

## The phenomenon

Droplet coalescence in a bath can be delayed by oscillating the bath vertically with an amplitude  $A$  and a frequency  $f$  between 20 Hz to 400 Hz : the droplet bounces on the interface [1,2]. The droplet deformation is enhanced by considering a low viscosity droplet on a high viscosity bath.

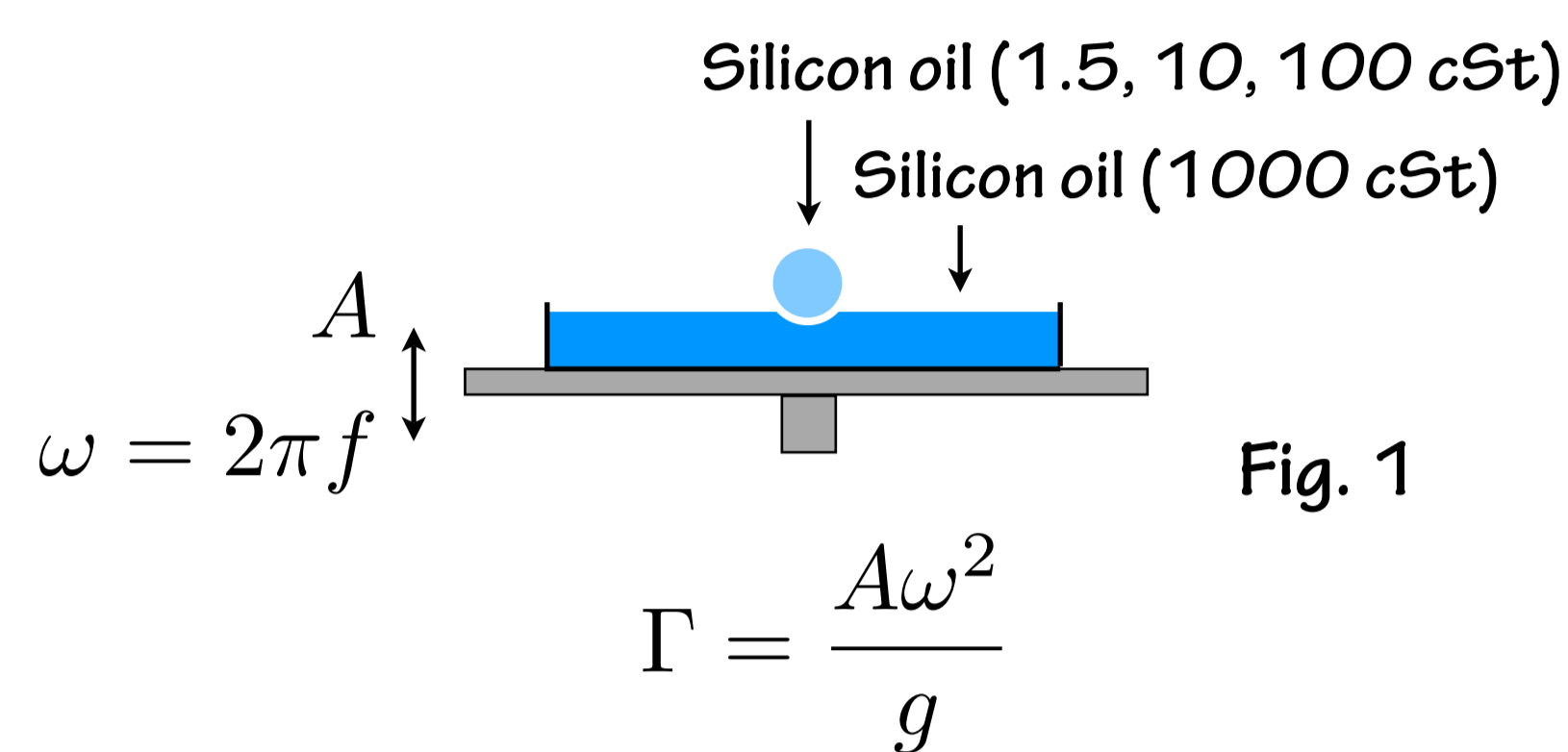


Fig. 1

## Bouncing threshold

The bouncing of the droplet only occurs when the reduced acceleration  $\Gamma$  is higher than a threshold value  $\Gamma_{th}$  which depends among other on the forcing frequency  $f$ , the droplet radius  $R$ , and the viscosity  $\nu$ . In Fig. 2, the threshold  $\Gamma_{th}$  is represented as a function of the forcing frequency for various droplets viscosities. For high viscosity droplets,  $\Gamma_{th}$  increases monotonically and scales as  $f^2$  [1]. For low viscosity droplets,  $\Gamma_{th}(f)$  presents extrema that suggest a resonance phenomenon.

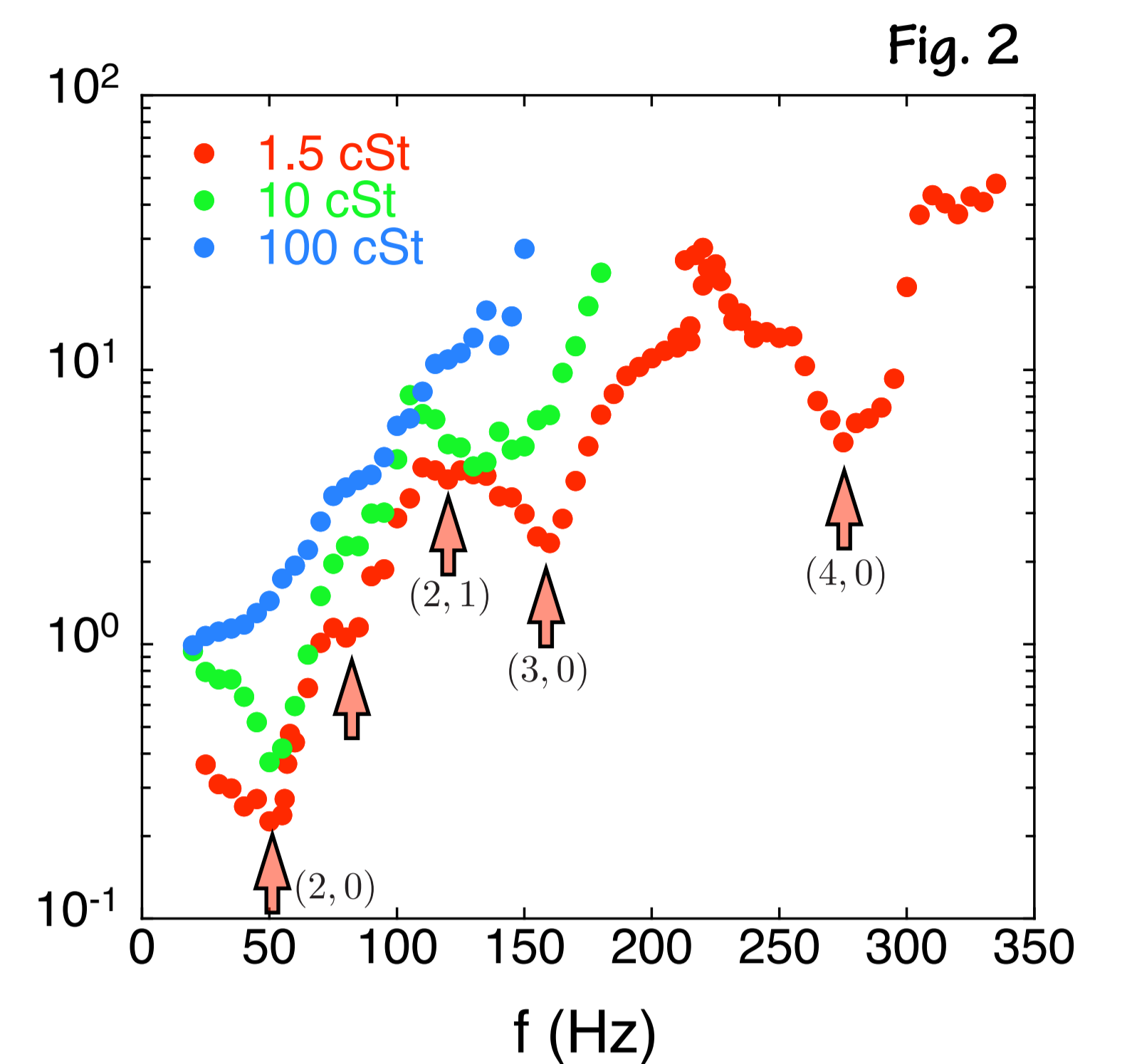
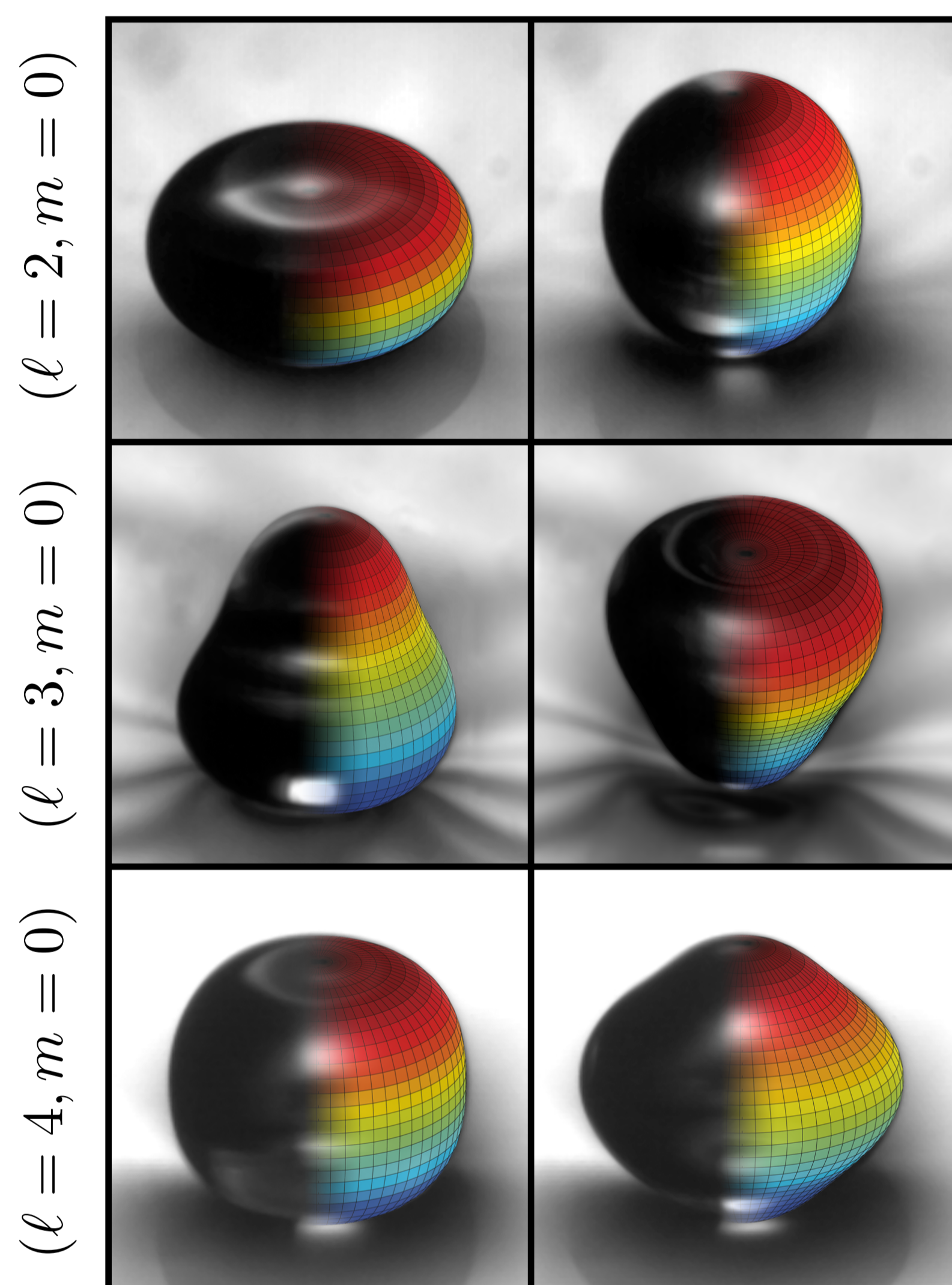


Fig. 2

## Resonant modes



Spherical harmonic solution superposed to the experimental pictures.

Fig. 3

Specific modes are observed at minima of the  $\Gamma_{th}(f)$  curve. They are analogous to the natural modes of deformation expressed by Rayleigh in terms of spherical harmonics  $Y_{l,m}$ . Modes  $m = 0$  and  $\ell = 2, 3$  and  $4$  are observed (Fig. 3).

The droplet may be considered as a damped driven harmonic oscillator :  
 - surface tension is the restoring force,  
 - viscosity is the damping process.

Natural frequencies scale as the capillary frequency  $f_c$  :

$$f_c = \sqrt{\sigma/m} \quad \text{with } m = 4\pi/3\rho R^3$$

More precisely, the dispersion relation specifies the natural "Rayleigh" frequency  $f_R$  related to a  $\ell$ -mode :

$$\left(\frac{f_R(\ell)}{f_c}\right)^2 = \frac{1}{3\pi}\ell(\ell-1)(\ell+2)$$

Usually,  $f_{Resonance}(\ell) = a f_R(\ell)$ ;  $a$  being a multiplication factor which depends on the geometry of the excitation [4, 5].

When  $f = a f_R(\ell)$ , the droplet stores the whole energy provided by the oscillating bath in deformation and dissipates it due to enhanced internal motion : the droplet resonates. It is impossible to make the droplet bouncing in the mode  $(\ell, m=0)$ ,  $\Gamma_{th}$  should diverge. Experimentally, a maximum is found at  $f = a f_R(\ell)$  as  $m \neq 0$  modes are excited.

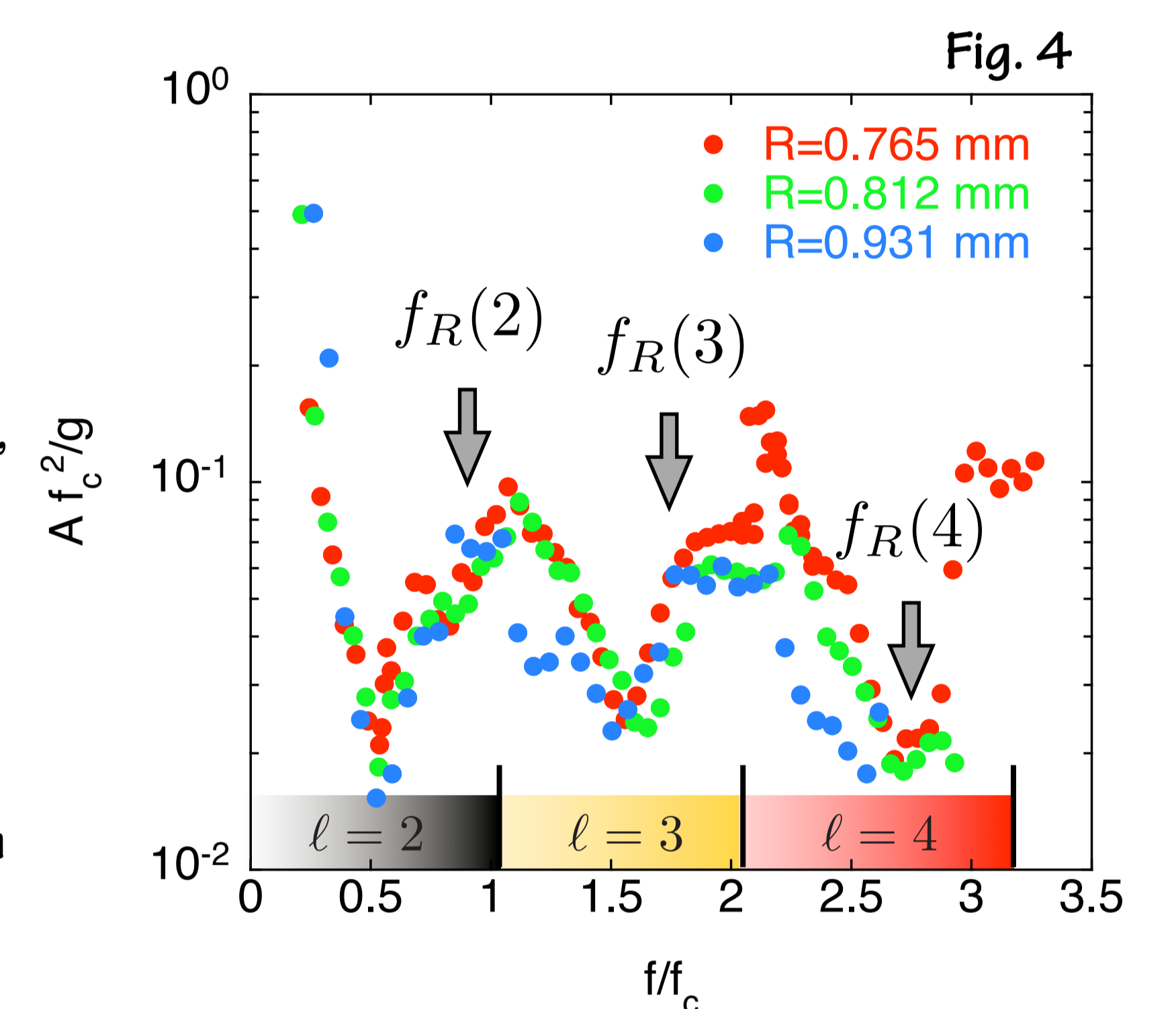
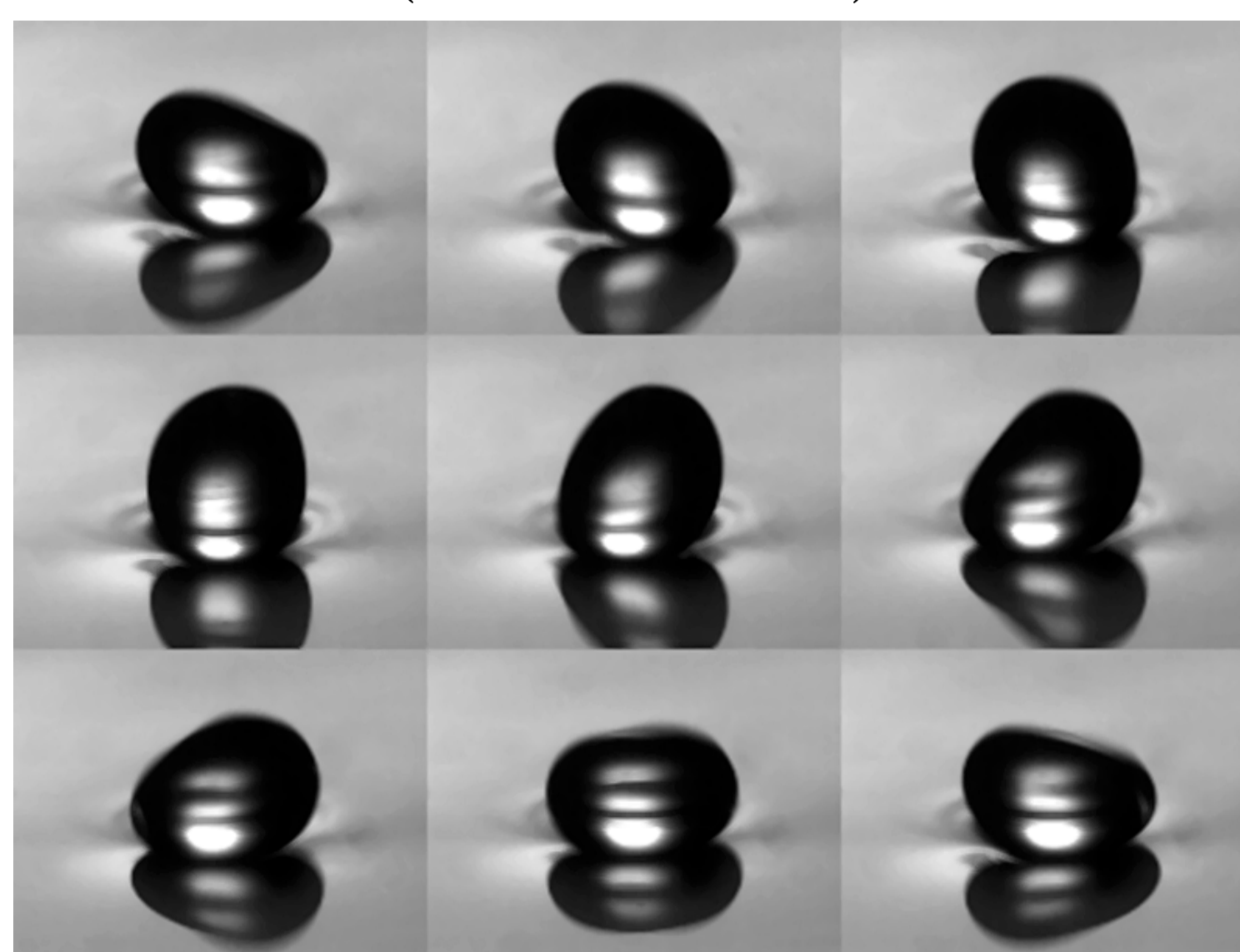


Fig. 4

## Displacement mode : The Roller

$(\ell = 2, m = 1)$



Snapshots of a Roller

Fig. 5

At the first maximum of the  $\Gamma_{th}(f)$  curve (115 Hz for 1.5 cSt droplet), the droplet moves along a linear trajectory. The mode of deformation, related to  $\ell = 2$  and  $m = 1$ , induces the internal rotation of the liquid.

The initial horizontal speed  $v$  of the droplet has been measured for various frequencies and amplitudes of the bath. The phenomenon only occurs above a cut-off frequency  $f_0 = 103$  Hz and an amplitude threshold  $A_0$ .

The speed  $v$  scales with  $(A-A_0)(f-f_0)$ , the reduced maximum speed of the bath.

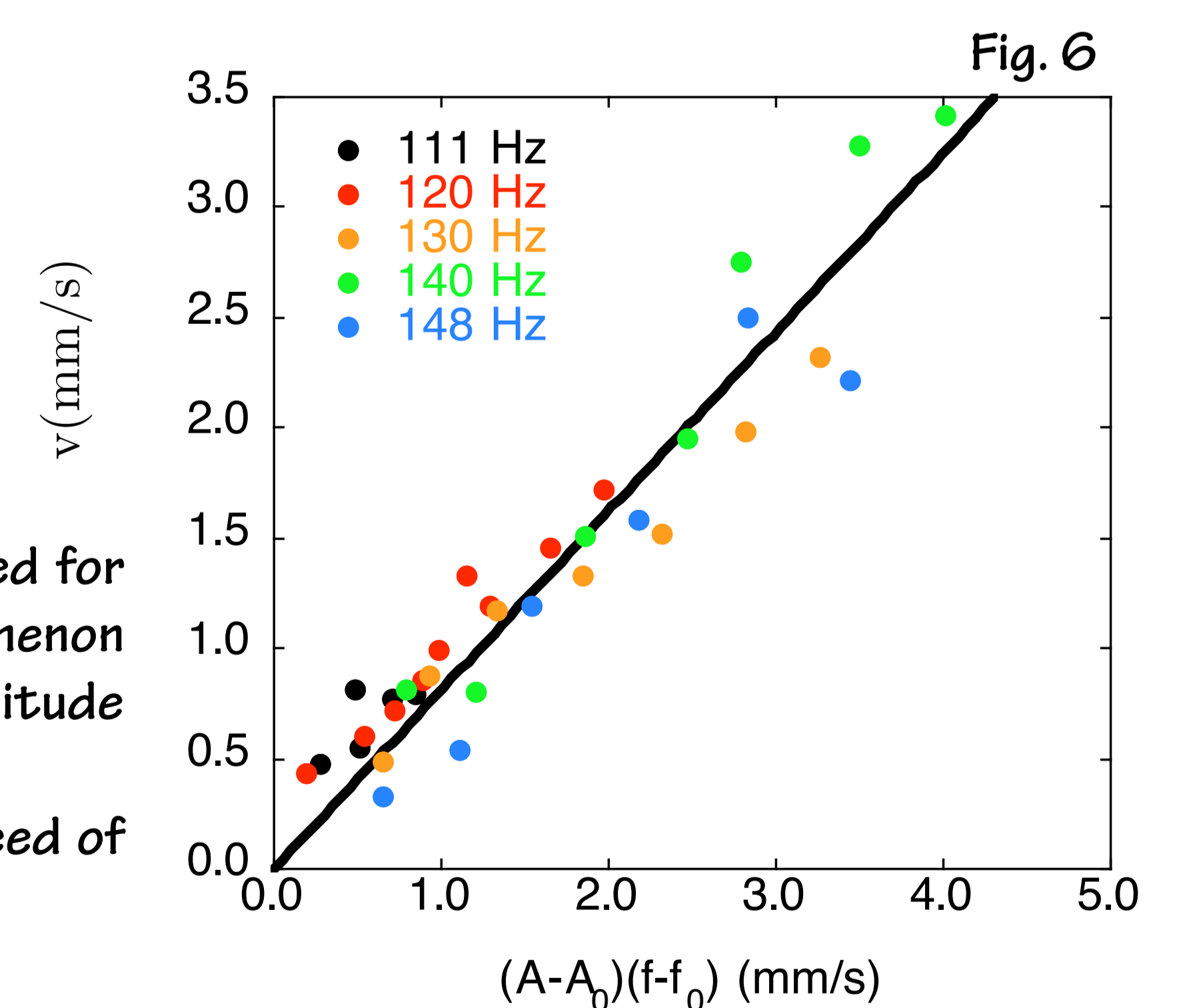


Fig. 6

## Conclusion

1. Extrema of the bouncing threshold curve of low viscosity droplets are related to a resonance phenomenon.
2. A model which explains the first minima for the mode of deformation  $(\ell = 2, m = 0)$  has been developed in [3].
3. A new mode of displacement for low viscosity droplets has been discovered that can be generalised to a wide range of droplet sizes, the Roller. This self-propelled mode allows manipulation without contact.

## References

- [1] Y. Couder, E. Fort, C. H. Cautier and A. Boudaoud, Phys. Rev. Lett. 94, 177801 (2005).
- [2] N. Vandewalle, D. Terwagne, K. Mulleners, T. Gilet and S. Dorbolo, Phys. Fluids 18, 091106 (2006).
- [3] T. Gilet, D. Terwagne, N. Vandewalle and S. Dorbolo, Phys. Rev. Lett. 100, 167802 (2008).
- [4] S. Courty, G. Lagubeau and T. Tixier, Phys. Rev. E 73, 045301(R) (2006).
- [5] X. Noblin, A. Buguin and F. Brochard-Wyart, Eur. Phys. J. E 14, 395-40 (2004).