm{BD}-13^{\circ} 4923$ | 0.997 | 0.924 | 0.697 | 0.933 |
| $\mathrm{BD}-13^{\circ} 4927$ | 0.994 | 0.897 | 0.695 | 0.925 |
| $\mathrm{BD}-13^{\circ} 4928$ | 0.990 | 0.829 | 0.534 | 0.883 |
| $\mathrm{BD}-13^{\circ} 4929$ | 0.995 | 0.819 | 0.521 | 0.878 |
| $\mathrm{BD}-13^{\circ} 4930$ | 0.993 | 0.911 | 0.690 | 0.928 |
| HD 168075 | 0.995 | 0.913 | 0.667 | 0.848 |
| HD 168076 | 0.995 | 0.925 | 0.690 | 0.932 |
| HD 168137 | 0.995 | 0.942 | 0.727 | 0.943 |
| HD 168183 | 0.991 | 0.876 | 0.508 | 0.891 |



Figure 16. Distribution of the number of detected SB systems as obtained from 10000 realizations of a cluster with nine O-type binaries.

We then applied the existing time sampling of our O-type star observations, associating each time series to an object of the simulated cluster, then deducing an observed binary fraction. Repeating this for 10000 simulated clusters, Fig. 16 shows the distribution of the detected binary fraction. It predicts that our campaign should have been able to achieve a detection rate of $0.93_{-0.10}^{+0.07}$ [ 0.85 confidence interval], leaving little room for undetected spectroscopic binaries.


Figure 17. Dispersion of the Не ${ }_{\mathrm{I}} \lambda 5876$ RV measurements versus their mean RV. Circles around the $\sigma-\mu$ point indicate the gravitationally bound binaries while squares show the objects with a composite spectrum. The average systemic velocities have been used instead of the mean RV for the two binaries with orbital solutions: BD $-13^{\circ} 4923$ and HD 168183 .

To conclude this section, we note that our simulations do not take into account line blending, so that our adopted detection threshold might be too optimistic for SB2 binaries with similar mass and flux. Yet, a semi-amplitude of the RV curve of several $10 \mathrm{~km} \mathrm{~s}^{-1}$ would already result in significant changes ( $>20$ per cent) of the line depths, so that the binary nature of those objects would be unveiled by their line profiles rather than by their RV variations. Only some of the long-period systems ( $q>0.8, P>1000 \mathrm{~d}$ ) would likely remain undetected. This additional shading of the parameter space decreases the total detection probabilities given in Table 8 and Fig. 16 by a couple of per cent, which does not qualitatively affect our conclusions.

## 5 DISCUSSION

### 5.1 Membership

Fig. 17 displays the observed $1 \sigma$ RV dispersions of the Не ${ }_{\mathrm{I}} \lambda 5876$ line in our sample, plotted versus the average RV measurements of the same line. In such a diagram, the vertical axis gives us an indication of the variability while the horizontal axis provides
us with some indication of the cluster systemic velocity. Beyond the four gravitationally bound binaries, it is BD $-13^{\circ} 4928$ that shows the largest apparent variability. Yet, as previously discussed, the star is undergoing rapid rotation, thus broadening the spectral lines and making the exact determination of the line centroid less accurate. The remaining stars present an rms dispersion of their RV measurements of a few $\mathrm{km} \mathrm{s}^{-1}$ at most and, indeed, all of them are considered to be isolated or very long period objects.
Both BD $-13^{\circ} 4930$ and HD 168076 show smaller peculiar velocities compared to the average cluster velocity $\left(\approx 17 \mathrm{~km} \mathrm{~s}^{-1}\right)$. Yet we note that HD 168076 is a possible long period binary, so that the observed shifts might be additional hints of an ongoing physical process rather than indications of non-membership. Assuming $\mathrm{BD}-13^{\circ} 4930$ is not a member of NGC 6611, the following discussion would mostly remain unaffected.

### 5.2 Binary fraction

With four definite SB2 gravitationally bound systems out of nine objects in our sample, the minimal binary fraction is $f_{\text {min }} \sim 0.44$. As discussed in Section 3.2, two objects display an apparently composite spectrum and could be long period binaries if physically bound. Summing up those binary candidates would increase the binary fraction $f$ up to 0.67 . From Table 8 , one can further conclude that there is little room for non-detection biases for the remaining three stars, making them likely single. For the sake of the present discussion, we will adopt a formal binary fraction of $f=0.55 \pm$ 0.11 in our sample.

Using the same reasoning as in Paper I, and assuming that the O stars in our sample are randomly drawn from a binomial parent distribution, we infer a minimal binary fraction for the parent distribution of 0.33 at the 0.01 significance level. We note that both the binary fraction and the constraints on the parent distribution parameters are compatible with the results obtained for NGC 6231 (see Paper I). Yet the detailed estimate of the detection biases done in Section 4 allows us to address another aspect of the parent distribution. Fig. 16 reveals that the likelihood to detect five or fewer binaries under the hypothesis that all the O stars in NGC 6611 are binaries amounts to 0.003 only. As a consequence, our observations
are not compatible with a contemporaneous multiplicity rate of 100 per cent.

### 5.3 Orbital parameters

While our campaign cannot allow us to derive the complete set of orbital parameters of all the SB2 systems, the periods and mass ratios are reasonably well constrained (Table 9). The binary fraction in NGC 6611 is similar to the one in NGC 6231 (Paper I). Yet, NGC 6611 lacks the overabundant short-period population seen in NGC 6231. This however might result from the small size of our sample. A Kolmogorov-Smirnof test does indeed not authorize us to reject the hypothesis that the two samples are taken from the same underlying distribution. Similarly, the distributions of mass ratios in the two clusters are in good agreement and are both compatible with a uniform distribution from $q=1$ to $0.2-0.3$ (the observational limit for secondary detection in massive spectroscopic binaries). Finally, we note that the location, in the period-eccentricity diagram, of the systems with known eccentricity would remain well within the area normally populated by O-type binaries.

### 5.4 Companion mass function

With two $\mathrm{O}+\mathrm{O}$ and two $\mathrm{O}+\mathrm{B}$ binaries, plus two additional $\mathrm{O}+\mathrm{OB}$ candidates, the secondary companion of an O-type star in NGC 6611 is strongly biased towards higher masses. Indeed, only three systems, the presumably single stars, could still hide undetected, lowmass companions. As we pointed out in Paper I, the mass of an O star companion cannot be randomly drawn from a normal initial mass function.

### 5.5 Distance

Using the new spectral classification listed in Table 9, we estimated the spectroscopic distance moduli of each star. This ultimate check allows not only to test the agreement between the spectral types and the object brightnesses, but also for an independent estimate of the cluster distance. For this purpose, we adopted the calibration of Martins \& Plez (2006) for O stars and the one of Humphreys

Table 9. Final spectroscopic classification and multiplicity properties of the studied O stars in NGC 6611. 'SB2O' means that an orbital solution has been obtained; 'SB2E', that the system displays eclipses in addition to its SB2 signature.

| Object | Mult. | Evans et al. | Sp. Type This work | $q=M_{2} / M_{1}$ | $P$ <br> (d) | $e$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Gravitationally-bound systems |  |  |  |  |  |  |
| BD $-13^{\circ} 4923$ | SB2O | O5 V((f+))+O | O4 V((f+))+O7.5 V | 0.6 | 13.3 | 0.3 |
| HD 168075 ${ }^{\text {a }}$ | SB2O | O6-7 V((f)) | O6.5 V((f))+B0-1 V | 0.487 | 43.6 | 0.17 |
| HD 168137 | SB2 | O8.5 V | O7 V+O8 V | 0.75-0.8 | $>30$ | Undef. |
| HD 168183 | SB2OE | B0 III | O9.5 III + B3-5 V/III | 0.3 | 4.015 | 0.05 |
| Composite systems |  |  |  |  |  |  |
| $\begin{aligned} & \mathrm{BD}-13^{\circ} 4929 \\ & \mathrm{HD} 168076 \end{aligned}$ | SB3 <br> composite | $\begin{aligned} & \text { O9 V } \\ & \text { O4III ((f+)) } \end{aligned}$ | $\begin{aligned} & \mathrm{O} 7 \mathrm{~V}+(\mathrm{B} 0.5 \mathrm{~V}+\mathrm{B} 0.5 \mathrm{~V}) \\ & 03.5 \mathrm{~V}((\mathrm{f}+))+\mathrm{O} .5 \mathrm{~V} \end{aligned}$ | $\begin{aligned} & 1.0+(1.0) \\ & 0.43 \end{aligned}$ | $\begin{aligned} & \text { Undef. }+(\sim 4) \\ & >1.5 \quad 10^{5} ? \end{aligned}$ | Undef. Undef. |
| Single stars |  |  |  |  |  |  |
| BD $-13^{\circ} 4927$ | sgl. | O7 II(f) | O7 II(f) | $n / a$ | $n / a$ | $n / a$ |
| $\mathrm{BD}-13^{\circ} 4928$ | sgl. | O9.5 Vn | 09.5 V | $n / a$ | $n / a$ | $n / a$ |
| $\mathrm{BD}-13^{\circ} 4930$ | sgl. | O9.7 IIIp | O9.5 Vp | $n / a$ | $n / a$ | $n / a$ |

[^1]Table 10. New spectroscopic DM for the O stars in this work.

| Object | $E(B-V)$ | $V_{0}$ | $M_{\mathrm{V}}^{\text {th. }}$ | DM |
| :--- | :--- | :--- | :--- | :--- |
| $\mathrm{BD}-13^{\circ} 4923$ | $1.26 \pm 0.02$ | $5.29 \pm 0.15$ | -5.85 | $11.14 \pm 0.15$ |
| HD 168076 | $0.71 \pm 0.10$ | $5.52 \pm 0.38$ | -5.94 | $11.46 \pm 0.38$ |
| HD 168075 | $0.81 \pm 0.11$ | $5.74 \pm 0.41$ | -5.65 | $11.39 \pm 0.41$ |
| HD 168183 | $0.60 \pm 0.10$ | $5.94 \pm 0.38$ | -5.20 | $11.14 \pm 0.38$ |
| $\mathrm{BD}-13^{\circ} 4929$ | $0.98 \pm 0.21$ | $6.17 \pm 0.81$ | -5.24 | $11.41 \pm 0.81$ |
| HD 168137 | $0.68 \pm 0.03$ | $6.42 \pm 0.12$ | -5.33 | $11.75 \pm 0.12$ |
| $\mathrm{BD}-13^{\circ} 4927$ | $0.94 \pm 0.61$ | $7.59 \pm 2.34$ | -5.54 | $13.13 \pm 2.34$ |
| $\mathrm{BD}-13^{\circ} 4928$ | $0.76 \pm 0.05$ | $7.19 \pm 0.20$ | -3.90 | $11.09 \pm 0.20$ |
| $\mathrm{BD}-13^{\circ} 4930$ | $0.60 \pm 0.08$ | $7.11 \pm 0.31$ | -4.34 | $11.45 \pm 0.31$ |

\& McElroy (1984) for B stars. For multiple objects, we weighted the intrinsic colour excess $E(B-V)$ and the bolometric correction by the expected contributions of the components to the total flux. Errors were estimated by means of the error propagation, adopting a dispersion corresponding to half of a subspectral type on the parameters taken from the calibrations. The reddening law is taken to be $A_{\mathrm{V}}=3.75 E(B-V)$ (Hillenbrand et al. 1993). Table 10 lists the visual magnitude corrected for reddening $\left(V_{0}\right)$, the expected absolute magnitude given the component spectral types ( $M_{V}^{\text {th. }}$ ) and the corresponding distance moduli (DM). Given the estimated errorbars, the latter are in very good agreement with each other. A weighted average gives $\overline{D M}=11.4 \pm 0.1$. This corresponds to a distance to NGC 6611 of $d=1.9 \pm 0.1 \mathrm{kpc}$, in excellent agreement with the estimate of Dufton et al. (2006).

## 6 SUMMARY

Using a set of about 100 medium- to high-resolution spectra collected over several years, we revisited the properties of the current O-type star population of NGC 6611 in the Eagle Nebula, with a particular emphasis on their spectral classification and multiplicity.

We unveiled several new binaries and confirmed two suspected candidates. We proposed the first SB2 solutions for BD $-13^{\circ} 4923$ and HD 168183. We further identified several objects, possible longperiod systems, displaying a composite spectrum. We detected the secondary companion signature for all the detected binaries allowing us to put important constraints on the mass-ratio distribution and on the distribution of secondary masses. As a support to our analysis, we further present a set of Monte Carlo simulations that assess the binary detection biases of our campaign and show that the latter do not affect our conclusions.

The minimal binary fraction is $f_{\min } \approx 0.44$ but can be as large as 0.67 if the two composite objects are confirmed to be gravitationally bound. In all cases, our current data set excludes a 100 per cent spectroscopic binary fraction in NGC 6611. Still, up to 75 per cent of the total O-type star population is to be found in multiple systems. The period and mass-ratio distributions obtained are both compatible with those derived in Paper I for NGC 6231, although the two samples suffer from low-number statistics. All the detected companions are O or early- to mid- B stars, suggesting again that the secondary companion of an O-type star cannot be randomly drawn from a classical mass function but has a mass strongly biased towards the upper range of the mass spectrum.

Finally, we could not detect any significant difference in the multiplicity properties of the O star population of NGC 6611 and 6231. The environmental properties in which those two populations evolved are however significantly different. On the one hand, NGC 6611 is likely very young, about $1-2$ Myr old as attested by the
presence of two O3-4 stars but no Wolf-Rayet object, and still very dusty. On the other hand, NGC 6231 is significantly older (around 3-4 Myr) and has already cleared out its dust.

The results presented here and the significant revision of the status of several objects despite the relative proximity of the cluster and the abundant literature on its O-star population shows the need to pursue our effort on several other nearby clusters. Only upon completion of our campaign will we be able to draw a global picture of the properties of the O-type star population in the nearby Universe.

## ACKNOWLEDGMENTS

The authors are grateful to R. Barbá for sharing results before publication and to the referee, D. Gies, for a careful revision of the manuscript. This paper relies on data taken at the La SillaParanal Observatory under program IDs 71.C-0513, 171.D-0237, 073.D-0234, 073.D-0609, 075.D-0061, 075.D-0369, 077.D-0146, 079.D-0564 and 082.D-0136(A). This paper also made use of the ADS, of the SIMBAD and WEBDA data bases and of the Vizier catalogue access tool (CDS, Strasbourg, France).

## REFERENCES

Barbá R. H., Gamen R., Arias J. I., Morrell N., Màiz Apellàniz J., Alfaro E., Walborn N., Sota A., 2009, in Rivinius T., Curé M., eds, The Interferometric View of Hot Stars. Rev. Mex. Conf. Ser., Mexico, in press
Bevington P., 1969, Data Reduction and Error Analysis for the Physical Sciences. McGraw-Hill Book Company, Inc., USA
Bonatto C., Santos J. F. C. Jr., Bica E., 2006, A\&A, 445, 567
Bosch G. L., Morrell N. I., Niemelä V. S., 1999, Rev. Mex. Astron. Astrofis., 35, 85
Conti P. S., 1973, ApJ, 179, 181
Conti P. S., Alschuler W. R., 1971, ApJ, 170, 325
Conti P. S., Leep E. M., Lorre J. J., 1977, ApJ, 214, 759
Duchêne G., Simon T., Eislöffel J., Bouvier J., 2001, A\&A, 379, 147
Dufton P. L. et al., 2006, A\&A, 457, 265
Evans C. J. et al., 2005, A\&A, 437, 467
Gosset E., Royer P., Rauw G., Manfroid J., Vreux J.-M., 2001, MNRAS, 327, 435
Guarcello M. G., Prisinzano L., Micela G., Damiani F., Peres G., Sciortino S., 2007, A\&A, 462, 245

Gvaramadze V. V., Bomans D. J., 2008, A\&A, 490, 1071
Heck A., Manfroid J., Mersch G., 1985, A\&AS, 59, 63
Hester J. J. et al., 1996, AJ, 111, 2349
Hillenbrand L. A., Massey P., Strom S. E., Merrill K. M., 1993, AJ, 106, 1906
Hipparcos Catalogue 1997, in Perryman M. A. C., ESA Special Publication Vol. 1200. ESA, Noordwijk

Hobbs L. M. et al., 2008, ApJ, 680, 1256
Humphreys R. M., McElroy D. B., 1984, ApJ, 284, 565
Hunter I. et al., 2007, A\&A, 466, 277
Martayan C., Floquet M., Hubert A. M., Neiner C., Frémat Y., Baade D., Fabregat J., 2008, A\&A, 489, 459
Martins F., Plez B., 2006, A\&A, 457, 637
Martins F., Schaerer D., Hillier D., 2005, A\&A, 436, 1049
Mathys G., 1988, A\&AS, 76, 427
Mathys G., 1989, A\&AS, 81, 237
Oliveira J. M., 2008, in Reipurth B., ed., ASP Monograph Publications, Handbook of Star Forming Regions, Vol. V. Astron. Soc. Pac., San Francisco, p. 599
Sana H., 2009, A\&A, 501, 291
Sana H., Gosset E., 2009, A\&A, submitted
Sana H., Gosset E., Nazé Y., Rauw G., Linder N., 2008, MNRAS, 386, 447 (Paper I)
Underhill A. B., 1994, ApJ, 420, 869
Walborn N. R., Fitzpatrick E. L., 1990, PASP, 102, 379

Walborn N. R. et al., 2002, AJ, 123, 2754
Walker M. F., 1961, ApJ, 133, 438
Wolff S. C., Strom S. E., Dror D., Venn K., 2007, AJ, 133, 1092

## SUPPORTING INFORMATION

Additional Supporting Information may be found in the online version of this article:
Table 3. Journal of the spectroscopic observations of the O-type stars in NGC 6611.

Please note: Wiley-Blackwell are not responsible for the content or functionality of any supporting information supplied by the authors. Any queries (other than missing material) should be directed to the corresponding author for the article.

This paper has been typeset from a $\mathrm{T}_{\mathrm{E}} \mathrm{X} / \mathrm{E} \mathrm{T}_{\mathrm{E}} \mathrm{X}$ file prepared by the author.


[^0]:    *E-mail: hsana@eso.org
    $\dagger$ FNRS, Belgium.

[^1]:    ${ }^{a}$ Orbital parameters for HD 168075 are taken from Barbá et al. (2009).

